

# Bilinear discretization of quadratic vector fields: integrability and geometry

Yuri B. Suris

(Technische Universität Berlin)

HU Seminar “Algebra, Geometry and Physics”  
Berlin, 27.04.2021



Discretization in  
Geometry and Dynamics  
SFB Transregio 109

# Part 1. Generalities

# The problem of integrable discretization. Hamiltonian approach (Birkhäuser, 2003)

Consider a completely integrable flow

$$\dot{x} = f(x) = \{H, x\} \quad (1)$$

with a Hamilton function  $H$  on a Poisson manifold  $\mathcal{P}$  with a Poisson bracket  $\{\cdot, \cdot\}$ . Thus, flow (1) possesses sufficiently many functionally independent integrals  $I_k(x)$  in involution.

The *problem of integrable discretization*: find a family of diffeomorphisms  $\mathcal{P} \rightarrow \mathcal{P}$ ,

$$\tilde{x} = \Phi(x; \epsilon), \quad (2)$$

depending smoothly on a small parameter  $\epsilon > 0$ , with the following properties:

1. The maps (2) *approximate* the flow (1):

$$\Phi(x; \epsilon) = x + \epsilon f(x) + O(\epsilon^2).$$

2. The maps (2) are *Poisson* w. r. t. the bracket  $\{\cdot, \cdot\}$  or some its deformation  $\{\cdot, \cdot\}_\epsilon = \{\cdot, \cdot\} + O(\epsilon)$ .
3. The maps (2) are *integrable*, i.e. possess the necessary number of independent integrals in involution,  
 $I_k(x; \epsilon) = I_k(x) + O(\epsilon)$ .

While integrable lattice systems (like Toda or Volterra lattices) can be discretized in a systematic way (based, e.g., on the  $r$ -matrix structure), there is no systematic way to obtain *decent* integrable discretizations for integrable systems of classical mechanics.

- ▶ R.Hirota, K.Kimura. *Discretization of the Euler top*. J. Phys. Soc. Japan **69** (2000) 627–630,
- ▶ K.Kimura, R.Hirota. *Discretization of the Lagrange top*. J. Phys. Soc. Japan **69** (2000) 3193–3199.

Reasons for this omission: discretization of the Euler top seemed to be an isolated curiosity; discretization of the Lagrange top seemed to be completely incomprehensible, if not even wrong.

Renewed interest stimulated by a talk by T. Ratiu at the Oberwolfach Workshop “Geometric Integration”, March 2006, who claimed that HK-type discretizations for the Clebsch system and for the Kovalevsky top are also integrable.

# Hirota-Kimura or Kahan?

- ▶ W. Kahan. *Unconventional numerical methods for trajectory calculations* (Unpublished lecture notes, 1993).

$$\dot{x} = Q(x) + Bx + c \rightsquigarrow (\tilde{x} - x)/\epsilon = Q(x, \tilde{x}) + B(x + \tilde{x})/2 + c,$$

where  $B \in \mathbb{R}^{n \times n}$ ,  $c \in \mathbb{R}^n$ , each component of  $Q : \mathbb{R}^n \rightarrow \mathbb{R}^n$  is a *quadratic* form, and  $Q(x, \tilde{x}) = (Q(x + \tilde{x}) - Q(x) - Q(\tilde{x}))/2$  is the corresponding symmetric *bilinear* function. Thus,

$$\dot{x}_k \rightsquigarrow (\tilde{x}_k - x_k)/\epsilon, \quad x_k^2 \rightsquigarrow x_k \tilde{x}_k, \quad x_j x_k \rightsquigarrow (x_j \tilde{x}_k + \tilde{x}_j x_k)/2.$$

Linear w.r.t.  $\tilde{x}$ , therefore defines a *rational* map  $\tilde{x} = \Phi_f(x, \epsilon)$ .

Obvious symmetry:  $x \leftrightarrow \tilde{x}$ ,  $\epsilon \mapsto -\epsilon$ , therefore  $\Phi_f$  *reversible*:

$$\Phi_f^{-1}(x, \epsilon) = \Phi_f(x, -\epsilon).$$

In particular,  $\Phi_f$  is *birational*, and  $\deg \Phi_f = \deg \Phi_f^{-1} = n$ .

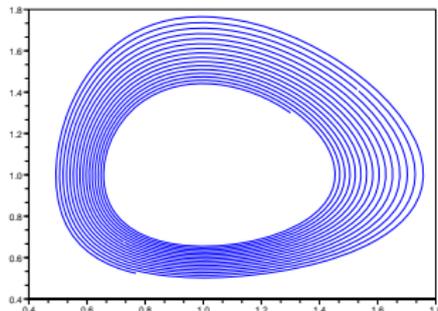
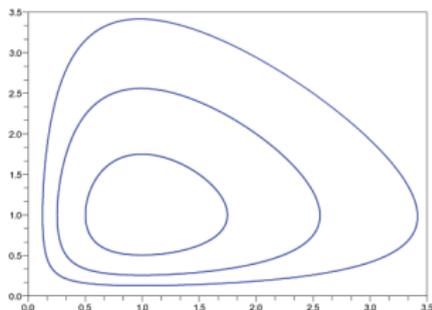
# Illustration: Lotka-Volterra system

Kahan's discretization for the Lotka-Volterra system:

$$\begin{cases} \dot{x} = x(1 - y), \\ \dot{y} = y(x - 1), \end{cases} \rightsquigarrow \begin{cases} \tilde{x} - x = \epsilon(\tilde{x} + x) - \epsilon(\tilde{x}y + x\tilde{y}), \\ \tilde{y} - y = \epsilon(\tilde{x}y + x\tilde{y}) - \epsilon(\tilde{y} + y). \end{cases}$$

Explicitly:

$$\begin{cases} \tilde{x} = x \frac{(1 + \epsilon)^2 - \epsilon(1 + \epsilon)x - \epsilon(1 - \epsilon)y}{1 - \epsilon^2 - \epsilon(1 - \epsilon)x + \epsilon(1 + \epsilon)y}, \\ \tilde{y} = y \frac{(1 - \epsilon)^2 + \epsilon(1 + \epsilon)x + \epsilon(1 - \epsilon)y}{1 - \epsilon^2 - \epsilon(1 - \epsilon)x + \epsilon(1 + \epsilon)y}. \end{cases}$$



Left: three orbits of Kahan's discretization with  $\epsilon = 0.1$ ,  
 right: one orbit of the explicit Euler with  $\epsilon = 0.01$ .

- ▶ J.M. Sanz-Serna. *An unconventional symplectic integrator of W.Kahan*. Applied Numer. Math. 1994, **16**, 245–250.

A sort of an explanation of a non-spiralling behavior: Kahan's discretization is symplectic w.r.t.  $dx \wedge dy/(xy)$ .

# Hirota-Kimura's discrete time Euler top

$$\begin{cases} \dot{x}_1 = \alpha_1 x_2 x_3, \\ \dot{x}_2 = \alpha_2 x_3 x_1, \\ \dot{x}_3 = \alpha_3 x_1 x_2, \end{cases} \rightsquigarrow \begin{cases} \tilde{x}_1 - x_1 = \epsilon \alpha_1 (\tilde{x}_2 x_3 + x_2 \tilde{x}_3), \\ \tilde{x}_2 - x_2 = \epsilon \alpha_2 (\tilde{x}_3 x_1 + x_3 \tilde{x}_1), \\ \tilde{x}_3 - x_3 = \epsilon \alpha_3 (\tilde{x}_1 x_2 + x_1 \tilde{x}_2). \end{cases}$$

Features:

- ▶ Equations are linear w.r.t.  $\tilde{x} = (\tilde{x}_1, \tilde{x}_2, \tilde{x}_3)^T$ :

$$A(x, \epsilon) \tilde{x} = x, \quad A(x, \epsilon) = \begin{pmatrix} 1 & -\epsilon \alpha_1 x_3 & -\epsilon \alpha_1 x_2 \\ -\epsilon \alpha_2 x_3 & 1 & -\epsilon \alpha_2 x_1 \\ -\epsilon \alpha_3 x_2 & -\epsilon \alpha_3 x_1 & 1 \end{pmatrix},$$

result in a rational map, which is *reversible* (therefore birational):

$$\tilde{x} = \Phi(x, \epsilon) = A^{-1}(x, \epsilon)x, \quad \Phi^{-1}(x, \epsilon) = \Phi(x, -\epsilon).$$

► Explicit formulas:

$$\left\{ \begin{array}{l} \tilde{x}_1 = \frac{x_1 + 2\epsilon\alpha_1x_2x_3 + \epsilon^2x_1(-\alpha_2\alpha_3x_1^2 + \alpha_3\alpha_1x_2^2 + \alpha_1\alpha_2x_3^2)}{\Delta(x, \epsilon)}, \\ \tilde{x}_2 = \frac{x_2 + 2\epsilon\alpha_2x_3x_1 + \epsilon^2x_2(\alpha_2\alpha_3x_1^2 - \alpha_3\alpha_1x_2^2 + \alpha_1\alpha_2x_3^2)}{\Delta(x, \epsilon)}, \\ \tilde{x}_3 = \frac{x_3 + 2\epsilon\alpha_3x_1x_2 + \epsilon^2x_3(\alpha_2\alpha_3x_1^2 + \alpha_3\alpha_1x_2^2 - \alpha_1\alpha_2x_3^2)}{\Delta(x, \epsilon)}, \end{array} \right.$$

where  $\Delta(x, \epsilon) = \det A(x, \epsilon)$

$$= 1 - \epsilon^2(\alpha_2\alpha_3x_1^2 + \alpha_3\alpha_1x_2^2 + \alpha_1\alpha_2x_3^2) - 2\epsilon^3\alpha_1\alpha_2\alpha_3x_1x_2x_3.$$

- ▶ Two independent integrals:

$$I_1(x, \epsilon) = \frac{1 - \epsilon^2 \alpha_2 \alpha_3 x_1^2}{1 - \epsilon^2 \alpha_3 \alpha_1 x_2^2}, \quad I_2(x, \epsilon) = \frac{1 - \epsilon^2 \alpha_3 \alpha_1 x_2^2}{1 - \epsilon^2 \alpha_1 \alpha_2 x_3^2}.$$

- ▶ Invariant volume form:

$$\omega = \frac{dx_1 \wedge dx_2 \wedge dx_3}{\phi(x)}, \quad \phi(x) = 1 - \epsilon^2 \alpha_j \alpha_k x_k^2$$

and bi-Hamiltonian structure found in:

- ▶ M. Petrera, Yu. S. *On the Hamiltonian structure of the Hirota-Kimura discretization of the Euler top.* Math. Nachr., 2010, **283**, 1654–1663.

# Hirota-Kimura's discrete time Lagrange top

Equations of motion of the Lagrange top:

$$\dot{m}_1 = (\alpha - 1)m_2m_3 + \gamma p_2,$$

$$\dot{m}_2 = (1 - \alpha)m_1m_3 - \gamma p_1,$$

$$\dot{m}_3 = 0,$$

$$\dot{p}_1 = \alpha p_2 m_3 - p_3 m_2,$$

$$\dot{p}_2 = p_3 m_1 - \alpha p_1 m_3,$$

$$\dot{p}_3 = p_1 m_2 - p_2 m_1.$$

It is Hamiltonian w.r.t. Lie-Poisson bracket of  $e(3)$ , has four functionally independent integrals in involution: two Casimir functions,

$$C_1 = p_1^2 + p_2^2 + p_3^2, \quad C_2 = m_1 p_1 + m_2 p_2 + m_3 p_3,$$

the Hamilton function, and the (trivial) “fourth integral”,

$$H_1 = \frac{1}{2}(m_1^2 + m_2^2 + \alpha m_3^2) + \gamma p_3, \quad H_2 = m_3.$$

Discretization:

$$\tilde{m}_1 - m_1 = \epsilon(\alpha - 1)(\tilde{m}_2 m_3 + m_2 \tilde{m}_3) + \epsilon\gamma(\rho_2 + \tilde{\rho}_2),$$

$$\tilde{m}_2 - m_2 = \epsilon(1 - \alpha)(\tilde{m}_1 m_3 + m_1 \tilde{m}_3) - \epsilon\gamma(\rho_1 + \tilde{\rho}_1),$$

$$\tilde{m}_3 - m_3 = 0,$$

$$\tilde{\rho}_1 - \rho_1 = \epsilon\alpha(\rho_2 \tilde{m}_3 + \tilde{\rho}_2 m_3) - \epsilon(\rho_3 \tilde{m}_2 + \tilde{\rho}_3 m_2),$$

$$\tilde{\rho}_2 - \rho_2 = \epsilon(\rho_3 \tilde{m}_1 + \tilde{\rho}_3 m_1) - \epsilon\alpha(\rho_1 \tilde{m}_3 + \tilde{\rho}_1 m_3),$$

$$\tilde{\rho}_3 - \rho_3 = \epsilon(\rho_1 \tilde{m}_2 + \tilde{\rho}_1 m_2 - \rho_2 \tilde{m}_1 - \tilde{\rho}_2 m_1).$$

As usual, get an explicit birational map  $(\tilde{m}, \tilde{\rho}) = \Phi(m, \rho, \epsilon)$ .

Trivial conserved quantity  $m_3 = \text{const}$ . Very difficult to find any further conserved quantity!

# Hirota-Kimura's method for finding integrals

**Incredible claim by HK:** for any initial point, there exist  $A, B, C \in \mathbb{R}$  such that

$$A(m_1^2 + m_2^2) + Bp_3^2 + Cp_3 = 1$$

along the orbit  $\Phi^i(p, m, \epsilon)$ ,  $i \in \mathbb{Z}$ .

How one could check this? Solve the system for the unknowns  $A, B, C$  for  $i = -1, 0, 1$ :

$$\begin{cases} A(\tilde{m}_1^2 + \tilde{m}_2^2) + B\tilde{p}_3^2 + C\tilde{p}_3 = 1, \\ A(m_1^2 + m_2^2) + Bp_3^2 + Cp_3 = 1, \\ A(\underline{m}_1^2 + \underline{m}_2^2) + B\underline{p}_3^2 + C\underline{p}_3 = 1 \end{cases}$$

with  $(\tilde{m}, \tilde{p}) = \Phi(m, p, \epsilon)$  and  $(\underline{m}, \underline{p}) = \Phi^{-1}(m, p, \epsilon)$ . Then check that  $A, B, C = A, B, C(m, p, \epsilon)$  are conserved quantities.

Why should this work???

**Definition.** For a given bijective map  $\Phi : \mathbb{R}^n \rightarrow \mathbb{R}^n$ , a set of functions  $\Psi = (\psi_1, \dots, \psi_\ell)$ , linearly independent over  $\mathbb{R}$ , is called a **HK-set**, if for every  $x_0 \in \mathbb{R}^n$  there exists a vector  $c = (c_1, \dots, c_\ell) \neq 0$ ,  $c = c(x_0)$ , such that

$$c_1 \psi_1(\Phi^j(x_0)) + \dots + c_\ell \psi_\ell(\Phi^j(x_0)) = 0 \quad \forall j \in \mathbb{Z}.$$

For a given  $x_0 \in \mathbb{R}^n$ , the set  $K_\Psi(x_0)$  of all vectors  $c(x_0) \in \mathbb{R}^\ell$  with this property is called the null-space of the HK-set  $\Psi$  (at the point  $x_0$ ). This is clearly a vector space.

**Dynamical consequence.** Existence of a HK-set  $\Psi$  with  $\dim K_\Psi(x_0) = d$  confines orbits of  $\Phi$  to  $(n - d)$ -dimensional invariant sets (similarly to the presence of  $d$  integrals).

**Proposition.** *If  $\Psi$  is a HK-set for a map  $\Phi$  with a  $d$ -dimensional null space then  $K_\Psi(\Phi(x_0)) = K_\Psi(x_0)$ , a  $Gr(d, \ell)$ -valued integral.*

Its Plücker coordinates are scalar integrals.

The most useful particular case:

**Corollary.** *Let  $\Psi$  be a HK-set for  $\Phi$  with  $\dim K_\Psi(x_0) = 1$  for all  $x_0 \in \mathbb{R}^n$ . Let  $K_\Psi(x_0) = [c_1(x_0) : \dots : c_\ell(x_0)] \in \mathbb{RP}^{\ell-1}$ . Then the functions  $c_j/c_k$  are integrals of motion for  $\Phi$ .*

The number of functionally independent integrals among them varies in examples (sometimes just = 1 and sometimes  $> 1$ ).

Results by Hirota and Kimura in the Lagrange top case:

**Theorem.** *The three sets of functions,*

$$\Psi_1 = (m_1^2 + m_2^2, p_3^2, p_3, 1),$$

$$\Psi_2 = (m_1 p_1 + m_2 p_2, p_3^2, p_3, 1),$$

$$\Psi_3 = (p_1^2 + p_2^2, p_3^2, p_3, 1),$$

*are HK-sets for the discrete time Lagrange top with one-dimensional null-spaces, each producing one independent integral.*

It follows that any orbit lies on a two-dimensional surface in  $\mathbb{R}^6$  which is intersection of three quadrics and a hyperplane  $m_3 = \text{const}$ .

# A simple integral (unnoticed by Hirota and Kimura)

**Theorem.** *The functions*

$$\Gamma = (\tilde{m}_1 p_1 - m_1 \tilde{p}_1, \tilde{m}_2 p_2 - m_2 \tilde{p}_2, \tilde{m}_3 p_3 - m_3 \tilde{p}_3)$$

*build a HK-set for the discrete time Lagrange top with one-dimensional null-space  $K_\Gamma(x) = [1 : 1 : J]$ ,*

$$J = \frac{(2\alpha - 1) + \epsilon^2(\alpha - 1)(m_1^2 + m_2^2) + \epsilon^2\gamma(m_1 p_1 + m_2 p_2)/m_3}{1 + \epsilon^2\alpha(1 - \alpha)m_3^2 - \epsilon^2\gamma p_3}.$$

**Theorem.** *The discrete time Lagrange top possesses an invariant volume form:*

$$\Phi^*\omega = \omega, \quad \omega = \frac{dm_1 \wedge dm_2 \wedge dm_3 \wedge dp_1 \wedge dp_2 \wedge dp_3}{\Delta(m, p)},$$

where

$$\Delta = 1 + \epsilon^2 \Delta^{(2)} + \epsilon^4 \Delta^{(4)} + \epsilon^6 \Delta^{(6)},$$

and  $\Delta^{(q)}$  are polynomials of degree  $q$  in  $(m, p)$ .

# Further examples of integrable HK-discretizations

Overview given in

- ▶ M. Petrera, A. Pfadler, Yu. S. *On integrability of Hirota-Kimura type discretizations*. Regular Chaotic Dyn., 2011, **16**, 245–289.

1. Reduced Nahm equations.
2. Three-wave interaction system.
3. Periodic Volterra chain of period  $N = 3, 4$ :

$$\dot{x}_k = x_k(x_{k+1} - x_{k-1}), \quad k \in \mathbb{Z}/N\mathbb{Z}$$

4. Dressing chain with  $N = 3$ :

$$\dot{x}_k + \dot{x}_{k+1} = x_{k+1}^2 - x_k^2 + \alpha_{k+1} - \alpha_k, \quad k \in \mathbb{Z}/N\mathbb{Z}, \quad N \text{ odd.}$$

5. System of two interacting Euler tops.
6. Kirchhof and Clebsch cases of rigid body in an ideal fluid.

# Clebsch system

Clebsch case of the motion of a rigid body in an ideal fluid:

$$\dot{m}_1 = (\omega_3 - \omega_2)p_2p_3,$$

$$\dot{m}_2 = (\omega_1 - \omega_3)p_3p_1,$$

$$\dot{m}_3 = (\omega_2 - \omega_1)p_1p_2,$$

$$\dot{p}_1 = m_3p_2 - m_2p_3,$$

$$\dot{p}_2 = m_1p_3 - m_3p_1,$$

$$\dot{p}_3 = m_2p_1 - m_1p_2.$$

It is Hamiltonian w.r.t. Lie-Poisson bracket of  $e(3)$ , has four functionally independent integrals in involution:

$$I_i = p_i^2 + \frac{m_j^2}{\omega_k - \omega_j} + \frac{m_k^2}{\omega_j - \omega_i}, \quad (i, j, k) = c.p.(1, 2, 3),$$

and  $H_4 = m_1p_1 + m_2p_2 + m_3p_3$ .

A Hirota-Kimura (or Kahan) style discretization:

$$\tilde{m}_1 - m_1 = \epsilon(\omega_3 - \omega_2)(\tilde{p}_2 p_3 + p_2 \tilde{p}_3),$$

$$\tilde{m}_2 - m_2 = \epsilon(\omega_1 - \omega_3)(\tilde{p}_3 p_1 + p_3 \tilde{p}_1),$$

$$\tilde{m}_3 - m_3 = \epsilon(\omega_2 - \omega_1)(\tilde{p}_1 p_2 + p_1 \tilde{p}_2),$$

$$\tilde{p}_1 - p_1 = \epsilon(\tilde{m}_3 p_2 + m_3 \tilde{p}_2) - \epsilon(\tilde{m}_2 p_3 + m_2 \tilde{p}_3),$$

$$\tilde{p}_2 - p_2 = \epsilon(\tilde{m}_1 p_3 + m_1 \tilde{p}_3) - \epsilon(\tilde{m}_3 p_1 + m_3 \tilde{p}_1),$$

$$\tilde{p}_3 - p_3 = \epsilon(\tilde{m}_2 p_1 + m_2 \tilde{p}_1) - \epsilon(\tilde{m}_1 p_2 + m_1 \tilde{p}_2).$$

A birational map of  $\mathbb{R}^6$  of degree 6:

$$\begin{pmatrix} \tilde{m} \\ \tilde{p} \end{pmatrix} = \Phi(m, p, \epsilon) = M^{-1}(m, p, \epsilon) \begin{pmatrix} m \\ p \end{pmatrix},$$

$$M(m, p, \epsilon) = \begin{pmatrix} 1 & 0 & 0 & 0 & \epsilon\omega_{23}p_3 & \epsilon\omega_{23}p_2 \\ 0 & 1 & 0 & \epsilon\omega_{31}p_3 & 0 & \epsilon\omega_{31}p_1 \\ 0 & 0 & 1 & \epsilon\omega_{12}p_2 & \epsilon\omega_{12}p_1 & 0 \\ 0 & \epsilon p_3 & -\epsilon p_2 & 1 & -\epsilon m_3 & \epsilon m_2 \\ -\epsilon p_3 & 0 & \epsilon p_1 & \epsilon m_3 & 1 & -\epsilon m_1 \\ \epsilon p_2 & -\epsilon p_1 & 0 & -\epsilon m_2 & \epsilon m_1 & 1 \end{pmatrix},$$

with  $\omega_{ij} = \omega_j - \omega_i$ . The usual reversibility:

$$\Phi^{-1}(m, p, \epsilon) = \Phi(m, p, -\epsilon).$$

Based on:

- ▶ M. Petrera, A. Pfadler, Yu. S. *On integrability of Hirota-Kimura type discretizations. Experimental study of the discrete Clebsch system.* *Experimental Math.*, 2009, **18**, 223–247.
- ▶ M. Petrera, Yu. S. *New results on integrability of the Kahan-Hirota-Kimura discretizations.* - In: *Nonlinear Systems and Their Remarkable Mathematical Structures*, CRC Press, 2018, 94–120.

**Theorem.** a) *The set of functions*

$$\Psi = (p_1^2, p_2^2, p_3^2, m_1^2, m_2^2, m_3^2, m_1 p_1, m_2 p_2, m_3 p_3, 1)$$

*is a HK-set for  $\Phi$ , with  $\dim K_\Psi(m, p) = 4$ . Thus, any orbit of  $\Phi$  lies on an intersection of four quadrics in  $\mathbb{R}^6$ .*

b) *The following four are HK-sets for  $\Phi$  with one-dimensional null-spaces:*

$$\Psi_0 = (p_1^2, p_2^2, p_3^2, 1),$$

$$\Psi_1 = (p_1^2, p_2^2, p_3^2, m_1^2, m_2^2, m_3^2, m_1 p_1),$$

$$\Psi_2 = (p_1^2, p_2^2, p_3^2, m_1^2, m_2^2, m_3^2, m_2 p_2),$$

$$\Psi_3 = (p_1^2, p_2^2, p_3^2, m_1^2, m_2^2, m_3^2, m_3 p_3).$$

*There holds:  $K_\Psi = K_{\Psi_0} \oplus K_{\Psi_1} \oplus K_{\Psi_2} \oplus K_{\Psi_3}$ .*

# Complexity issues

The claims in part b) refer to solutions of the following systems:

$$(c_1 p_1^2 + c_2 p_2^2 + c_3 p_3^2) \circ \Phi^i = 1,$$

(to be solved for 3 consecutive values of  $i$ , e.g.,  $i = -1, 0, 1$ ),  
and

$$(\alpha_1 p_1^2 + \alpha_2 p_2^2 + \alpha_3 p_3^2 + \alpha_4 m_1^2 + \alpha_5 m_2^2 + \alpha_6 m_3^2) \circ \Phi^i = m_1 p_1 \circ \Phi^i,$$

etc. (to be solved for 6 consecutive values of  $i$ , e.g.,  $i \in [-2, 3]$ ).

This is a serious challenge for symbolic computations (for  $\Phi^3$  we are dealing with polynomials of degree 216 in 6 variables which is prohibitively complex). Various tricks invented to reduce the range of  $i$ .

# Integral for non-integrable Kahan discretizations

- ▶ E. Celledoni, R.I. McLachlan, B. Owren, G.R.W. Quispel. *Geometric properties of Kahan's method*. J. Phys. A, 2013, **46**, 025201.

**Theorem.** Let  $f(x) = J\nabla H(x)$ , with  $J \in so(n)$ , Hamilton function  $H : \mathbb{R}^n \rightarrow \mathbb{R}$  of  $\deg = 3$ . Then  $\Phi_f(x, \epsilon)$  admits a rational integral:

$$\tilde{H}(x, \epsilon) = H(x) + \frac{\epsilon}{3} (\nabla H(x))^T \left( I - \frac{\epsilon}{2} f'(x) \right)^{-1} f(x),$$

and an invariant volume form

$$\frac{dx_1 \wedge \dots \wedge dx_n}{\det \left( I - \frac{\epsilon}{2} f'(x) \right)}.$$

Degree of denominator  $\det(I - \frac{\epsilon}{2} f'(x))$  is  $n$ , degree of numerator of  $\tilde{H}(x, \epsilon)$  is  $n + 1$ .

. Part 2. Integrability of planar quadratic birational maps

# Why planar?

- ▶ Planar algebraic geometry is much simpler.
- ▶ Structure of the group of birational maps of  $\mathbb{P}^n$  is unknown for  $n \geq 3$ . For  $n = 2$ , generated by quadratic maps (M. Noether theorem).
- ▶ For  $n \geq 3$ , many new phenomena. For instance, there does not hold necessarily that  $\deg \Phi^{-1} = \deg \Phi$ . (Kahan maps have this property and thus are very special!)

# Planar birational maps

- ▶ Consider a birational map

$$\phi: \mathbb{CP}^2 \rightarrow \mathbb{CP}^2, \quad [x : y : z] \mapsto [X : Y : Z],$$

$X, Y, Z$  homogeneous polynomials of  $\deg = d$  without a non-trivial (polynomial) common factor.

- ▶ *Indeterminacy set* (finitely many points, are blown up by  $\phi$ ):

$$\mathcal{I}(\phi) = \{X = Y = Z = 0\}.$$

- ▶ *Critical set* ( $\dim = 1$ , is blown down by  $\phi$ ):

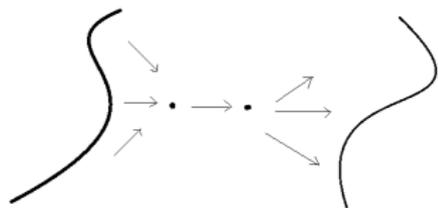
$$\mathcal{C}(\phi) = \{\det \partial(X, Y, Z) / \partial(x, y, z) = 0\}.$$

# Degree lowering and singularity confinement

A component  $V \subset \mathcal{C}(\phi)$  is a *degree lowering curve*, if for some  $n \in \mathbb{N}$  we have  $\phi^n(V) \in \mathcal{I}(\phi)$ . A *singularity confinement pattern* is a sequence

$$\mathcal{C}(\phi) \supset V \rightarrow \phi(V) \rightarrow \cdots \rightarrow \phi^n(V) \rightarrow \phi^{n+1}(V) \subset \mathcal{C}(\phi^{-1}).$$

A presence of such a curve is necessary and sufficient for  $\deg(\phi^n) < (\deg \phi)^n$ .



**Definition.** *Dynamical degree* and *algebraic entropy* of  $\phi$  are

$$\lambda_1(\phi) = \lim_{n \rightarrow \infty} (\deg(\phi^n))^{1/n} \leq d \quad \text{and} \quad h(\phi) = \log(\lambda_1(\phi)) \leq \log(d).$$

Inequalities strict iff there exist degree lowering curves.

How drastic can be the degree drop of iterations  $\phi^n$ ?

**Definition.** A birational map  $\phi$  is *integrable* if  $h(\phi) = 0$ .

## Birational quadratic maps of $\mathbb{P}^2$

A generic birational map  $\phi : \mathbb{P}^2 \dashrightarrow \mathbb{P}^2$  of  $\deg = 2$  can be represented as  $\phi = A_1 \circ \sigma \circ A_2$ , where  $A_1, A_2 \in \text{Aut}(\mathbb{P}^2)$ , and

$$\sigma : [x : y : z] \rightarrow [yz : xz : xy].$$

The dimension of this orbit is 14.

A generic map from this set, not an involution, can be described by a pair of bilinear (Kahan type) relations:

$$\begin{aligned}\tilde{x} - x &= a_1 + a_2(x + \tilde{x}) + a_3(y + \tilde{y}) + a_4x\tilde{x} + a_5y\tilde{y} + a_6x\tilde{y} + a_7y\tilde{x}, \\ \tilde{y} - y &= b_1 + b_2(x + \tilde{x}) + b_3(y + \tilde{y}) + b_4x\tilde{x} + b_5y\tilde{y} + b_6x\tilde{y} + b_7y\tilde{x}.\end{aligned}$$

# Singularities of birational quadratic maps of $\mathbb{P}^2$

- ▶ Singularities:  $\mathcal{I}(\phi) = \{p_1, p_2, p_3\}$ ,  $\mathcal{I}(\phi^{-1}) = \{q_1, q_2, q_3\}$ .
- ▶  $\phi$  blows up points  $p_1, p_2, p_3$  to lines  $(q_2 q_3)$ ,  $(q_1 q_3)$ ,  $(q_1 q_2)$ , resp.
- ▶  $\phi$  blows down lines  $(p_2 p_3)$ ,  $(p_1 p_3)$ ,  $(p_1 p_2)$  to points  $q_1, q_2, q_3$ , resp.

# Lifting to automorphism

**Definition.** Map  $\phi$  is *confining*, if all three lines  $(p_j p_k)$  are *degree lowering* (i.e., participate in *singularity confinement patterns*):

$$(p_j p_k) \rightarrow q_i \rightarrow \phi(q_i) \rightarrow \cdots \rightarrow \phi^{n_i-1}(q_i) = p_{\sigma_i} \rightarrow (q_{\sigma_j} q_{\sigma_k}).$$

*Orbit data* of a confining  $\phi$  consist of  $(n_1, n_2, n_3), (\sigma_1, \sigma_2, \sigma_3)$ .

A confining map  $\phi$  can be lifted to an automorphism  $\hat{\phi}$  of a surface  $S$  obtained from  $\mathbb{P}^2$  by blowing up all participating points.

Dynamical degree  $\lambda_1(\phi)$  can be found as the spectral radius of the action of  $\hat{\phi}^*$  on  $\text{Pic}(S)$ .

**Theorem** [Bedford, Kim' 2004]. For a confining map,  $\lambda_1(\phi)$  depends only on the orbit data associated to  $\phi$ .

## Example of integrable planar birational map: Kahan discretization of Hamiltonian systems

For  $n = 2$ , consider  $f(x, y) = J\nabla H(x, y)$ , with  $J = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$ .  
 $\Phi_f$  is a birational planar map with an invariant measure and an integral  $\Rightarrow$  completely integrable. Integral:

$$\tilde{H}(x, y, \epsilon) = \frac{C(x, y, \epsilon)}{D(x, y, \epsilon)},$$

where  $\deg C = 3$ ,  $\deg D = 2$ . Level sets:

$$\mathcal{E}_\lambda = \{(x, y) : C(x, y, \epsilon) - \lambda D(x, y, \epsilon) = 0\},$$

a pencil of cubic curves, characterized by its nine *base points*.  
On each invariant curve,  $\Phi_f$  induces a translation (relative to the addition law on this curve).

# Complexification, projectivization

Pencil

$$\bar{\mathcal{E}}_\lambda = \left\{ [x : y : z] \in \mathbb{CP}^2 : \bar{C}(x, y, z, \epsilon) - \lambda z \bar{D}(x, y, z, \epsilon) = 0 \right\}.$$

spanned by two curves,

$$\bar{\mathcal{E}}_0 = \left\{ [x : y : z] \in \mathbb{CP}^2 : \bar{C}(x, y, z, \epsilon) = 0 \right\},$$

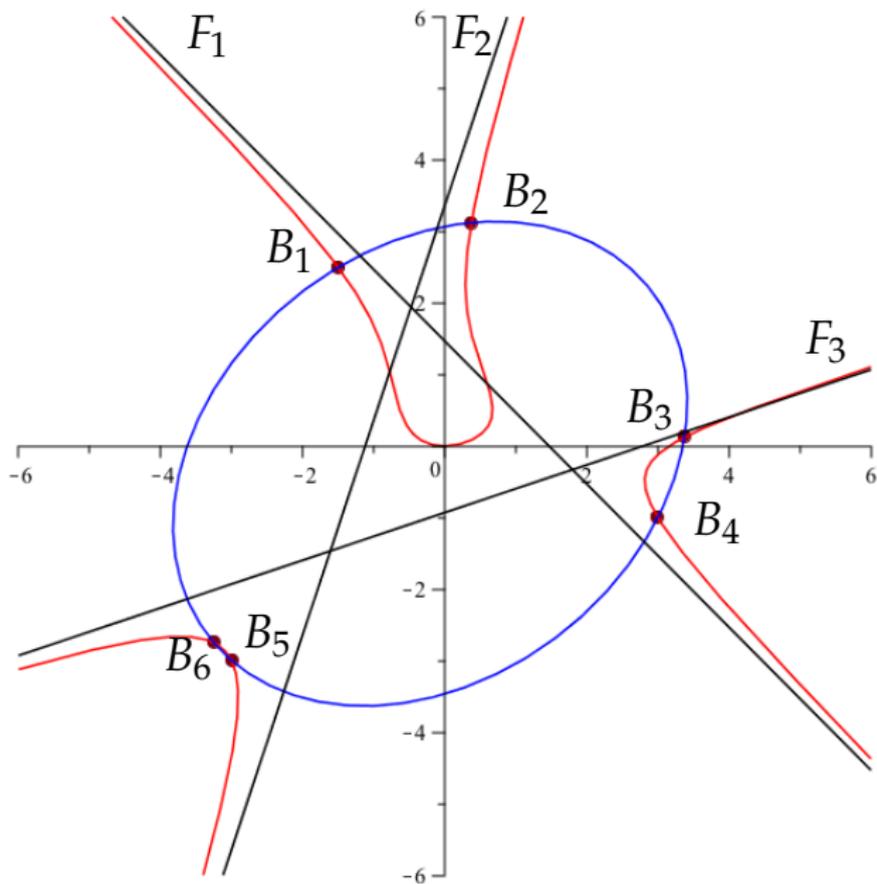
assumed nonsingular, and

$$\bar{\mathcal{E}}_\infty = \left\{ [x : y : z] \in \mathbb{CP}^2 : z \bar{D}(x, y, z, \epsilon) = 0 \right\}$$

reducible, consisting of conic  $\{\bar{D}(x, y, z, \epsilon) = 0\}$  and the line at infinity  $\{z = 0\}$ . Three base points at infinity:

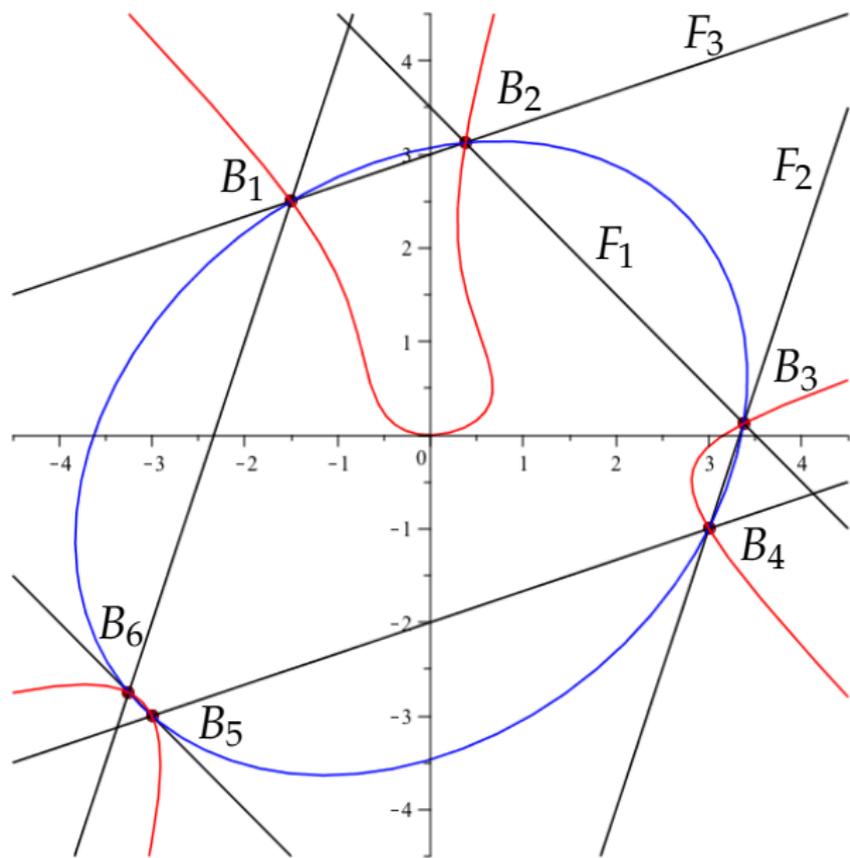
$$\{F_1, F_2, F_3\} = \bar{\mathcal{E}}_0 \cap \{z = 0\},$$

and six further base points  $\{B_1, \dots, B_6\} = \bar{\mathcal{E}}_0 \cap \{\bar{D} = 0\}$ .



- ▶ M. Petrera, J. Smirin, Yu. S. *Geometry of the Kahan discretizations of planar quadratic Hamiltonian systems.* Proc. R. Soc. A **476** (2019) 20180761

**Theorem.** *A pencil of elliptic curves consists of invariant curves for Kahan's discretization of a planar quadratic Hamiltonian vector field iff the hexagon through the six finite base points has three pairs of parallel sides which pass through the three base points at infinity.*



# Manin involutions for cubic curves

**Definition.** Consider a nonsingular cubic curve  $\bar{\mathcal{E}}$  in  $\mathbb{CP}^2$ .

- For a point  $P_0 \in \bar{\mathcal{E}}$ , the *Manin involution*  $I_{\bar{\mathcal{E}}, P_0} : \bar{\mathcal{E}} \rightarrow \bar{\mathcal{E}}$  is defined as follows:
  - ▶ For  $P \neq P_0$ , the point  $\bar{P} = I_{\bar{\mathcal{E}}, P_0}(P)$  is the unique third intersection point of  $\bar{\mathcal{E}}$  with the line  $(P_0P)$ ;
  - ▶ For  $P = P_0$ , the point  $\bar{P} = I_{\bar{\mathcal{E}}, P_0}(P)$  is the unique second intersection point of  $\bar{\mathcal{E}}$  with the tangent line to  $\bar{\mathcal{E}}$  at  $P = P_0$ .
- For two distinct points  $P_0, P_1 \in \bar{\mathcal{E}}$ , the *Manin transformation*  $M_{\bar{\mathcal{E}}, P_0, P_1} : \bar{\mathcal{E}} \rightarrow \bar{\mathcal{E}}$  is defined as

$$M_{\bar{\mathcal{E}}, P_0, P_1} = I_{\bar{\mathcal{E}}, P_1} \circ I_{\bar{\mathcal{E}}, P_0}.$$

With a natural addition law on  $\bar{\mathcal{E}}$ :

$$I_{\bar{\mathcal{E}}, P_0}(P) = -(P_0 + P), \quad M_{\bar{\mathcal{E}}, P_0, P_1}(P) = P + P_0 - P_1.$$

# Manin involutions for cubic pencils

**Definition.** Consider a pencil  $\mathfrak{E} = \{\bar{\mathcal{E}}_\lambda\}$  of cubic curves in  $\mathbb{CP}^2$  with at least one nonsingular member.

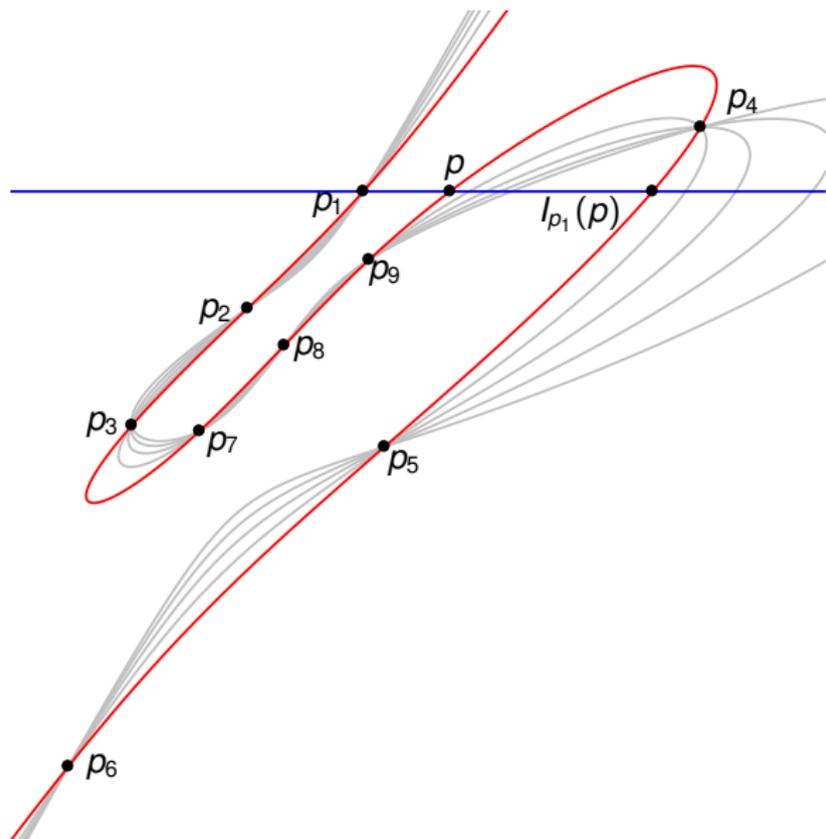
- Let  $B$  be a base point of the pencil. The *Manin involution*  $I_{\mathfrak{E},B} : \mathbb{CP}^2 \dashrightarrow \mathbb{CP}^2$  is a birational map defined as follows. For any  $P \in \mathbb{CP}^2$ , not a base point of  $\mathfrak{E}$ , let  $\bar{\mathcal{E}}_\lambda$  be the unique curve of  $\mathfrak{E}$  through  $P$ . Set

$$I_{\mathfrak{E},B}(P) = I_{\bar{\mathcal{E}}_\lambda,B}(P).$$

- Let  $B_1, B_2$  be two distinct base points of the pencil. The *Manin transformation*  $M_{\mathfrak{E},B_1,B_2} : \mathbb{CP}^2 \dashrightarrow \mathbb{CP}^2$  is a birational map defined as

$$M_{\mathfrak{E},B_1,B_2} = I_{\mathfrak{E},B_2} \circ I_{\mathfrak{E},B_1}.$$

# Manin involutions for cubic pencils



## Direct statement. Proof.

First one shows the Kahan map  $\Phi_f$  is a Manin transformation in six different ways:

$$\begin{aligned}\Phi_f &= l_{\mathcal{C}, B_1} \circ l_{\mathcal{C}, F_1} = l_{\mathcal{C}, F_1} \circ l_{\mathcal{C}, B_4} \\ &= l_{\mathcal{C}, B_5} \circ l_{\mathcal{C}, F_2} = l_{\mathcal{C}, F_2} \circ l_{\mathcal{C}, B_2} \\ &= l_{\mathcal{C}, B_3} \circ l_{\mathcal{C}, F_3} = l_{\mathcal{C}, F_3} \circ l_{\mathcal{C}, B_6}.\end{aligned}$$

Thus (on any invariant curve of  $\mathcal{C}$ ):

$$F_1 - B_1 = B_2 - F_2 = F_3 - B_3 = B_4 - F_1 = F_2 - B_5 = B_6 - F_3,$$

and

$$F_1 + F_2 + F_3 = O.$$

Have, e.g.:

$$B_1 + B_2 = F_1 + F_2 = -F_3 \quad \Leftrightarrow \quad B_1 + B_2 + F_3 = O.$$

Thus, line  $(B_1 B_2)$  passes through  $F_3$ . □

## Inverse statement. Proof.

Prescribe arbitrary nine coefficients of the side lines of the hexagon (three slopes  $\mu_1, \mu_2, \mu_3$  and six heights  $b_1, \dots, b_6$ ):

$$\begin{aligned}(B_1B_2) : y &= \mu_3x + b_1, & (B_4B_5) : y &= \mu_3x + b_4, \\(B_2B_3) : y &= \mu_1x + b_2, & (B_5B_6) : y &= \mu_1x + b_5, \\(B_3B_4) : y &= \mu_2x + b_3, & (B_6B_1) : y &= \mu_2x + b_6.\end{aligned}$$

This defines nine points  $B_1, \dots, B_6$  and  $F_1, F_2, F_3$ , therefore a pencil  $\mathfrak{E}$  of cubic curves with those nine base points. Set

$$\begin{aligned}\Phi &= l_{\mathfrak{E}, B_1} \circ l_{\mathfrak{E}, F_1} = l_{\mathfrak{E}, F_1} \circ l_{\mathfrak{E}, B_4} \\ &= l_{\mathfrak{E}, B_5} \circ l_{\mathfrak{E}, F_2} = l_{\mathfrak{E}, F_2} \circ l_{\mathfrak{E}, B_2} \\ &= l_{\mathfrak{E}, B_3} \circ l_{\mathfrak{E}, F_3} = l_{\mathfrak{E}, F_3} \circ l_{\mathfrak{E}, B_6}.\end{aligned}$$

This is a birational map of  $\mathbb{C}\mathbb{P}^2$  of degree 2. Check that this is a Kahan discretization of  $f = J\nabla H$  with  $\deg H = 3$ .

# Inverse statement. Proof.

Explicit expression:

$$\begin{aligned} H(x, y) = & \frac{2\mu_{12}}{b_{14}\mu_{23}\mu_{13}} \left( \frac{1}{3}(\mu_3 x - y)^3 + \frac{1}{2}(b_1 + b_4)(\mu_3 x - y)^2 + b_1 b_4(\mu_3 x - y) \right) \\ & - \frac{2\mu_{23}}{b_{25}\mu_{12}\mu_{13}} \left( \frac{1}{3}(\mu_1 x - y)^3 + \frac{1}{2}(b_2 + b_5)(\mu_1 x - y)^2 + b_2 b_5(\mu_1 x - y) \right) \\ & + \frac{2\mu_{13}}{b_{36}\mu_{12}\mu_{23}} \left( \frac{1}{3}(\mu_2 x - y)^3 + \frac{1}{2}(b_3 + b_6)(\mu_2 x - y)^2 + b_3 b_6(\mu_2 x - y) \right), \end{aligned}$$

where  $b_{ij} = b_i - b_j$ ,  $\mu_{ij} = \mu_i - \mu_j$ .

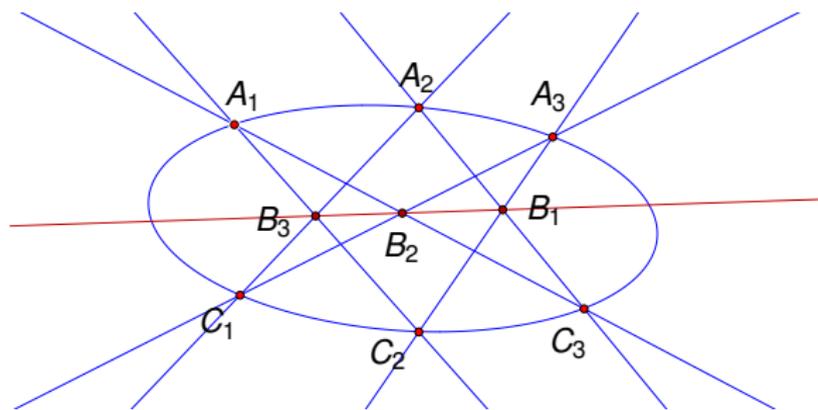
Geometry implies dynamics!



# Projective generalization of Hamiltonian case

Pascal configuration: six points  $A_1, A_2, A_3, C_1, C_2, C_3$  on a conic  $\mathcal{C}$ , and three intersection points on a line  $\ell$ :

$$B_1 = (A_2C_3) \cap (A_3C_2), \quad B_2 = (A_3C_1) \cap (A_1C_3), \quad B_3 = (A_1C_2) \cap (A_2C_1).$$



Consider pencil  $\mathfrak{E}$  of cubic curves passing through the nine points  $A_i, C_i, B_i$  (contains a reducible cubic  $\mathcal{C} \cup \ell$ ).

**Theorem** [S. 2020]. *The map*

$$\begin{aligned}\Phi &= I_{\mathcal{C}, A_1} \circ I_{\mathcal{C}, B_1} = I_{\mathcal{C}, B_1} \circ I_{\mathcal{C}, C_1} \\ &= I_{\mathcal{C}, A_2} \circ I_{\mathcal{C}, B_2} = I_{\mathcal{C}, B_2} \circ I_{\mathcal{C}, C_2} \\ &= I_{\mathcal{C}, A_3} \circ I_{\mathcal{C}, B_3} = I_{\mathcal{C}, B_3} \circ I_{\mathcal{C}, C_3}\end{aligned}$$

*is a birational map of degree 2 with*

- ▶  $\mathcal{I}(\Phi) = \{C_1, C_2, C_3\}$ , blown up to lines  $c_1 = (A_2A_3)$ ,  $c_2 = (A_3A_1)$ ,  $c_3 = (A_1A_2)$ ,
- ▶  $\mathcal{C}(\Phi)$  consisting of three lines  $a_1 = (C_2C_3)$ ,  $a_2 = (C_3C_1)$ ,  $a_3 = (C_1C_2)$ , blown down to points  $A_1, A_2, A_3$ .

*Singularity confinement patterns of the map  $\Phi$ :*

$$\begin{aligned}(C_2C_3) &\rightarrow A_1 \rightarrow B_1 \rightarrow C_1 \rightarrow (A_2A_3), \\ (C_3C_1) &\rightarrow A_2 \rightarrow B_2 \rightarrow C_2 \rightarrow (A_3A_1), \\ (C_1C_2) &\rightarrow A_3 \rightarrow B_3 \rightarrow C_3 \rightarrow (A_1A_2).\end{aligned}$$

To show: why the six Manin transformations correspond to one and the same translation on any curve of the pencil:

$$A_1 - B_1 = B_1 - C_1 = A_2 - B_2 = B_2 - C_2 = A_3 - B_3 = B_3 - C_3.$$

Collinearities of Pascal configuration are translated to:

$$\begin{aligned}A_2 + C_3 + B_1 &= O, & A_3 + C_2 + B_1 &= O, \\A_3 + C_1 + B_2 &= O, & A_1 + C_3 + B_2 &= O, \\A_1 + C_2 + B_3 &= O, & A_2 + C_1 + B_3 &= O,\end{aligned}$$

and

$$B_1 + B_2 + B_3 = O.$$

$$\begin{aligned}\text{Now: } A_1 + C_1 &= -(C_2 + B_3) - (A_3 + B_2) \\ &= -(A_3 + C_2) - (B_2 + B_3) = B_1 + B_1,\end{aligned}$$

which proves that  $A_1 - B_1 = B_1 - C_1$ . Similarly,

$$A_2 + C_1 = -B_3 = B_1 + B_2,$$

which proves that  $B_1 - C_1 = A_2 - B_2$ .

All other equations follow in the same way. □

## An early example

R. Penrose, C. Smith. *A quadratic mapping with invariant cubic curve*. Math. Proc. Camb. Phil. Soc. **89** (1981), 89–105:

$$\Phi : \begin{bmatrix} x_0 \\ x_1 \\ x_2 \end{bmatrix} \mapsto \begin{bmatrix} x_0(x_0 + ax_1 + a^{-1}x_2) \\ x_1(x_1 + ax_2 + a^{-1}x_0) \\ x_2(x_2 + ax_0 + a^{-1}x_1) \end{bmatrix}$$

with

$$A_1 = [0 : 1 : -a], \quad C_1 = [0 : a : -1], \quad B_1 = [0 : 1 : -1]$$

(and others cyclically). Upon a projective transformation sending  $B_1, B_2, B_3$  to infinity, get a Kahan discretization of a Hamiltonian vector field with  $H(x, y) = xy(1 - x - y)$  with the time step  $\epsilon = (a - 1)/(a + 1)$ .

## Further examples: $(\gamma_1, \gamma_2, \gamma_3)$ -family of 2d quadratic systems

$$\begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix} = \frac{1}{l_1^{\gamma_1-1} l_2^{\gamma_2-1} l_3^{\gamma_3-1}} J \nabla H,$$

where

$$J = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}, \quad H(x, y) = (l_1(x, y))^{\gamma_1} (l_2(x, y))^{\gamma_2} (l_3(x, y))^{\gamma_3},$$

$l_i(x, y) = a_i x + b_i y$  are linear forms, and  $\gamma_1, \gamma_2, \gamma_3 \in \mathbb{R}$ .

# Origin: reduced Nahm equations for symmetric monopoles [Hitchin, Manton, Murray' 1995]

- Tetrahedral symmetry,  $(\gamma_1, \gamma_2, \gamma_3) = (1, 1, 1)$ :

$$\begin{cases} \dot{x} = x^2 - y^2, \\ \dot{y} = -2xy, \end{cases} \quad H_1(x, y) = \frac{y}{3}(3x^2 - y^2).$$

- Octahedral symmetry,  $(\gamma_1, \gamma_2, \gamma_3) = (1, 1, 2)$ :

$$\begin{cases} \dot{x} = x^2 - 6y^2, \\ \dot{y} = -3xy - 2y^2, \end{cases} \quad H_2(x, y) = \frac{y}{2}(2x + 3y)(x - y)^2.$$

- Octahedral symmetry,  $(\gamma_1, \gamma_2, \gamma_3) = (1, 2, 3)$ :

$$\begin{cases} \dot{x} = 2x^2 - y^2, \\ \dot{y} = -10xy + y^2, \end{cases} \quad H_3(x, y) = \frac{y}{6}(3x - y)^2(4x + y)^3.$$

In all three cases all level sets  $H_i(x, y) = c$  are elliptic curves.

# The $(\gamma_1, \gamma_2, \gamma_3)$ -family: discretization

Hirota-Kimura-Kahan discretizations are integrable [Petrera, Pfadler, S. 2011]:

$$\begin{cases} \tilde{x} - x = \epsilon(\tilde{x}x - \tilde{y}y), \\ \tilde{y} - y = -\epsilon(\tilde{x}y + x\tilde{y}), \end{cases}$$

$$\begin{cases} \tilde{x} - x = \epsilon(2\tilde{x}x - 12\tilde{y}y), \\ \tilde{y} - y = -\epsilon(3\tilde{x}y + 3x\tilde{y} + 4\tilde{y}y), \end{cases}$$

$$\begin{cases} \tilde{x} - x = \epsilon(2\tilde{x}x - \tilde{y}y), \\ \tilde{y} - y = \epsilon(-5\tilde{x}y - 5x\tilde{y} + \tilde{y}y). \end{cases}$$

In all three cases, the map admits an invariant pencil of elliptic curves, of degrees 3, 4, and 6, respectively.

# The $(\gamma_1, \gamma_2, \gamma_3)$ -family: classification of integrable cases through discretization

**Theorem** [Zander' 2020]. *The only three cases when the Kahan discretization of the  $(\gamma_1, \gamma_2, \gamma_3)$ -system is confining, are  $(\gamma_1, \gamma_2, \gamma_3) = (1, 1, 1)$ ,  $(1, 1, 2)$ , and  $(1, 2, 3)$ . The orbit data in these three cases are:  $(\sigma_1, \sigma_2, \sigma_3) = (1, 2, 3)$  and, respectively,*

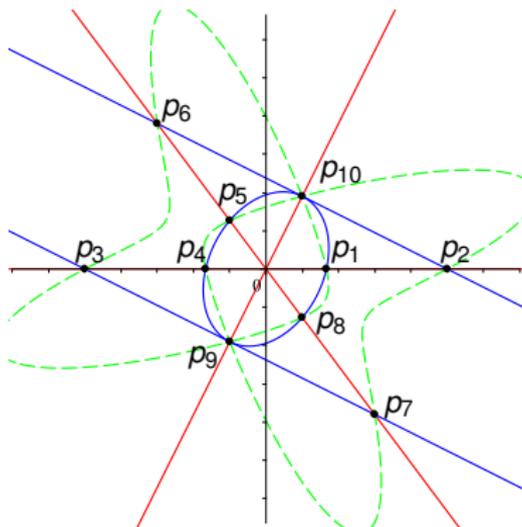
$$(n_1, n_2, n_3) = (3, 3, 3), (4, 4, 2), \text{ and } (6, 3, 2).$$

Observe: these  $(n_1, n_2, n_3)$  are the only positive integer solutions of

$$\frac{1}{n_1} + \frac{1}{n_2} + \frac{1}{n_3} = 1.$$

**Puzzle:** what do lengths of singularity confinement patterns have to do with tilings of the plane by congruent triangles???

# Kahan discretization for $(\gamma_1, \gamma_2, \gamma_3) = (1, 1, 2)$



# Kahan discretization for $(\gamma_1, \gamma_2, \gamma_3) = (1, 1, 2)$

- ▶ Invariant pencil consists of quartic curves with two double points:  $\mathfrak{E} = \mathcal{P}(4; p_1, \dots, p_8, p_9^2, p_{10}^2)$ .
- ▶  $\mathcal{I}(\phi) = \{p_4, p_8, p_{10}\}$ ,  $\mathcal{I}(\phi^{-1}) = \{p_1, p_5, p_9\}$ .
- ▶ Singularity confinement patterns:

$$(p_8 p_{10}) \rightarrow p_1 \rightarrow p_2 \rightarrow p_3 \rightarrow p_4 \rightarrow (p_5 p_9)$$

$$(p_4 p_{10}) \rightarrow p_5 \rightarrow p_6 \rightarrow p_7 \rightarrow p_8 \rightarrow (p_1 p_9)$$

$$(p_4 p_8) \rightarrow p_9 \rightarrow p_{10} \rightarrow (p_1 p_5)$$

- ▶ What is the geometric representation?

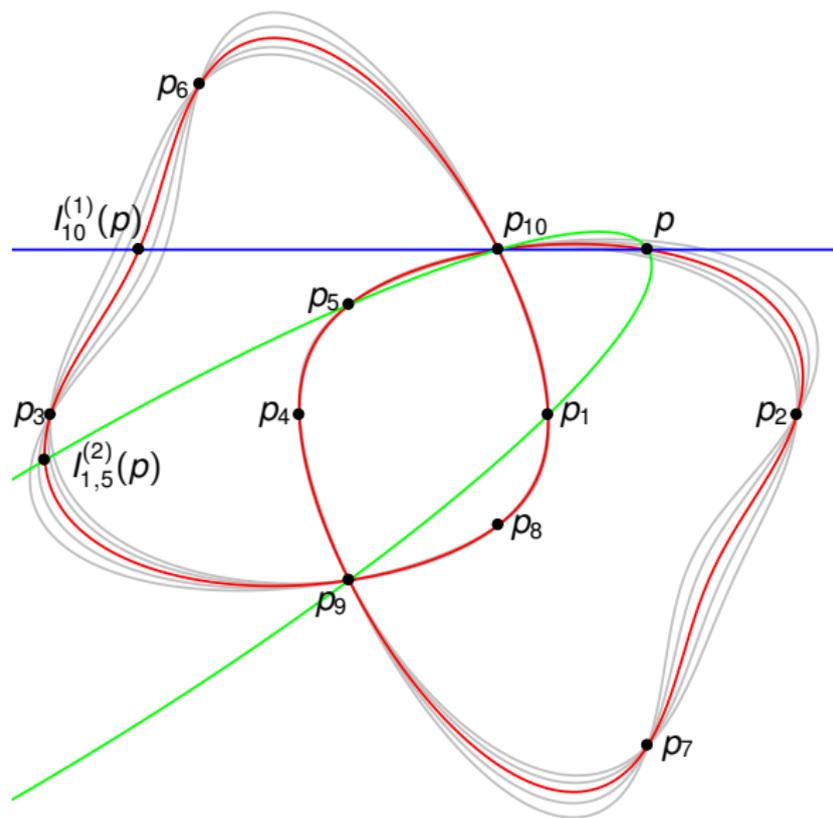
# Involutions for quartic pencils with two double points

Manin involutions for  $\mathcal{E} = \mathcal{P}(4; p_1, \dots, p_8, p_9^2, p_{10}^2)$ :

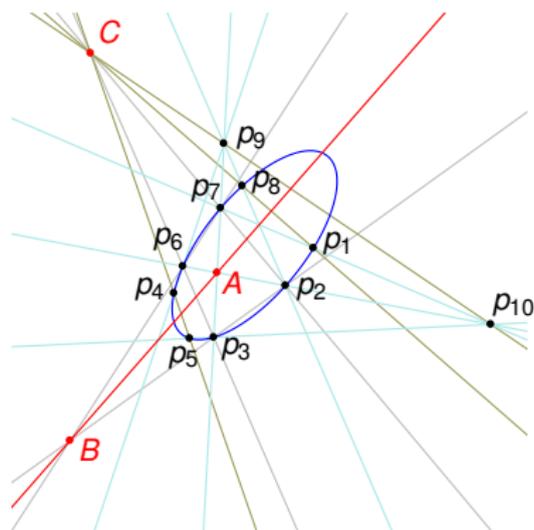
- ▶  $I_k^{(1)}$ ,  $k \in \{9, 10\}$ :  $I_k^{(1)}(p)$  is the third intersection point of the quartic through  $p$  with the line  $(pp_k)$ .
- ▶  $I_{i,j}^{(2)}$ ,  $i, j \in \{1, \dots, 8\}$ :  $I_{i,j}^{(2)}(p)$  is the sixth intersection point of the quartic through  $p$  with the conic through  $p_9, p_{10}, p_i, p_j, p$ .

Are derived from Manin involutions for a cubic pencil upon a quadratic Cremona transformation resolving both double points.

# Involutions for quartic pencils with two double points



# Quadratic Manin maps for special quartic pencils



Geometry of base points of a *projectively symmetric quartic pencil* with two double points  $\mathfrak{E} = \mathcal{P}(4; p_1, \dots, p_8, p_9^2, p_{10}^2)$ .

# Quadratic Manin maps for special quartic pencils

**Theorem** [Petrera, S., Wei, Zander' 2021].

1. The projective involution  $\sigma$  can be represented as

$$\sigma = I_{1,8}^{(2)} = I_{2,7}^{(2)} = I_{3,6}^{(2)} = I_{4,5}^{(2)}.$$

2. The map

$$\phi = I_{i,k}^{(2)} \circ I_{j,k}^{(2)} = I_9^{(1)} \circ \sigma = \sigma \circ I_{10}^{(1)},$$

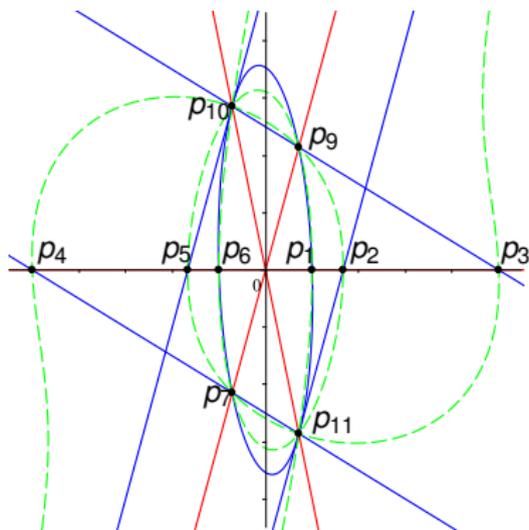
$(i, j) \in \{(1, 2), (2, 3), (3, 4), (5, 6), (6, 7), (7, 8)\}$  and  $k \in \{1, \dots, 8\}$  distinct from  $i, j$ , is a birational map of **degree 2**, with the singularity confinement patterns:

$$(p_8 p_{10}) \rightarrow p_1 \rightarrow p_2 \rightarrow p_3 \rightarrow p_4 \rightarrow (p_5 p_9),$$

$$(p_4 p_{10}) \rightarrow p_5 \rightarrow p_6 \rightarrow p_7 \rightarrow p_8 \rightarrow (p_1 p_9),$$

$$(p_4 p_8) \rightarrow p_9 \rightarrow p_{10} \rightarrow (p_1 p_5).$$

# Kahan discretization for $(\gamma_1, \gamma_2, \gamma_3) = (1, 2, 3)$



# Kahan discretization for $(\gamma_1, \gamma_2, \gamma_3) = (1, 2, 3)$

- ▶ Invariant pencil of sextic curves with 3 double points and 2 triple points:  $\mathcal{E} = \mathcal{P}(6; p_1, \dots, p_6, p_7^2, p_8^2, p_9^2, p_{10}^3, p_{11}^3)$ .
- ▶  $\mathcal{I}(\phi) = \{p_6, p_9, p_{11}\}$ ,  $\mathcal{I}(\phi^{-1}) = \{p_1, p_7, p_{10}\}$ .
- ▶ Singularity confinement patterns:

$$(p_9 p_{11}) \rightarrow p_1 \rightarrow p_2 \rightarrow p_3 \rightarrow p_4 \rightarrow p_5 \rightarrow p_6 \rightarrow (p_7 p_{10}),$$

$$(p_6 p_{11}) \rightarrow p_7 \rightarrow p_8 \rightarrow p_9 \rightarrow (p_1 p_{10}),$$

$$(p_6 p_9) \rightarrow p_{10} \rightarrow p_{11} \rightarrow (p_1 p_7).$$

- ▶ What is the geometric representation?

# Kahan discretization for $(\gamma_1, \gamma_2, \gamma_3) = (1, 2, 3)$

Manin involutions for  $\mathfrak{E} = \mathcal{P}(6; p_1, \dots, p_6, p_7^2, p_8^2, p_9^2, p_{10}^3, p_{11}^3)$ :

- ▶  $I_{i,j,k}^{(4)}$ ,  $i, j \in \{1, \dots, 6\}$ ,  $k \in \{7, 8, 9\}$ : e.g.,  $I_{i,j,9}^{(4)}$  is defined in terms of intersection of  $\mathfrak{E}$  with quartics of the pencil

$$\mathcal{P}(4; p_i, p_j, p_7, p_8, p_9^2, p_{10}^2, p_{11}^2).$$

- ▶  $I_{i,k}^{(3)}$ ,  $i \in \{1, \dots, 6\}$ ,  $k \in \{10, 11\}$ : e.g.,  $I_{i,10}^{(3)}$  is defined in terms of intersection of  $\mathfrak{E}$  with cubics of the pencil

$$\mathcal{P}(3; p_i, p_7, p_8, p_9, p_{10}^2, p_{11}).$$

**Theorem** [Petrera, S, Wei, Zander' 2021]. *The map  $\phi$  can be represented as compositions of (suitably defined) Manin involutions in the following ways:*

$$\phi = I_{i,k,m}^{(4)} \circ I_{j,k,m}^{(4)} = I_{i,n}^{(3)} \circ I_{j,n}^{(3)}$$

for any  $(i, j) \in \{(1, 2), (2, 3), (3, 4), (4, 5), (5, 6)\}$ ,  
 $k \in \{1, \dots, 6\} \setminus \{i, j\}$ , and  $m \in \{7, 8, 9\}$ ,  $n \in \{10, 11\}$ .

# Conclusions, work in progress and open problems

- ▶ Classification of integrable cases of Kahan discretization for the  $(\gamma_1, \gamma_2, \gamma_3)$ -family.
- ▶ Geometric construction of Manin involutions for pencils of elliptic curves of degree 4 and 6.
- ▶ Integrable Kahan discretizations for  $(\gamma_1, \gamma_2, \gamma_3) = (1, 1, 1)$ ,  $(1, 1, 2)$ ,  $(1, 2, 3)$  are Manin maps for pencils of elliptic curves of degree 3, 4, 6, resp.
- ▶ Special geometry of base points ensures  $\text{deg} = 2$  for certain Manin maps.
- ▶ Work in progress: singularity structure and geometric description for higher-dimensional examples, e.g., Kahan discretization of the Euler top (3D,  $g = 1$ ) or the Clebsch system (6D,  $g = 2$ ).