

**SOMMERSEMESTER 2016 - HÖHERE ANALYSIS II  
LINEARE PARTIELLE DIFFERENTIALGLEICHUNGEN**

**Homework #6 Solutions**

**Problem 1.** Consider the Neumann problem for the Laplace operator in a bounded domain of class  $C^1$ . Define the Green function  $N(x, y)$  (also sometimes called the Neumann function) and give a formal solution formula for Neumann problem

$$\begin{aligned} -\Delta u &= f \in L_2(\Omega) , \\ \partial_\nu u \Big|_{\partial\Omega} &= g \in L_2(\partial\Omega) . \end{aligned}$$

If  $\Psi(x)$  is a fundamental solution for the Laplace operator, and  $u \in C^2(\bar{\Omega})$ , then according to Proposition 4.5 we have with  $G(x, y) = \Psi(x - y)$  that

$$u(x) = - \int_{\Omega} G(x, y) \Delta u(y) dy - \int_{\partial\Omega} \partial_{\nu(y)} G(x, y) u(y) dS(y) + \int_{\partial\Omega} G(x, y) \partial_{\nu(y)} u(y) dS(y) .$$

In the case of the Neumann problem we do not have any information about the function  $u$  on the boundary whereas  $\Delta u$  in  $\Omega$  and  $\nu \cdot \nabla u$  on  $\partial\Omega$  are known. Hence, the Neumann function  $N(x, y)$  is a fundamental solution which satisfies the condition that

$$\nu(y) \cdot \nabla_y N(x, y) = 0 \quad \text{for all } x \in \Omega, y \in \partial\Omega .$$

Note that if this condition is satisfied, the formula above reads as

$$u(x) = - \int_{\Omega} N(x, y) f(y) dy + \int_{\partial\Omega} N(x, y) g(y) dS(y) .$$

*Construction of the Neumann function.* Let  $\Phi$  be the fundamental solution introduced in Definition 4.1. Suppose that the function  $H(x, y)$  satisfies the following Neumann boundary value problem

$$\begin{aligned} \Delta_y H(x, y) &= 0 \quad \text{for all } x, y \in \Omega, \\ \partial_{\nu(y)} H(x, y) &= -\partial_{\nu(y)} \Phi(x - y) \quad \text{for all } y \in \partial\Omega, x \in \Omega . \end{aligned}$$

Then the Neumann function is given by  $N(x, y) = \Phi(x - y) + H(x, y)$ .

**Problem 2.** Derive the Green function  $G(x, y)$  for the Dirichlet problem for the Laplacian in the ball  $B_R(0)$  ( $R > 0$  fixed) in the case  $d \geq 3$  and derive a formula for the solution of the Dirichlet problem

$$-\Delta u = 0 \text{ in } B_R(0) \quad u = g \in C(\partial B_R(0)) .$$

State a Theorem similar to Theorem 4.8 and prove it. Hint: For  $x \neq 0$  set

$$G(x, y) = \frac{1}{d(d-2)\omega_d} [ |x-y|^{2-d} - a|bx-y|^{2-d} ]$$

and choose  $a \in \mathbb{R}$  and  $b > R/|x|$  such that  $G(x, y) = 0$  for all  $0 \neq x \in B_R(0)$  and  $y \in \partial B_R(0)$ .

*Solution.* Note that for  $x \in B_r(0)$  the function  $G(x, y)$  satisfies  $-\Delta_y G(x, y) = \delta_x$  as long as  $bx \notin B_R(0)$ . Fix  $x \in B_R(0) \setminus \{0\}$ . Then  $G(x, y) = 0$  holds if for all  $|y| = R$

$$a^{2/(2-d)}[b^2|x|^2 - 2bx \cdot y + R^2] = |x|^2 - 2x \cdot y + R^2 .$$

Choose now  $a^{(2-d)/2} = b$ . Then the equation is true if and only if

$$b^2|x|^2 + R^2 = b(|x|^2 + R^2) \quad \text{or} \quad b(b-1)|x|^2 = (b-1)R^2 .$$

Hence,  $b = R^2/|x|^2$  and the Green function for the ball is

$$\begin{aligned} G(x, y) &= \frac{1}{d(d-2)\omega_d} \left[ |x-y|^{2-d} - \left( \frac{|x|}{R} \right)^{2-d} \left| \frac{R^2}{|x|^2}x - y \right|^{2-d} \right] \\ &= \frac{1}{d(d-2)\omega_d} \left[ |x-y|^{2-d} - \left| \frac{R}{|x|}x - \frac{|x|}{R}y \right|^{2-d} \right] \\ &= \frac{1}{d(d-2)\omega_d} \left[ |x-y|^{2-d} - (R^2 - 2x \cdot y + |x|^2|y|^2/R^2)^{2-d} \right] , \end{aligned}$$

which is well-defined also for  $x = 0$ .

Compute now, for  $x \in B_R(0)$  and  $|y| = R$ , using the formula for the gradient of the fundamental solution given before Theorem 4.2,

$$\begin{aligned} \nu(y) \cdot \nabla_y G(x, y) &= \frac{1}{d\omega_d} \left[ \frac{y \cdot x - y \cdot y}{R|x-y|^d} - \frac{y}{R} \left( \frac{|x|}{R} \right)^{2-d} \frac{R^2 x/|x|^2 - y}{|R^2 x/|x|^2 - y|^d} \right] \\ &= \frac{1}{d\omega_d} \left[ \frac{y \cdot x - R^2}{R|x-y|^d} - \frac{|x|^d}{R^{3-d}} \frac{R^2 x \cdot y - R^2|x|^2}{|R^2 x - y|^d} \right] \\ &= \frac{1}{d\omega_d} \left[ \frac{y \cdot x - R^2}{R|x-y|^d} - \frac{1}{R} \frac{x \cdot y - |x|^2}{|x-y|^d} \right] = \frac{1}{d\omega_d R} \frac{|x|^2 - R^2}{|x-y|^d} \end{aligned}$$

**Theorem.** Suppose that  $g \in C(\partial B_R(0))$ . Then the function

$$u(x) = \frac{R^2 - |x|^2}{d\omega_d R} \int_{\partial B_R(0)} \frac{g(y)}{|x-y|^d} dS(y)$$

has the regularity  $u \in C^\infty(B_R(0)) \cap L_\infty(B_R(0))$ , satisfies the equation  $-\Delta u = 0$  in  $B_R(0)$ , and  $u(x) = g(x)$  for all  $x \in \partial B_R(0)$  in the sense that for all  $\underline{x} \in \partial B_R(0)$

$$\lim_{x_l \rightarrow \underline{x}} u(x_l) = g(\underline{x})$$

for all sequences  $\{x_l\} \subset B_R(0)$  converging to  $\underline{x} \in \partial B_R(0)$ .

*Proof.* For all  $x \in B_R(0)$  one verifies that  $\partial_j u(x)$  can be computed by the differentiation into the integral which can be done since  $x \neq y$ . This can be done for all derivatives of higher order as well and gives  $u \in C^\infty(B_R(0))$ . Next compute for  $x \in B_R(0)$  and  $|y| = R$ ,

using the quotient rule of differential calculus,

$$\begin{aligned} \Delta_x \frac{|x|^2 - R^2}{|x - y|^d} &= \nabla_x \cdot \nabla_x \frac{|x|^2 - R^2}{|x - y|^d} = \nabla_x \cdot \frac{2x|x - y|^d - d(|x|^2 - R^2)|x - y|^{d-2}(x - y)}{|x - y|^{2d}} \\ &= \nabla_x \cdot \frac{2x}{|x - y|^d} - \nabla_x \cdot \frac{d(|x|^2 - R^2)(x - y)}{|x - y|^{d+2}} \\ &= \frac{2d}{|x - y|^d} - \frac{2d(|x|^2 - x \cdot y)}{|x - y|^{d+2}} - \frac{2dx \cdot (x - y) - d^2(|x|^2 - R^2)}{|x - y|^{d+2}} + \frac{d(d + 2)(|x|^2 - R^2)}{|x - y|^{d+2}} \\ &= \frac{2d(x - y)^2}{|x - y|^{d+2}} - \frac{4d[|x|^2 - x \cdot y] - 2d(|x|^2 - R^2)}{|x - y|^{d+2}} = \frac{2d[|y|^2 - R^2]}{|x - y|^{d+2}} = 0. \end{aligned}$$

This shows that  $-\Delta u = 0$ , pointwise in  $B_R(0)$ . Next we will show that  $u \in L_\infty(\Omega)$ . For that purpose set

$$K(x, y) = \frac{1}{d\omega_d R} \frac{R^2 - |x|^2}{|x - y|^d}$$

and observe that for  $K > 0$  for  $x \in B_R(0)$ . Then, choosing  $g \equiv 1$  and using Proposition 4.5 with  $\Omega = B_R(0)$  we obtain

$$\int_{\partial B_R(0)} K(x, y) dS(y) = 1$$

From here the proof is completed as the proof of Theorem 4.7. □

A perhaps simpler way of proving  $\Delta_x K(x, y) = 0$  for  $x \in B_R(0)$  and  $|y| = R$  is based on the fact that  $G(x, y) = G(y, x)$  whenever  $x \neq y$  and that  $K(x, y) = \partial_{\nu(y)} G(x, y)$  for  $y \in \partial B_R(0)$ .

**Problem 3.** *The chain rule for Sobolev functions.* Suppose that  $f \in C^1(\mathbb{R}) \cap W_\infty^1(\mathbb{R})$  and let  $u \in W_p^1(\Omega)$  for some  $p \in [1, \infty)$  where  $\Omega \subset \mathbb{R}^d$  is open and bounded. Prove that  $f \circ u \in W_p^1(\Omega)$  and  $\partial_j(f \circ u) = f'(u)\partial_j u$  in  $L_p(\Omega)$  for  $j = 1, \dots, d$ . Does this chain rule also hold in the case that  $p = \infty$ ?

*Proof.* Fix  $p \in [1, \infty)$ . Note that  $f \circ u \in L_p(\Omega)$  since  $f \circ u \in L_\infty(\Omega)$  and  $\Omega$  is bounded. By Theorem 3.13 there exists a sequence  $u_m \in C^\infty(\Omega) \cap W_p^1(\Omega)$  such that  $\|u_m - u\|_{W_p^1(\Omega)} \rightarrow 0$  as  $m \rightarrow \infty$ . Note that this convergence implies that  $u_m \rightarrow u$  almost everywhere. The chain rule is true for differentiable functions, hence we have

$$\partial_j(f \circ u) = f'(u_m)\partial_j u_m \quad \text{for } m = 1, 2, \dots \text{ and } j = 1, 2, \dots, d.$$

In what follows the positive integer  $j \leq d$  is fixed. The sequence  $f \circ u_m \in W_p^1(\Omega)$  converges to  $f \circ u$  in  $L_p(\Omega)$ . Indeed, by the mean value theorem of differential calculus

$$\|f \circ u - f \circ u_m\|_{L_p(\Omega)}^p \leq \max_{\xi \in \mathbb{R}} |f'(\xi)| \int_\Omega |u(x) - u_m(x)| dx.$$

Convergence in  $L_p(\Omega)$  implies convergence in  $\mathcal{D}'(\Omega)$ . By Proposition 3.7 we know that

$$\partial_j(f \circ u_m) \longrightarrow \partial_j(f \circ u) \quad \text{in } \mathcal{D}'(\Omega).$$

Furthermore, we know that  $f'(u)\partial_j u \in L_p(\Omega)$  since  $f'(u) \in L_\infty(\Omega)$ . Using the triangle inequality in  $L_p(\Omega)$  we obtain

$$\begin{aligned} & \left( \int_{\Omega} |f'(u_m)\partial_j u_m - f'(u)\partial_j u|^p dx \right)^{1/p} \\ &= \left( \int_{\Omega} |f'(u_m)\partial_j(u_m - u) + (f'(u_m) - f'(u))\partial_j u|^p dx \right)^{1/p} \\ &\leq \left( \int_{\Omega} |f'(u_m)\partial_j(u_m - u)|^p dx \right)^{1/p} + \left( \int_{\Omega} |(f'(u_m) - f'(u))\partial_j u|^p dx \right)^{1/p} \\ &\leq \|f'(u_m)\|_{L_\infty(\Omega)} \|u_m - u\|_{W_p^1(\Omega)} + \left( \int_{\Omega} |(f'(u_m) - f'(u))\partial_j u|^p dx \right)^{1/p} \longrightarrow 0 \end{aligned}$$

as  $m \rightarrow \infty$  where the continuity of  $f'$  and the Lebesgue dominated convergence theorem has been used. Here the observation that the continuity of  $f'$  implies that  $f'(u_m) \rightarrow f'(u)$  almost everywhere in  $\Omega$  is crucial. This shows that

$$f'(u_m)\partial_j u_m \longrightarrow f'(u)\partial_j u \quad \text{in } L_p(\Omega)$$

and by the uniqueness of the limit one obtains

$$f'(u)\partial_j u = \partial_j(f \circ u) \in L_p(\Omega).$$

□

If  $u \in W_\infty^1(\Omega)$ , then there does not need to exist a sequence of smooth functions converging to  $u$  in  $W_\infty^1(\Omega)$ . However, since  $\Omega$  is bounded, we know that  $u \in W_p^1(\Omega)$  for any  $p \in [1, \infty)$  and we can use the statement proved above to state that  $\partial_j(f \circ u) = f'(u)\partial_j u$  in  $L_p(\Omega)$ . However, the right-hand side is also in  $L_\infty(\Omega)$  and that shows that the chain rule holds also for compositions  $f \circ u$  with  $f \in C^1(\mathbb{R}) \cap W_\infty^1(\mathbb{R})$  and  $u \in W_\infty^1(\Omega)$ .

An interesting application of this chain rule pertains to problem 4.1 and results in a proof of the divergence theorem for bounded regions whose boundary is locally the graph of a Lipschitz function.