## Permutohedra for knots and quivers



## Piotr Sułkowski

Faculty of Physics, University of Warsaw
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Jakub Jankowski ${ }^{1,3}$, Piotr Kucharski ${ }^{2,3}$, Hélder Larraguível ${ }^{3}$, Dmitry Noshchenko ${ }^{3}$, and Piotr Sutkowski ${ }^{2,3}$<br>${ }^{1}$ Institute of Theoretical Physics, University of Wroctaw, PL-50204 Wroctaw, Poland<br>${ }^{2}$ Walter Burke Institute for Theoretical Physics, California Institute of Technology, Pasadena, CA 91125, USA<br>${ }^{3}$ Faculty of Physics, University of Warsaw, ul. Pasteura 5, 02-093 Warsaw, Poland<br>E-mail: jakub.jankowski@uwr.edu.pl, piotrek@caltech.edu, helder.larraguivel@fuw.edu.pl, dmitry.noshchenko@fuw.edu.pl, psulkows@fuw.edu.pl

AbStract: The knots-quivers correspondence states that various characteristics of a knot are encoded in the corresponding quiver and the moduli space of its representations. However, this correspondence is not a bijection: more than one quiver may be assigned to a given knot and encode the same information. In this work we study this phenomenon systematically and show that it is generic rather than exceptional. First, we find conditions that characterize equivalent quivers. Then we show that equivalent quivers arise in families that have the structure of permutohedra, and the set of all equivalent quivers for a given knot is parameterized by vertices of a graph made of several permutohedra glued together. These graphs can be also interpreted as webs of dual $3 \mathrm{~d} \mathcal{N}=2$ theories. All these results are intimately related to properties of homological diagrams for knots, as well as to multi-cover skein relations that arise in counting of holomorphic curves with boundaries on Lagrangian branes in Calabi-Yau three-folds.

## Permutohedron

Permutohedron: ( $n-1$ )-dimensional polytope whose vertices represent permutations of $n$ objects, and edges correspond to transpositions of adjacent neighbours.


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## Compactification and dualities

Superstring theory - effective theory in 4-dim follows from compactification of 10-dim string theory on a Calabi-Yau manifold.

# 10-dim $=R^{4} \times$ 


(Beyond) Standard Model theory

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(Liouville, Toda)
3d-3d correspondence:

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\begin{aligned}
& \text { 6-dim }=\mathbf{R}^{\mathbf{3}} \times(\text { 3-manifold }) \\
& \mathrm{v}=2 \text { sUSY gauge theory } \begin{array}{c}
\text { Chern-Simons theory } \\
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## Chern-Simons theory

Chern-Simons gauge theory - 3-dim TQFT [Witten, 1989]:

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Polynomial knot invariants from Wilson loop observables:

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P_{R}(a, q)=\left\langle\operatorname{Tr}_{R} e^{\oint A}\right\rangle=\int \mathcal{D} A\left(\operatorname{Tr}_{R} e^{\oint A}\right) e^{\frac{i k}{4 \pi} S}
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HOMFLY-PT polynomial for $\operatorname{SU}(\mathrm{N})$ gauge group:

$$
q=e^{\frac{2 \pi}{k+N}}, \quad a=q^{N}
$$

Jones polynomial for $\operatorname{SU}(2)$, Alexander polynomial for $a=1$

## Topological strings and open-closed duality

Chern-Simons theory on $\mathbf{S}^{3}$ arises as an effective description of A-model open topological string theory in deformed conifold $T^{*} \mathbf{S}^{3}$, with appropriate boundary conditions (N branes) on S3 (Witten, 1993).

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Chern-Simons


After a geometric transition, in the 't Hooft limit, $T^{*} \mathbf{S}^{3}$ is replaced by the resolved conifold $X$ (with non-trivial $\mathbf{S}^{2}$ ), $N$ branes disappear, and we are left with A-model closed topological string theory (Gopakumar-Vafa, 1998).

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Z^{\text {closed }}=\exp \left(\sum_{g=0}^{\infty} g_{s}^{2 g-2} F_{g}(Q)\right)
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Knots arise once we introduce an extra lagrangian brane, which intersects $\mathbf{S}^{3}$, along a knot $K$. This brane survives the geometric transition.

## M-theory, knots and BPS states

Embed the above system in M-theory. Chern-Simons theory on $\mathbf{S}^{3}$ engineered by $N$ M5-branes in deformed conifold $T^{*} \mathbf{S}^{3}$. A knot $K$ engineered by extra M5-branes on lagrangian $L_{K}$. What is effective SUSY theory in 3 spacetime dimensions?

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\text { space-time : } & \mathbb{R} \times T^{*} \mathbf{S}^{3} \times M_{4} \\
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Knot invariants of $K$, computed by Chern-Simons theory on the initial $\mathbf{S}^{3}$, are encoded in (conjecturally) integral BPS invariants (Labastida-Marino-Ooguri-Vafa, 2000) in the effective SUSY theory on ( $\mathbf{R} \times \mathrm{D}$ ).

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## Ooguri-Vafa (LMOV) invariants

Brane amplitude as generating function of colored polynomials:

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\sum_{R} P_{R}(a, q) \operatorname{Tr}_{R} V=\exp \left(\sum_{n=1}^{\infty} \sum_{R} \frac{1}{n} f_{R}\left(a^{n}, q^{n}\right) \operatorname{Tr}_{R} V^{n}\right)
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with $f_{R}$ enumerating bound states of D2-D4 branes:

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f_{R}(a, q)=\sum_{i, j} N_{R, i, j} \frac{a^{i} q^{j}}{q-q^{-1}}, \quad N_{R, i, j} \in \mathbb{Z}
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BPS integralities in terms of HOMFLY-PT polynomials:

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f_{S^{3}}(a, q)=P_{S^{3}}(a, q)-P_{\square}(a, q) P_{S^{2}}(a, q)+\frac{1}{3} P_{\square}(a, q)^{3}-\frac{1}{3} P_{\square}\left(a^{3}, q^{3}\right)
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For symmetric representations: $P(x)=\prod\left(1-x^{r} a^{i} q^{j+2 k+1}\right)^{N_{r, i, j}}$ $r \geq 1 ; i, j ; k \geq 0$

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Kucharski-Reineke-Stosic-PS (arXiv: 1707.02991, 1707.04017)
BPS states enumerated by LMOV invariants are bound states of certain "elementary" states, whose interactions are encoded in a quiver diagram. Nodes of a quiver correspond to those "elementary" states, and arrows to interactions.

Calabi-Yau description


Spacetime description

## Quiver representation theory



Consider moduli space of maps $\mathbb{C}^{d_{i}} \rightarrow \mathbb{C}^{d_{j}}$. It is characterized by motivic Donaldson-Thomas invariants: $\Omega_{d_{1}, \ldots, d_{m} ; j} \in \mathbb{N}$

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& =\prod_{\left(d_{1}, \ldots, d_{m}\right) \neq 0} \prod_{j \in \mathbb{Z}} \prod_{k \geq 0}\left(1-\left(x_{1}^{d_{1}} \cdots x_{m}^{d_{m}}\right) q^{j+2 k+1}\right)^{(-1)^{j+1} \Omega_{d_{1}, \ldots, d_{m} ; j}}
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## Knots-quivers correspondence

With appropriate identification of variables, generating function of colored HOMFLY-PT polynomials can be written in the form of motivic generating function, for some particular symmetric matrix $C$

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P(x)=\sum_{r=0}^{\infty} \bar{P}_{r}(a, q) x^{r}=\sum_{d_{1}, \ldots, d_{m} \geq 0} q^{\sum_{i, j} C_{i, j} d_{i} d_{j}} x^{d_{1}+\ldots+d_{m}} \frac{\prod_{i=1}^{m} q^{l_{i} d_{i}} a^{a_{i} d_{i}}(-1)^{t_{i} d_{i}}}{\prod_{i=1}^{m}\left(q^{2} ; q^{2}\right) d_{i}} \\
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Note: infinite number of colored polynomials / LMOV invariants encoded in a finite number of parameters of a matrix $C$.

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| Algebra of BPS states | Cohom. Hall Algebra |

## Quivers and HOMFLY-PT homology

Recall - colored HOMFLY-PT polynomials arise as Euler characteristics of coloured HOMFLY-PT homologies:

$$
P_{r}(a, q)=P_{r}(a, q,-1)=\sum_{i, j, k} a^{i} q^{j}(-1)^{k} \operatorname{dim} \mathcal{H}_{i j k}^{S^{r}}(K) .
$$

Some homological information is encoded in superpolynomials:

$$
P_{r}(a, q, t)=\sum_{i, j, k} a^{i} q^{j} t^{k} \operatorname{dim} \mathcal{H}_{i j k}^{S^{r}}(K) \equiv \sum_{i \in \mathscr{G}_{r}(K)} a^{a_{i}^{(r)}} q^{q_{i}^{(r)}} t^{t_{i}^{(r)}}
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Homological diagram ${ }^{0}$ for figure-8 knot


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Homological diagram ${ }^{0}$ for figure-8 knot


$$
x_{i}=x \lambda_{i}, \quad \lambda_{i}=a^{a_{i}} q^{q_{i}-t_{i}}(-t)^{t_{i}}
$$

Relation to quivers:

$$
t_{i}^{(1)} \equiv t_{i}=C_{i, i}
$$

## Examples

Colored polynomial for trefoil: $P_{r}(a, q)=\frac{a^{2 r}}{q^{2 r}} \sum_{k=0}^{r}\left[\begin{array}{l}r \\ k\end{array}\right]^{2 k(r+1)} \prod_{i=1}^{k}\left(1-a^{2} q^{2(i-2)}\right)$.

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Quiver form follows from:

$$
\left[\begin{array}{c}
r \\
k
\end{array}\right]\left(\frac{a^{2}}{q^{2}} ; q^{2}\right)_{k}=\sum_{i=0}^{k} \frac{\left(q^{2} ; q^{2}\right)_{r}\left(-\frac{a^{2}}{q^{2}}\right)^{i} q^{i(i-1)}}{\left(q^{2} ; q^{2}\right)_{r-k}\left(q^{2} ; q^{2}\right)_{i}\left(q^{2} ; q^{2}\right)_{k-i}} .
$$

We find: $\quad C^{T_{2,3}}=\left[\begin{array}{lll}0 & 1 & 1 \\ 1 & 2 & 2 \\ 1 & 2 & 3\end{array}\right]$


## Examples - torus knots

Colored polynomial for trefoil: $P_{r}(a, q)=\frac{a^{2 r}}{q^{2 r}} \sum_{k=0}^{r}\left[\begin{array}{l}r \\ k\end{array}\right] q^{2 k(r+1)} \prod_{i=1}^{k}\left(1-a^{2} q^{2(i-2)}\right)$
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$$

We find: $\quad C^{T_{2,3}}=\left[\begin{array}{lll}0 & 1 & 1 \\ 1 & 2 & 2 \\ 1 & 2 & 3\end{array}\right]$

$(2,5)$ torus knot: $\quad C^{T_{2,5}}=\left[\begin{array}{lllll}0 & 1 & 1 & 3 & 3 \\ 1 & 2 & 2 & 3 & 3 \\ 1 & 2 & 3 & 4 & 4 \\ 3 & 3 & 4 & 4 & 4 \\ 3 & 3 & 4 & 4 & 5\end{array}\right]$

## Examples - torus knots

$$
C^{T_{3,4}}=\left[\begin{array}{lllll}
0 & 1 & 2 & 3 & 5 \\
1 & 2 & 3 & 3 & 5 \\
2 & 3 & 4 & 4 & 5 \\
3 & 3 & 4 & 4 & 5 \\
5 & 5 & 5 & 5 & 6
\end{array}\right]
$$

## Examples -62 knot

$$
C^{6_{2}}=\left[\begin{array}{ccccccccccc}
-2 & -2 & -1 & -1 & -1 & -1 & 0 & -1 & 1 & 1 & 1 \\
-2 & -1 & -1 & 0 & 0 & 0 & 1 & 0 & 1 & 2 & 2 \\
-1 & -1 & 0 & 1 & 0 & 0 & 1 & 0 & 1 & 2 & 2 \\
-1 & 0 & 1 & 0 & 0 & 0 & 1 & 0 & 2 & 1 & 1 \\
-1 & 0 & 0 & 0 & 1 & 1 & 1 & 1 & 2 & 2 & 2 \\
-1 & 0 & 0 & 0 & 1 & 1 & 1 & 1 & 2 & 2 & 2 \\
0 & 1 & 1 & 1 & 1 & 1 & 2 & 1 & 2 & 2 & 2 \\
-1 & 0 & 0 & 0 & 1 & 1 & 1 & 2 & 2 & 3 & 3 \\
1 & 1 & 1 & 2 & 2 & 2 & 2 & 2 & 3 & 3 & 3 \\
1 & 2 & 2 & 1 & 2 & 2 & 2 & 3 & 3 & 3 & 3 \\
1 & 2 & 2 & 1 & 2 & 2 & 2 & 3 & 3 & 3 & 4
\end{array}\right]
$$

## Examples - 63 knot

$$
C^{6_{3}}=\left[\begin{array}{ccccccccccccc}
0 & 0 & 0 & -1 & -1 & 0 & 0 & -1 & -1 & 0 & 0 & -1 & -1 \\
0 & 1 & 0 & -1 & -2 & 1 & 0 & -1 & -2 & 1 & 1 & 0 & -1 \\
0 & 0 & 0 & -1 & -2 & 1 & 0 & 0 & -2 & 1 & 1 & 0 & 0 \\
-1 & -1 & -1 & -2 & -3 & 0 & -1 & -2 & -3 & -1 & 0 & -2 & -2 \\
-1 & -2 & -2 & -3 & -3 & -1 & -1 & -2 & -3 & -1 & -1 & -2 & -2 \\
0 & 1 & 1 & 0 & -1 & 2 & 1 & 0 & -1 & 2 & 1 & 1 & -1 \\
0 & 0 & 0 & -1 & -1 & 1 & 1 & 0 & -1 & 2 & 1 & 1 & 0 \\
-1 & -1 & 0 & -2 & -2 & 0 & 0 & -1 & -2 & 0 & 0 & -1 & -2 \\
-1 & -2 & -2 & -3 & -3 & -1 & -1 & -2 & -2 & 0 & -1 & -1 & -2 \\
0 & 1 & 1 & -1 & -1 & 2 & 2 & 0 & 0 & 3 & 2 & 1 & 0 \\
0 & 1 & 1 & 0 & -1 & 1 & 1 & 0 & -1 & 2 & 2 & 1 & 0 \\
-1 & 0 & 0 & -2 & -2 & 1 & 1 & -1 & -1 & 1 & 1 & 0 & -1 \\
-1 & -1 & 0 & -2 & -2 & -1 & 0 & -2 & -2 & 0 & 0 & -1 & -1
\end{array}\right]
$$

## Quivers determined for:

- all knots up to 6 crossings
- $(2,2 p+1)$ torus knots, for all $p$
- $(2,2 p)$ torus links
- $(3,3 p+1)$ and $(3,3 p+2)$ torus knots
- an infinite family of twist knots
- an infinite family of rational and arborescent knots (Stosic-Wedrich)


## How unique is the correspondence?

| Knot |  |
| :---: | :---: |
| Unknot | $0_{1}$ |
|  | $3_{1}$ |
|  | $5_{1}$ |
|  | $7_{1}$ |
| Torus knots $T_{2,2 p+1}$ | $9_{1}$ |
|  | $11_{1}$ |
|  | $13_{1}$ |
|  | $15_{1}$ |
| $\vdots$ | $(2 p+1)_{1}$ |
|  | $4_{1}$ |
| Twists knots $T K_{2\|p\|+2}$ | $6_{1}$ |
|  | $8_{1}$ |
| Twists knots $T K_{2 p+1}$ | $5_{2}$ |
|  | $7_{2}$ |
| Stand-alone examples | $6_{2}$ |
|  | $6_{3}$ |
|  | $7_{3}$ |

## Equivalent quivers

| 1 |
| :---: |
| 1 |
| 3 |
| 13 |
| 68 |
| 405 |
| 2684 |
| 19557 |
| $\vdots$ |
| $\sim 2 p!$ |
| 2 |
| 141 |
| 36555 |
| 12 |
| 1983 |
| 3534 |
| 142368 |
| 109636 |

## Local equivalence

## arXiv: 2105.11806

Theorem 6. Consider a quiver $Q$ corresponding to the knot $K$ and another symmetric quiver $Q^{\prime}$ such that $Q_{0}^{\prime}=Q_{0}$ and $\lambda_{i}^{\prime}=\lambda_{i} \forall i \in Q_{0}$ ( $\lambda_{i}$ comes from the knots-quivers change of variables). If $Q$ and $Q^{\prime}$ are related by a sequence of disjoint transpositions, each exchanging non-diagonal elements

$$
C_{a b} \leftrightarrow C_{c d}, \quad C_{b a} \leftrightarrow C_{d c},
$$

for some pairwise different $a, b, c, d, \in Q_{0}$, such that

$$
\lambda_{a} \lambda_{b}=\lambda_{c} \lambda_{d}
$$

and

$$
C_{a b}=C_{c d}-1, \quad C_{a i}+C_{b i}=C_{c i}+C_{d i}-\delta_{c i}-\delta_{d i}, \quad \forall i \in Q_{0},
$$

or

$$
C_{c d}=C_{a b}-1, \quad C_{c i}+C_{d i}=C_{a i}+C_{b i}-\delta_{a i}-\delta_{b i}, \quad \forall i \in Q_{0},
$$

then $Q$ and $Q^{\prime}$ are equivalent in the sense of the definition 4.

## Local equivalence

Proof: follows from comparison of quiver generating series $P(x)$ for $Q$ and $Q^{\prime}$. Agreement at the order $x^{2}$ leads to the center of mass condition (i.e. the center of mass for nodes ( $a, b$ ) coincides with the center of mass for nodes ( $c, d)$ ).

$$
\lambda_{a} \lambda_{b}=\lambda_{c} \lambda_{d}
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Agreement at the order $x^{3}$ leads to the conditions of the form:

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$$

The agreement at the order $x^{3}$ asserts the agreement to all orders, as follows from the multi-cover skein relation.
[Ekholm-Kucharski-Longhi, arXiv: 1910.06193]

## Local equivalence

Verifying systematically conditions from the above theorem, we can identify all equivalent quivers associated to a given knot.


$$
\left[\begin{array}{ccccc}
0 & -1 & -1 & 0 & 0 \\
-1 & -2 & -2 & -1 & 0 \\
-1 & -2 & -1 & -1 & 0 \\
0 & -1 & -1 & 1 & 1 \\
0 & 0 & 0 & 1 & 2
\end{array}\right] 0-\left[\begin{array}{ccccc}
0 & -1 & -1 & 0 & 0 \\
-1 & -2 & -2 & -1 & -1 \\
-1 & -2 & -1 & 0 & 0 \\
0 & -1 & 0 & 1 & 1 \\
0 & -1 & 0 & 1 & 2
\end{array}\right]
$$

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$\left[\begin{array}{lllll}0 & 1 & 1 & 3 & 2 \\ 1 & 2 & 3 & 3 & 3 \\ 1 & 3 & 3 & 4 & 4 \\ 3 & 3 & 4 & 4 & 4 \\ 2 & 3 & 4 & 4 & 5\end{array}\right] \quad\left[\begin{array}{lllll}0 & 1 & 1 & 3 & 3 \\ 1 & 2 & 2 & 3 & 3 \\ 1 & 2 & 3 & 4 & 4 \\ 3 & 3 & 4 & 4 & 4 \\ 3 & 3 & 4 & 4 & 5\end{array}\right] \quad\left[\begin{array}{lllll}0 & 1 & 1 & 3 & 3 \\ 1 & 2 & 2 & 3 & 4 \\ 1 & 2 & 3 & 3 & 4 \\ 3 & 3 & 3 & 4 & 4 \\ 3 & 4 & 4 & 4 & 5\end{array}\right]$

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## Global structure

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Instead of analyzing quiver matrices, consider the structure of quiver generating series. We find that in general it takes form

$$
P_{K}(x, a, q, t)=\left.\sum_{\check{d}_{1}, \ldots, \check{d}_{m-n} \geq 0}(-q)^{\sum_{i, j} \check{C}_{i j} \check{d}_{i} \check{d}_{j}} \frac{\check{x}_{1}^{\breve{d}_{1}} \cdots \check{x}_{m-n}^{\check{d}_{m-n}}}{\left(q^{2} ; q^{2}\right)_{\check{d}_{1}} \cdots\left(q^{2} ; q^{2}\right)_{\check{d}_{m-n}}} \Pi_{\check{d}_{1}, \ldots, \check{d}_{n}}\right|_{\check{x}_{i}=x \check{\lambda}_{i}}
$$

where $\check{C}$ is a matrix of a subquiver, and the last piece takes form

$$
\frac{\Pi_{\breve{d}_{1}, \ldots, \check{d}_{n}}}{\left(q^{2} ; q^{2}\right)_{\breve{d}_{1}} \cdots\left(q^{2} ; q^{2}\right)_{\breve{d}_{n}}}=\sum_{\check{d}_{1}=\alpha_{1}+\beta_{1}} \ldots \sum_{\check{d}_{n}=\alpha_{n}+\beta_{n}} \frac{(-q)^{2 \sum_{i<j} \beta_{i} \alpha_{j}+\pi_{2}\left(\alpha_{1}, \ldots, \alpha_{n} ; \beta_{1}, \ldots, \beta_{n}\right)} \kappa^{\beta_{1}+\ldots+\beta_{n}}}{\left(q^{2} ; q^{2}\right)_{\alpha_{1}}\left(q^{2} ; q^{2}\right)_{\beta_{1}} \cdots\left(q^{2} ; q^{2}\right)_{\alpha_{n}}\left(q^{2} ; q^{2}\right)_{\beta_{n}}}
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$$

Recall: a permutation $\sigma$ is determined by a set of its inversions, i.e. a set of all pairs $(\sigma(i), \sigma(j))$ such that $i<j$ and $\sigma(i)>\sigma(j)$. Such a set of inversions is encoded in the term $\sum_{i<j} \beta_{i} \alpha_{j}$

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It follows that various identifications of indices lead to a permutohedron of equivalent quivers!

## Global structure

Such structures arise from the following formula for

$$
\Pi_{\check{d}_{1}, \ldots, \check{d}_{n}}=\left(\xi ; q^{2}\right)_{\check{d}_{1}+\ldots+\check{d}_{n}}
$$

(or its generalizations):

$$
\begin{aligned}
& \frac{\left(\xi ; q^{2}\right)_{\check{d}_{1}+\ldots+\check{d}_{n}}^{\left(q^{2} ; q^{2}\right)_{\breve{d}_{1}} \cdots\left(q^{2} ; q^{2}\right)_{\check{d}_{n}}}=}{} \sum_{\alpha_{1}+\beta_{1}=\check{d}_{1}} \ldots \sum_{\alpha_{n}+\beta_{n}=\check{d}_{n}}(-q)^{\beta_{1}^{2}+\ldots+\beta_{n}^{2}+2 \sum_{i=1}^{n-1} \beta_{i+1}\left(\check{d}_{1}+\ldots+\check{d}_{i}\right)} \times \\
& \times \frac{\left(\xi q^{-1}\right)^{\beta_{1}+\cdots+\beta_{n}}}{\left(q^{2} ; q^{2}\right)_{\alpha_{1}}\left(q^{2} ; q^{2}\right)_{\beta_{1}} \cdots\left(q^{2} ; q^{2}\right)_{\alpha_{n}}\left(q^{2} ; q^{2}\right)_{\beta_{n}}},
\end{aligned}
$$

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(or its generalizations):

$$
\begin{aligned}
\frac{\left(\xi ; q^{2}\right)_{\check{d}_{1}+\ldots}+\ldots \check{d}_{n}}{\left(q^{2} ; q^{2}\right)_{\check{d}_{1}}^{\cdots} \cdots\left(q^{2} ; q^{2}\right)_{\check{d}_{n}}}= & \sum_{\alpha_{1}+\beta_{1}=\check{d}_{1}} \cdots \sum_{\alpha_{n}+\beta_{n}=\check{d}_{n}}(-q)^{\beta_{1}^{2}+\ldots+\beta_{n}^{2}+2 \sum_{i=1}^{n-1} \beta_{i+1}\left(\tilde{d}_{1}+\ldots+\check{d}_{i}\right)} \times \\
& \times \frac{\left(q^{2} ; q^{2}\right)_{\alpha_{1}}\left(q^{2} ; q^{2}\right)_{\beta_{1}} \cdots\left(q^{2} ; q^{2}\right)_{\alpha_{n}}\left(q^{2} ; q^{2}\right)_{\beta_{n}}}{\beta_{1}+\cdots+\beta_{n}},
\end{aligned}
$$

Such form of quiver generating functions follows from constraints in the local equivalence theorem.

## Global structure

We refer to the subquiver mentioned above as a "prequiver".
The full quiver is determined from a sub quiver by permutation and a pair of integers ( $k, l$ ), in the operation called "splitting".


## Global structure

In general, there are several equivalent formulas for a given HOMFLY-PT generating function, and each of them gives rise to one permutohedron. Altogether we obtain a large permutohedron graph.

## Global structure - examples

In this case permutohedron graph is made of 3 permutohedra $\Pi_{3}$.


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## Global structure - examples

For $(2,2 p+1)$ torus knots, permutohedron graph is made of two chains of larger and larger permutohedra.


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## 91



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## Global structure - examples

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## Summary

- Knots-quivers correspondence, motivated by string theory, relates knot theory and quiver representation theory
- It turns out that many quivers may be associated to a given knot
- They are parametrized by vertices of a permutohedron graph
- This indicates some interesting structure of the underlying HOMFLY-PT homology, and of the corresponding LMOV (motivic DT) invariants



## FNP

## Summary

Future directions and related developments:

- identify permutohedra for rational and arborescent knots (following Stosic-Wedrich, arXiv: 1711.03333, 2004.10837)
- develop open topological string interpretation (following Ekholm-Kucharski-Longhi, arXiv: 1811.03110, 1910.06193)
- conduct analogous analysis for other underlying toric Calabi-Yau manifolds (following Kimura-Panfil-Sugimoto-PS , arXiv: 1811.03556, 2011.06783)



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