

**SOMMERSEMESTER 2016 - HÖHERE ANALYSIS II  
LINEARE PARTIELLE DIFFERENTIALGLEICHUNGEN**

**Homework #8** due 06/14

**Problem 1.** Consider the Dirichlet problem for the Laplace equation in the unit disk, that is

$$-\Delta u = 0 \quad \text{in } B_1(0) \subset \mathbb{R}^2, \quad u = g \quad \text{in } \partial B_1(0).$$

Suppose that  $g$  can be expanded into a Fourier series  $g(\phi) = \sum_{k=0}^{\infty} a_k \cos(k\phi)$  and look for a solution of the form

$$u(r, \phi) = \sum_{k=0}^{\infty} b_k(r) \cos(k\phi),$$

where  $(r, \phi)$  are polar coordinates.

a.) Show that the  $b_k$  have to satisfy the equation

$$b_k''(r) + \frac{1}{r} b_k'(r) - \frac{k^2}{r^2} b_k = 0 \quad \text{for } k = 1, 2, \dots$$

Solve this differential equation. Since it is of second order, there are two linearly independent solutions. However, one has to be discarded. Why?

*Solution.* Recall that the Laplacian in polar coordinates is given by

$$\Delta = \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \phi^2}.$$

Differentiating the series  $u(r, \phi)$  formally gives

$$\Delta u = \sum_{k=0}^{\infty} \left[ b_k''(r) + \frac{1}{r} b_k'(r) - \frac{k^2}{r^2} b_k(r) \right] \cos(k\phi)$$

Since the functions  $\cos(k\phi)$  are an orthogonal set in  $L_2(0, 2\pi)$ , the series vanishes only if each  $b_k$  satisfies the differential equation

$$b_k''(r) + \frac{1}{r} b_k'(r) - \frac{k^2}{r^2} b_k(r) = 0.$$

The general solutions of this equation are  $b_0(r) = c_1 + c_2 \log r$  and  $b_k(r) = c_1 r^k + c_2 r^{-k}$  for  $k = 1, 2, \dots$ . In each case we set  $c_2 = 0$  and discard the unbounded solution. We obtain

$$(1) \quad u(r, \phi) = \sum_{k=0}^{\infty} a_k r^k \cos(k\phi).$$

Given that the series for  $g$  is convergent in the  $L_2$ -sense, that is  $\sum_{k=0}^{\infty} a_k^2 < \infty$ , the series for  $u(r, \phi)$  is absolutely convergent for all  $r < 1$ .

b.) Give a condition on the sequence  $\{a_k\}$  which guarantees that the energy  $\int_{B_1(0)} |\nabla u|^2 dx$

of the series solution is finite.

*Solution.* Observe that

$$|\nabla u|^2 = |\partial_r u|^2 + \frac{1}{r^2} |\partial_\phi u|^2 .$$

Hence, using the orthogonality of the cosine and sine functions and  $\int_0^{2\pi} \cos^2(k\phi) d\phi = \int_0^{2\pi} \sin^2(k\phi) d\phi = \pi$  for  $k = 1, 2, \dots$ , one obtains

$$\begin{aligned} \int_{B_1(0)} |\nabla u|^2 dx &= \int_0^{2\pi} \int_0^1 \left[ |\partial_r u|^2 + \frac{1}{r^2} |\partial_\phi u|^2 \right] r dr d\phi \\ &= \pi \sum_{k=1}^{\infty} a_k^2 k^2 \int_0^1 r^{2k-1} dr + \pi \sum_{k=1}^{\infty} a_k^2 k^2 \int_0^1 r^{2k-1} dr = \pi \sum_{k=1}^{\infty} a_k^2 k . \end{aligned}$$

This integral is finite if the sum  $\sum_{k=1}^{\infty} a_k^2 k < \infty$ .

c\*.) Can you find a  $g \in C(\partial B_1(0))$  such that the corresponding solution  $u$  does not have finite energy ?

**Problem 2.** Use the series expansion derived in Problem 1 to obtain the Poisson integral formula in  $d = 2$ . In higher dimensions this formula has been discussed in Homework 6, Problem 2.

*Solution.* If

$$g(\phi) = \sum_{k=0}^{\infty} a_k \cos(k\phi)$$

then

$$a_0 = \frac{1}{2\pi} \int_0^{2\pi} g(\theta) d\theta \quad \text{and} \quad a_k = \frac{1}{\pi} \int_0^{2\pi} g(\theta) \cos(k\theta) d\theta \quad \text{for } k \geq 1 ,$$

and

$$\int_0^{2\pi} g(\theta) \sin(k\theta) d\theta = 0 \quad \text{for } k = 1, 2 .$$

Moving these expression into formula one gives

$$\begin{aligned} u(r, \phi) &= \frac{1}{2\pi} \int_0^{2\pi} g(\theta) d\theta + \frac{1}{\pi} \sum_{k=1}^{\infty} r^k \left[ \cos(k\phi) \int_0^{2\pi} g(\theta) \cos(k\theta) d\theta + \sin(k\phi) \int_0^{2\pi} g(\theta) \sin(k\theta) d\theta \right] \\ &= \frac{1}{2\pi} \int_0^{2\pi} g(\theta) d\theta + \frac{1}{\pi} \sum_{k=1}^{\infty} r^k \int_0^{2\pi} g(\theta) \cos(k(\theta - \phi)) d\theta \\ &= \frac{1}{2\pi} \int_0^{2\pi} g(\theta) d\theta + \frac{1}{2\pi} \int_0^{2\pi} \sum_{k=1}^{\infty} r^k [e^{ik(\theta-\phi)} + e^{-ik(\theta-\phi)}] g(\theta) d\theta \\ &= \frac{1}{2\pi} \int_0^{2\pi} g(\theta) \left[ \frac{1}{1 - re^{i(\theta-\phi)}} + \frac{re^{-i(\theta-\phi)}}{1 - re^{-i(\theta-\phi)}} \right] d\theta \\ &= \frac{1}{2\pi} \int_0^{2\pi} g(\theta) \frac{1 - re^{-i(\theta-\phi)} + (1 - re^{i(\theta-\phi)})re^{-i(\theta-\phi)}}{1 - 2r \cos(\theta - \phi) + r^2} d\theta \\ &= \frac{1}{2\pi} \int_0^{2\pi} g(\theta) \frac{1 - r^2}{1 - 2r \cos(\theta - \phi) + r^2} d\theta \end{aligned}$$

Returning to Cartesian coordinates gives

$$u(x) = \frac{1 - |x|^2}{2\pi} \int_S \frac{g(y)}{|x - y|^2} dS(y) .$$

which is the Poisson formula for the unit circle  $S$ .

**Problem 3.** Consider the function  $f(x) = \chi_{[-a,a]} \in L_2(\mathbb{R})$ .

a.) Compute the distributional derivative  $f'(x)$  and show that  $\Delta^h f \rightarrow 0$  almost everywhere.

*Solution.* Let  $\varphi \in C_0^\infty(\mathbb{R})$ . Then

$$\frac{d}{dx} f(\varphi) = - \int_a^a \varphi'(x) dx = -\varphi(a) + \varphi(-a)$$

and hence  $f' = \delta_{-a} + \delta_a$ . Note that  $f' = 0$  almost everywhere. For all  $x \in \mathbb{R} \setminus \{a, -a\}$  we have

$$\Delta^h f(x) = \frac{f(x+h) - f(x)}{h} = 0$$

as soon as  $|h| < \min\{|x-a|, |x+a|\}$ . Hence,  $\Delta^h f \rightarrow 0$  almost everywhere.

b.) Show that the difference quotients  $\Delta^h f$  do not converge to zero in  $L_2(\mathbb{R})$ .

*Solution.* Suppose that  $h > 0$ . Then

$$\Delta^h f(x) = \begin{cases} 0 & \text{for } x > a \\ -1/h & \text{for } a-h < x \leq a \\ 0 & \text{for } -a \leq x \leq a-h \\ 1/h & \text{for } -a-h \leq x < -a \\ 0 & \text{for } x < -a-h \end{cases}$$

and

$$\int_{\mathbb{R}} |\Delta^h f(x)|^2 dx = \int_{a-h}^a \frac{1}{h^2} dx + \int_{-a-h}^{-a} \frac{1}{h^2} dx = \frac{2}{h}$$

which proves that  $\Delta^h f$  does not converge to zero in  $L_2(\mathbb{R})$ . A similar computation can be done for  $h < 0$ .