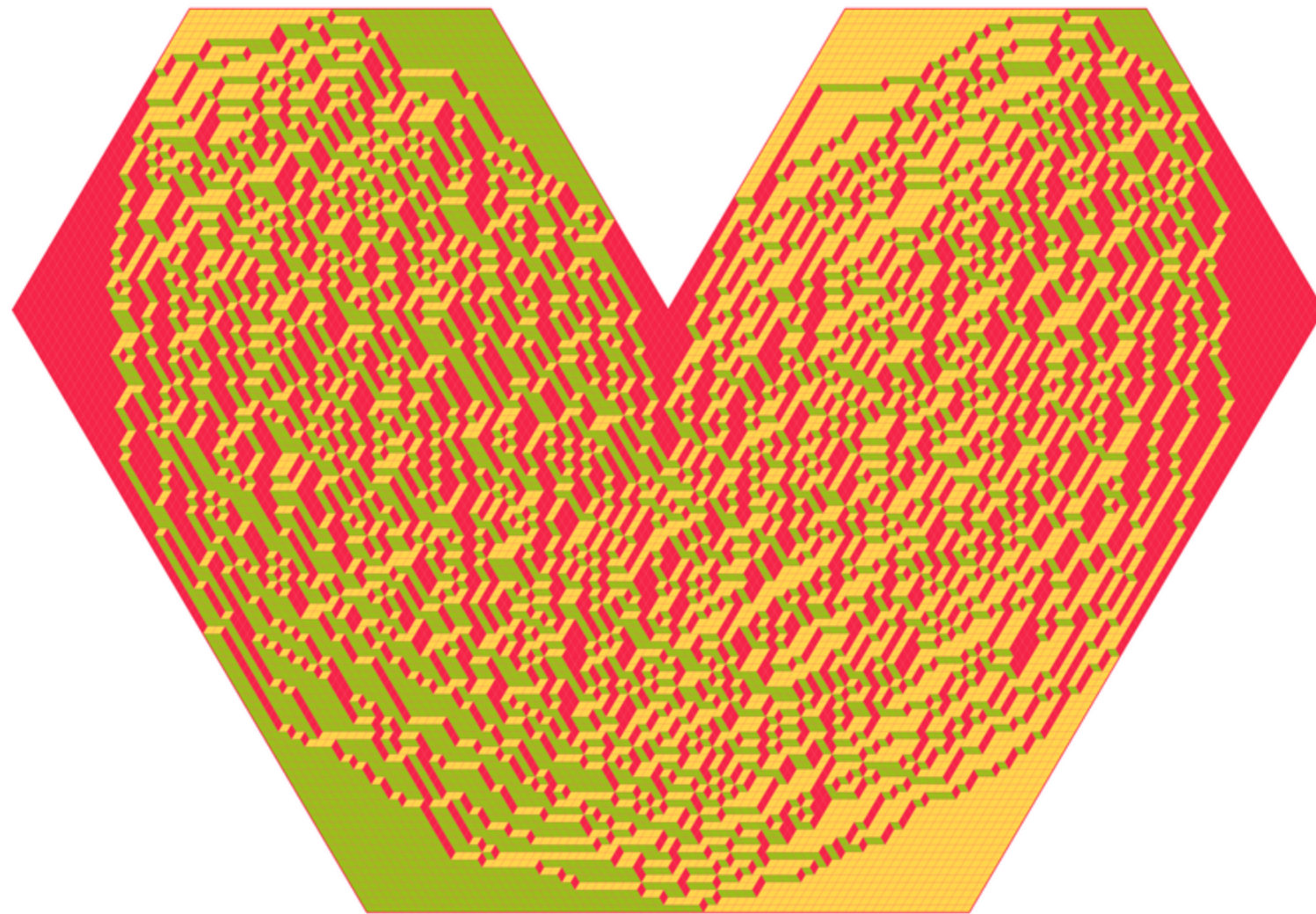


# Large random matrices are everywhere



Gaëtan Borot

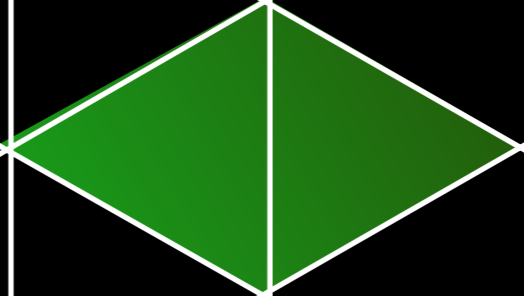
HU Institut für Mathematik & Institut für Physik



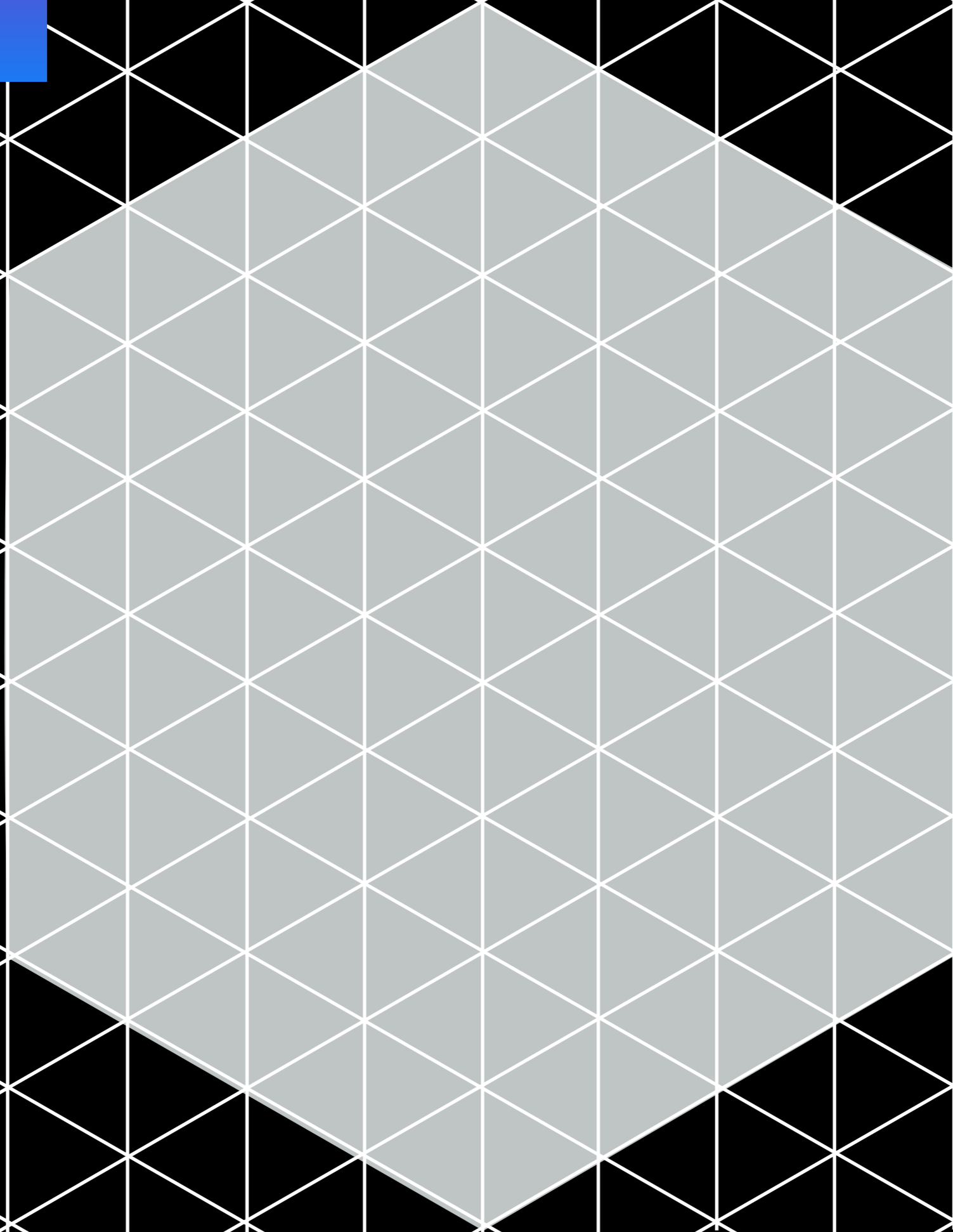
**1**

# Tilings and sugar melting

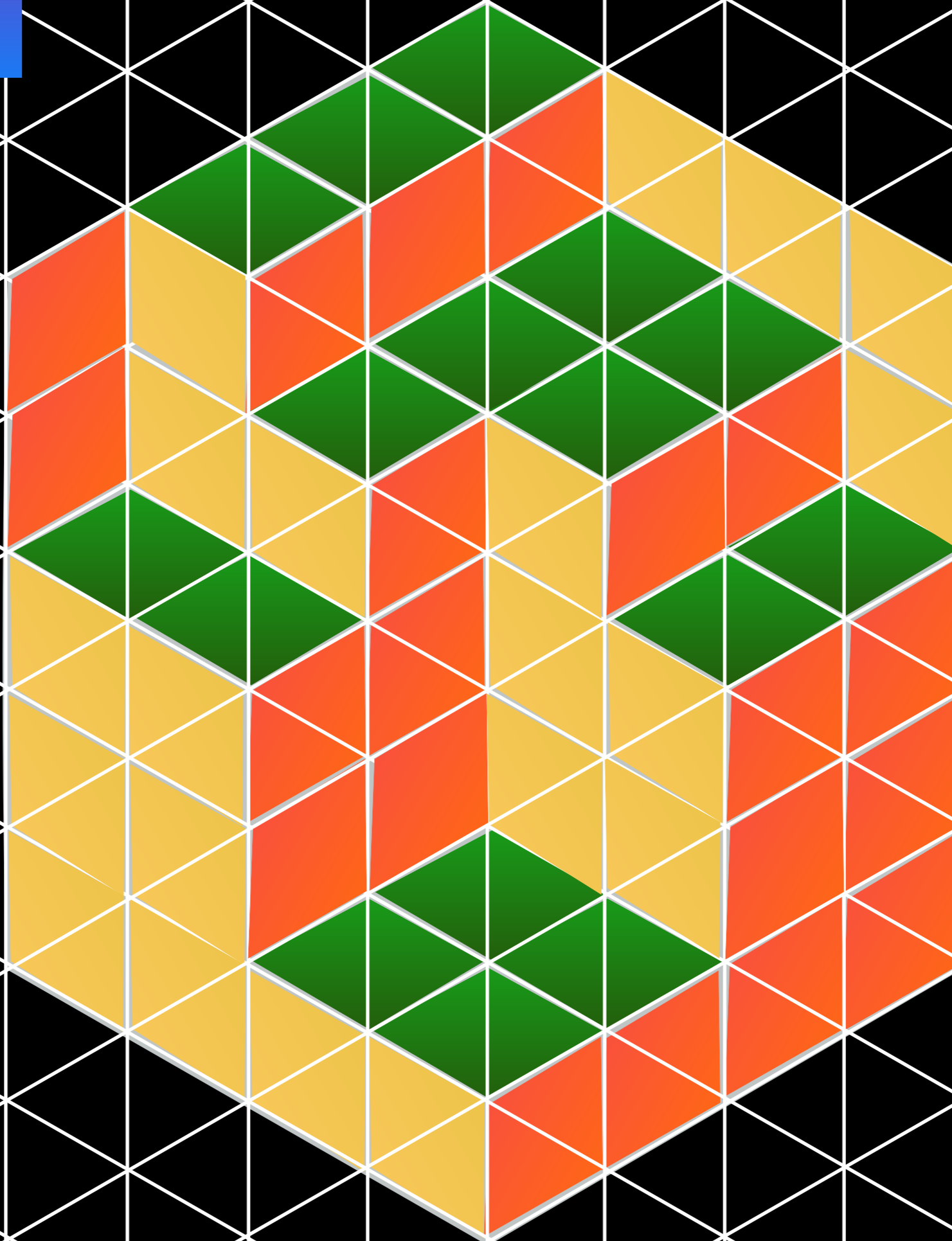
# 3 types of tiles



A domain

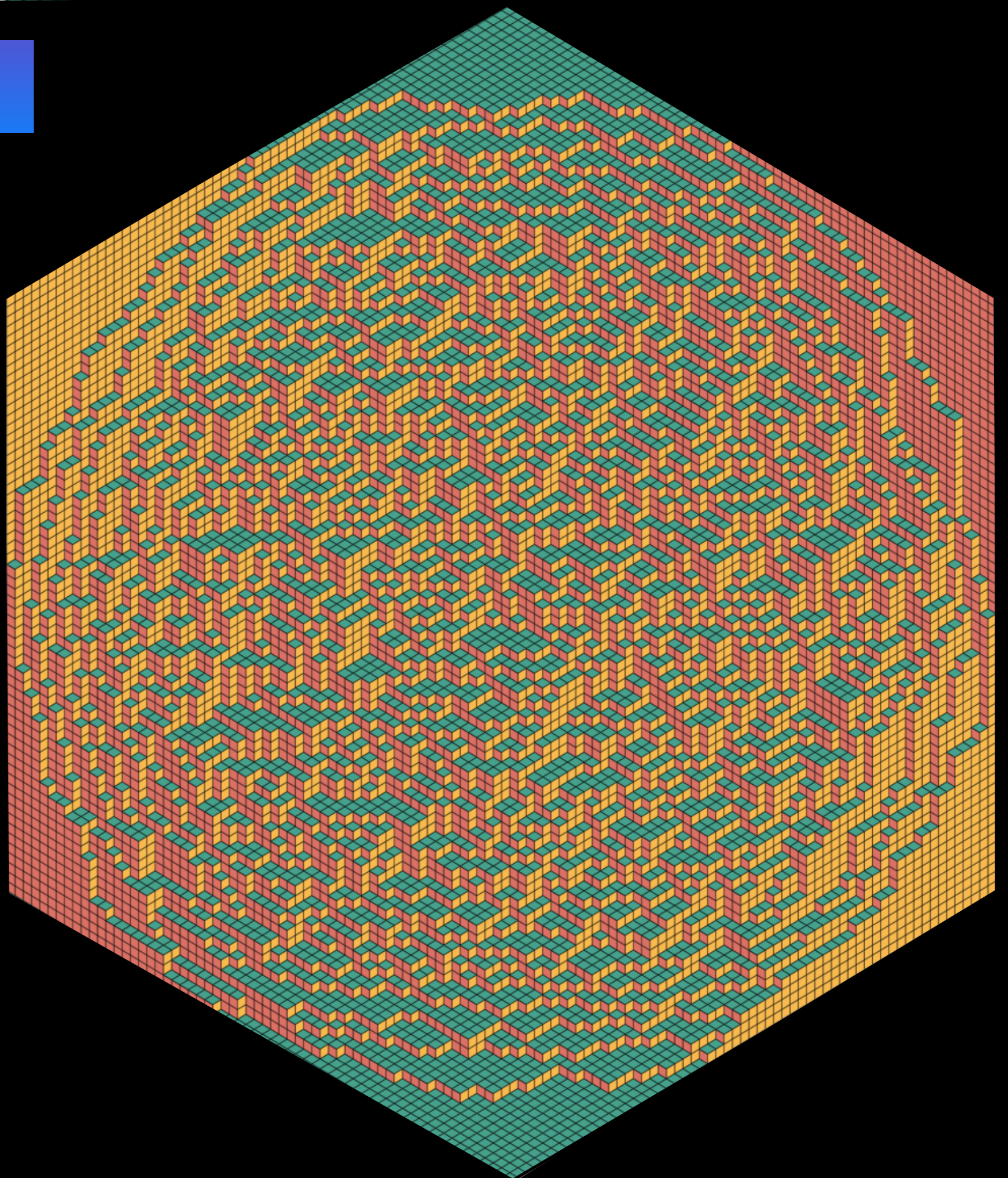


A tiling



# A large tiling

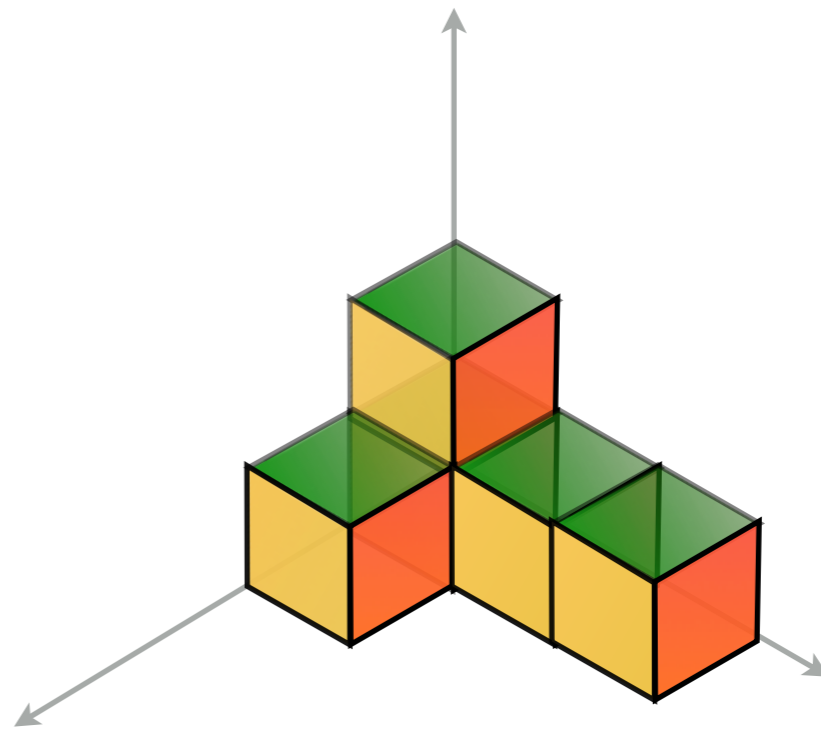
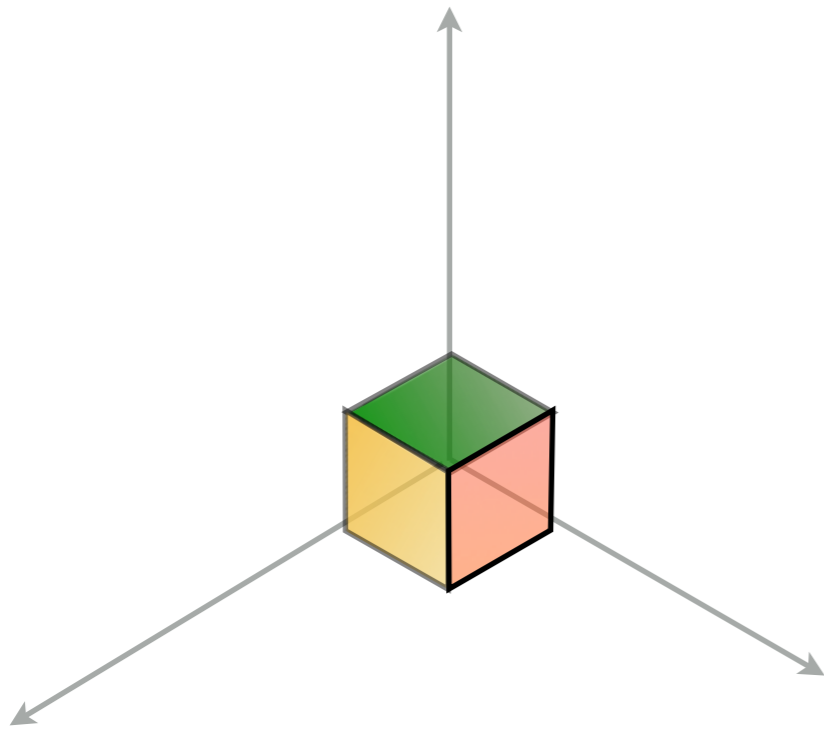
*Simulation by  
L. Petrov*



# 1. Random tilings

- These are 2d projections of a 3d picture

2d tiling  $\longleftrightarrow$  Piling up cubes in the corner of a room



# 1. Random tilings

- These are 2d projections of a 3d picture

2d tiling  $\longleftrightarrow$  Piling up cubes in the corner of a room

- Model for random piling / sugar melting

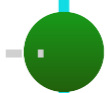
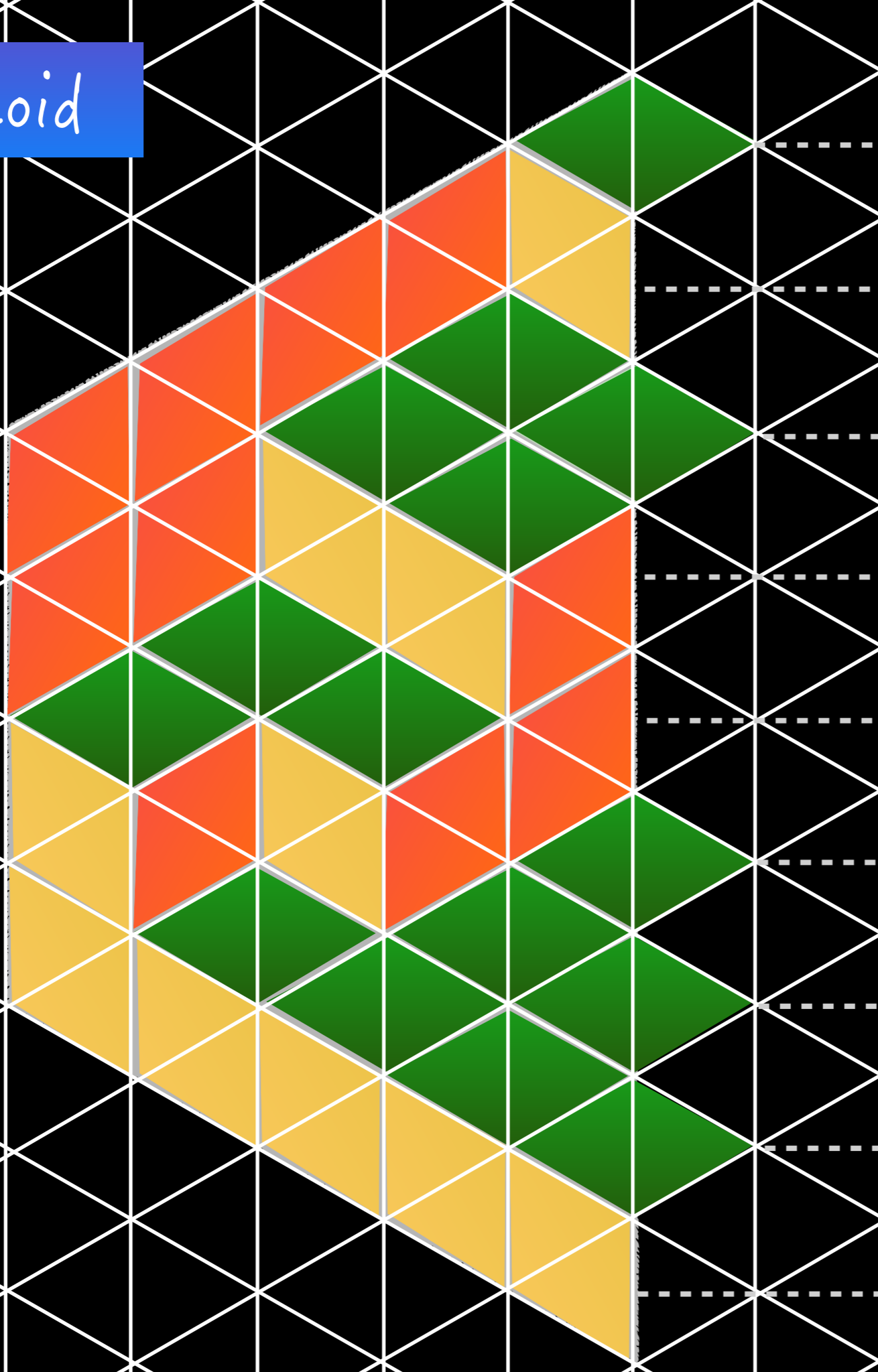
- Choose N randomly, with probability proportional to  $\exp\left(-\frac{\epsilon N}{k_B T}\right)$

(Boltzmann law:  $\epsilon$  energy per cube, T temperature)

- Choose a piling with N cubes, randomly



Trapezoid



$l_1$



$l_2$



$l_3$



$l_4$



$l_5$



# 1. Random tilings

(Gelfand-Tsetlin 1950, Cohn-Larsen-Propp 1998)

There are  $\prod_{i < j} \frac{\ell_i - \ell_j}{i - j}$  tilings of the trapezoid

having green tiles sticking out at positions  $\ell_1 > \ell_2 > \dots$

$\implies$  There are  $P(\ell) = \left( \prod_{i < j} \frac{\ell_i - \ell_j}{i - j} \right)^2$  tilings of the hexagon

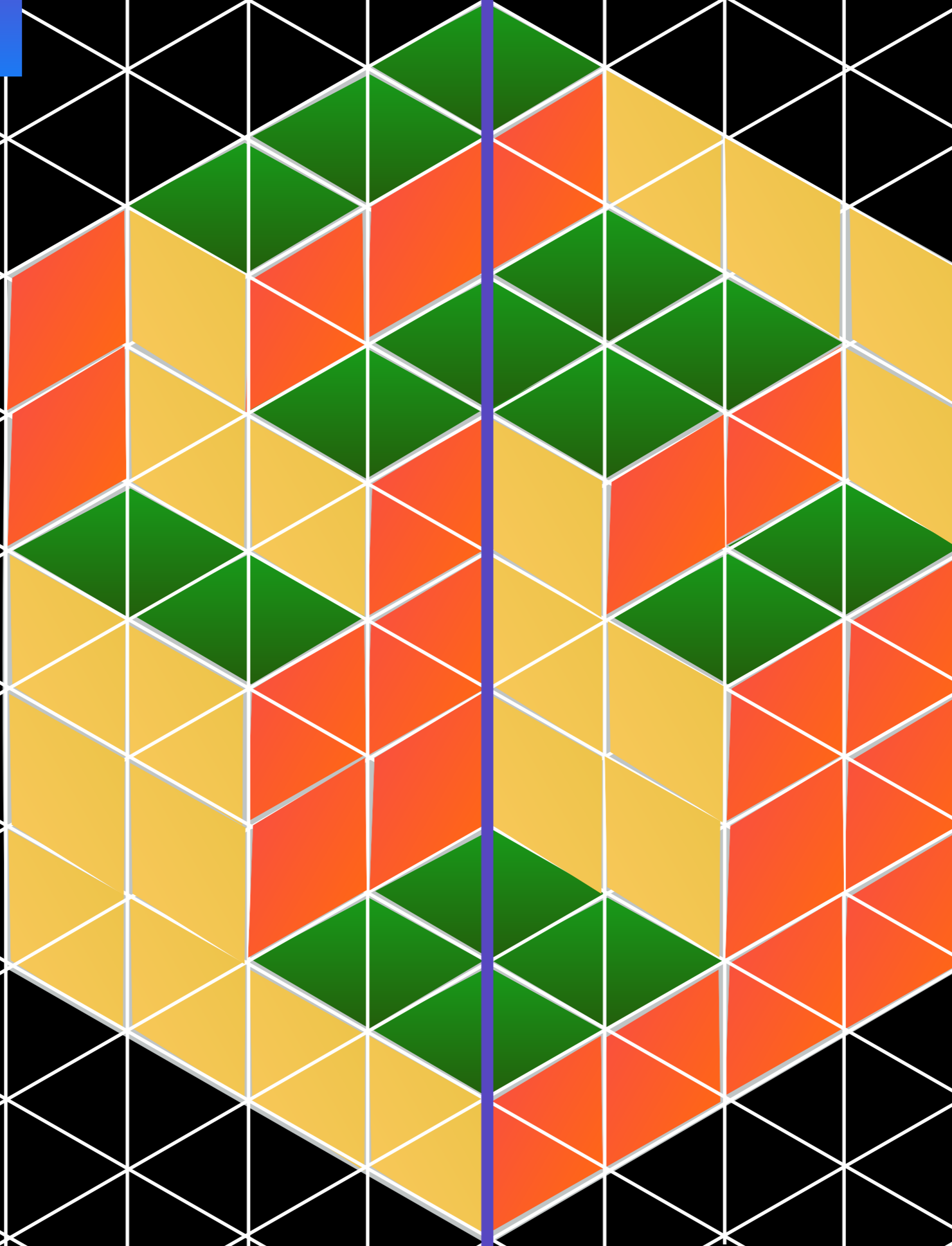
having green tiles crossing the vertical section

at positions  $\ell_1 > \ell_2 > \dots$

- Random choice of tiling of the hexagon

$\implies$  Probability to see  $\ell_1, \ell_2, \dots$  is proportional to  $P(\ell)$

A tiling



# 1. Random tilings

Random choice of tiling  $P(\ell) = \left( \prod_{i < j} \frac{\ell_i - \ell_j}{i - j} \right)^2$

Probability to observe  $\ell_i, \ell_j$  close to each other is small

→  $\ell_1, \ell_2, \dots$  are not independent from each other

They rather tend not to be close to each other !

# 1. Random tilings

For very large tilings chosen at random

**Observation:** Arctic circle phenomenon

there is a (non-random) curve  $c$  such that  
with probability  $\sim 1$  when  $N$  is large

- outside  $c$ : frozen tiles
- inside  $c$ : fluctuating tiles (surface looks rough)

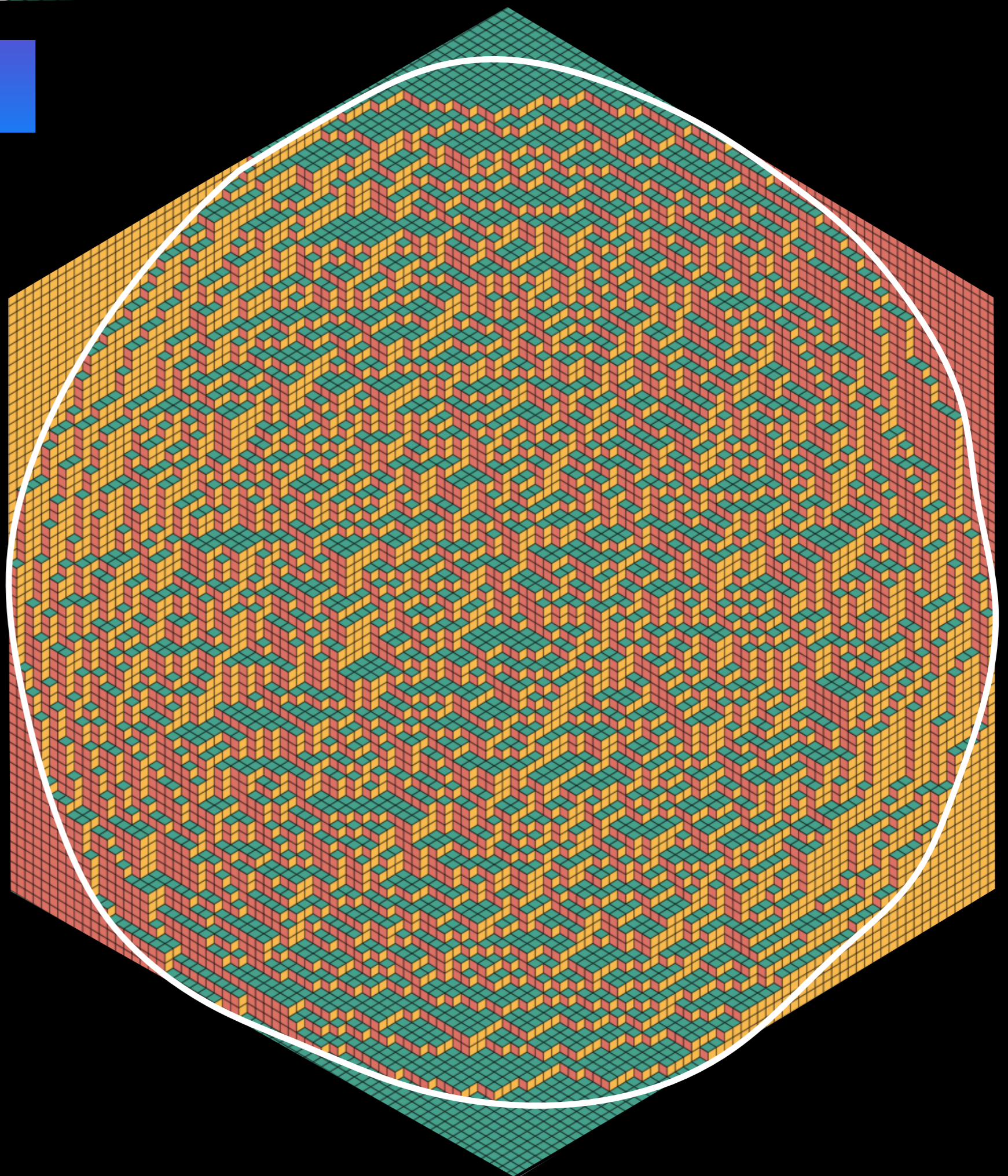
## Questions

- Describe  $c$  / the distribution of  $\frac{\ell_1}{\sqrt{N}}$
- describe the law of (microscopic) fluctuations inside  $c$

# A large tiling

Simulation by  
L. Petrov

C



**2**

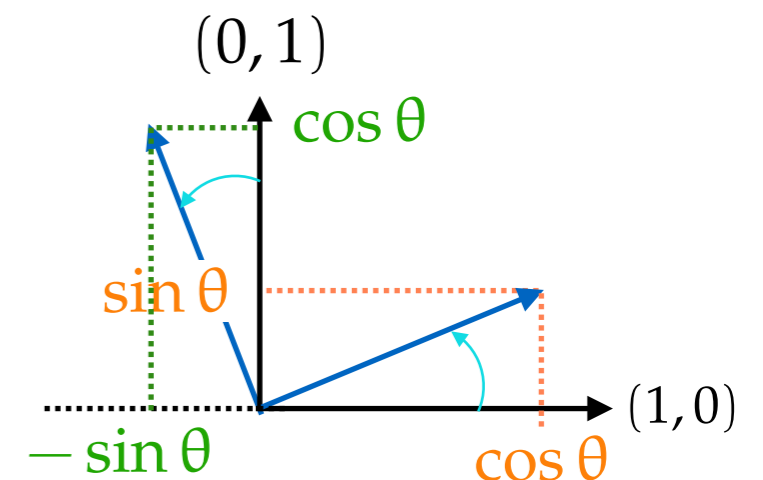
**Spectrum of (random) matrices**

## 2. Spectrum of (random) matrices

- A **matrix** is a table filled with numbers
- A matrix of size  $n \times n$  represents a linear transformation

$$\begin{pmatrix} A(1,1) & A(1,2) & \cdots & A(1,n) \\ A(2,1) & A(2,2) & \cdots & A(2,n) \\ \vdots & \vdots & \ddots & \vdots \\ A(n,1) & A(n,2) & \cdots & A(n,n) \end{pmatrix} \begin{array}{l} x_1 \mapsto A(1,1)x_1 + A(1,2)x_2 + \cdots + A(1,n)x_n \\ x_2 \mapsto A(2,1)x_1 + A(2,2)x_2 + \cdots + A(2,n)x_n \\ \vdots \\ x_n \mapsto A(n,1)x_1 + A(n,2)x_2 + \cdots + A(n,n)x_n \end{array}$$

- Example: for  $n = 2$   $\begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix}$   
represents the rotation of angle  $\theta$  in the plane



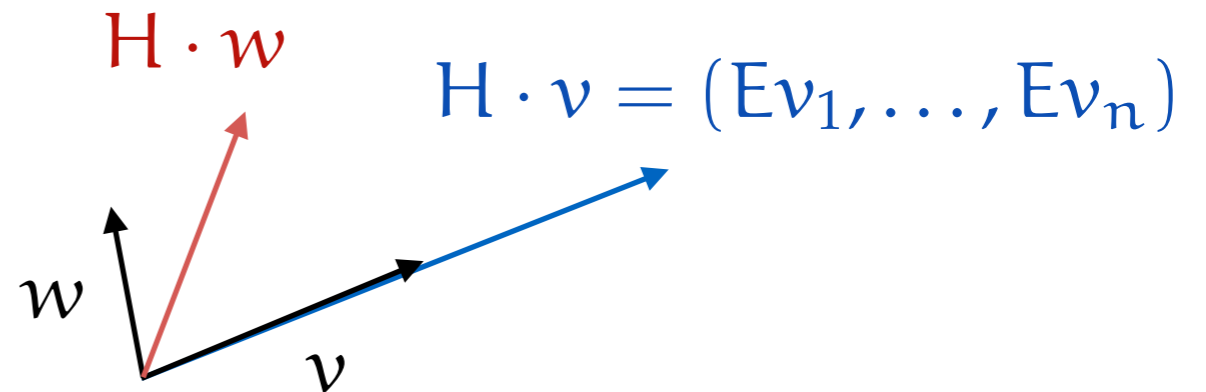


## 2. Spectrum of (random) matrices

- Given a matrix  $H$  of size  $n \times n$

$v = (v_1, \dots, v_n)$  is an eigenstate with eigenvalue  $E$

if  $H$  acts on  $v$  as  
rescaling by the factor  $E$



- Under certain assumptions on  $H$  (hermitian)

there are exactly  $n$  eigenstates (up to scale)

and  $n$  corresponding eigenvalues  $E_1 \geq \dots \geq E_n$

= spectrum of  $H$

## 2. Spectrum of (random) matrices

- **Random** (hermitian) matrix of size  $n \times n$   
→ **random spectrum**  $E_1 \geq \dots \geq E_n$
- When the random model does not have a preferred direction  
(1900 ... Dyson, Wigner, Mehta ... 1960)

Probability to find  $E_i$  near  $x_i$  at precision  $\delta \ll 1$

is proportional to  $\delta^n \cdot \rho(x) \cdot \prod_{i < j} (x_i - x_j)^2$

**Questions** When  $n$  is large

- How do  $E_1 \geq \dots \geq E_n$  distribute ?
- How does  $E_1 = \max(E_1, \dots, E_n)$  behave ?

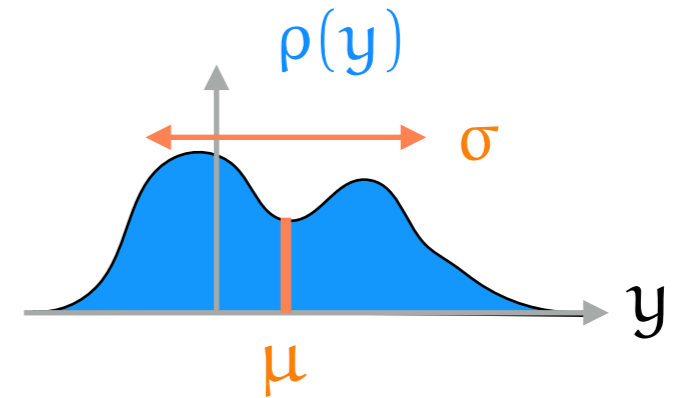
**3**

**What happens for  
independent random variables ?**

# 3. Independent random variables

Consider a random variable  $Y$

$$\mu = \text{Mean}(Y) \quad \sigma = \text{Variance}(Y)$$

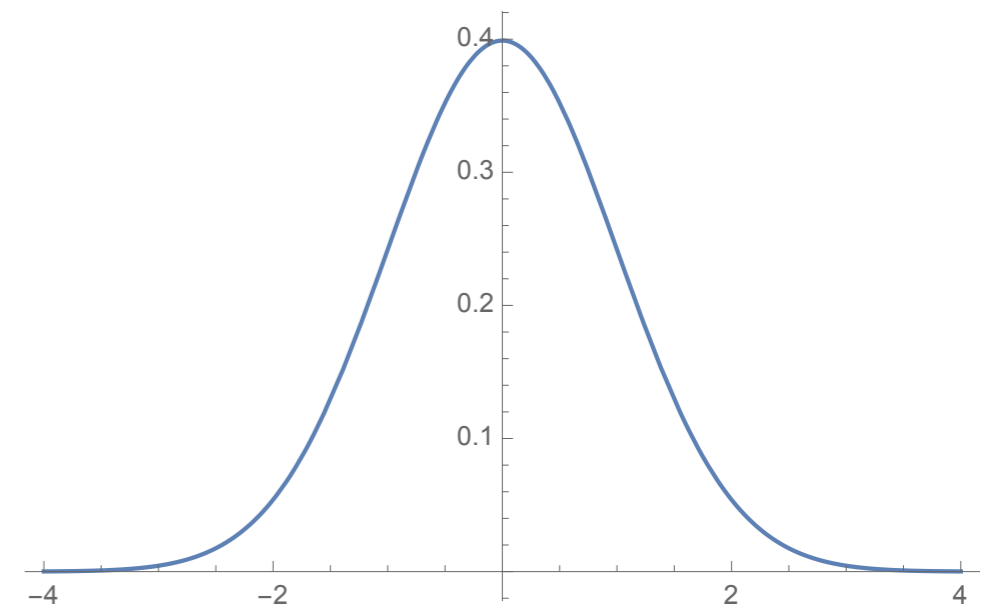


Make  $n$  independent measurements of it:  $y_1, \dots, y_n$

## Central limit theorem

$$y_1 + \dots + y_n \underset{n \text{ large}}{\approx} n\mu + \sqrt{n}\sigma s$$

$s$  random, with Gaussian distribution



(Gauß law)

# 3. Independent random variables

Make  $n$  independent measurements :  $y_1, \dots, y_n$

$$m_n = \max(y_1, \dots, y_n)$$

## Fluctuations of the maximum

(Fisher-Tippett 1928, Gumbel 1935, Gnedenko 1942)

Assume  $\rho(y) \underset{y \text{ large}}{\approx} y^\alpha e^{-\beta y^\gamma}$  with  $\beta, \gamma > 0$

Then  $m_n \underset{n \text{ large}}{\approx} a_n + b_n \xi$

Non-random  $a_n = \left(\frac{\ln(n)}{\beta}\right)^{1/\gamma}$  and  $b_n = (a_n)^{1-\gamma} \ll a_n$

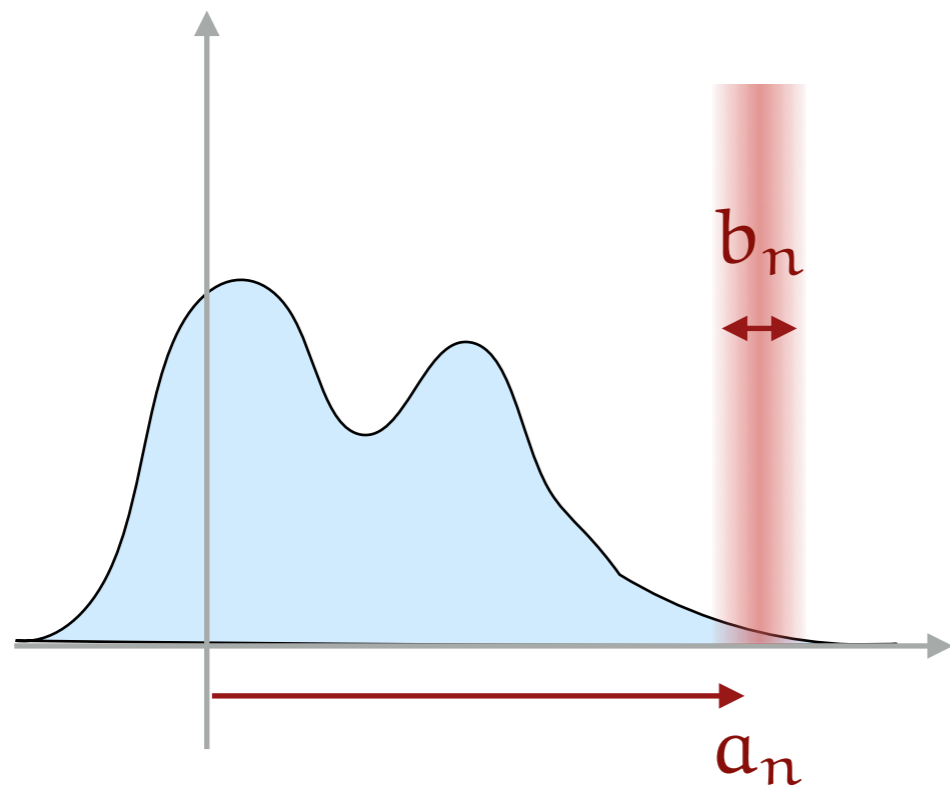
$\xi$  random with Gumbel distribution  $\rho(\xi) = \exp(-\xi - \exp(-\xi))$

# 3. Independent random variables

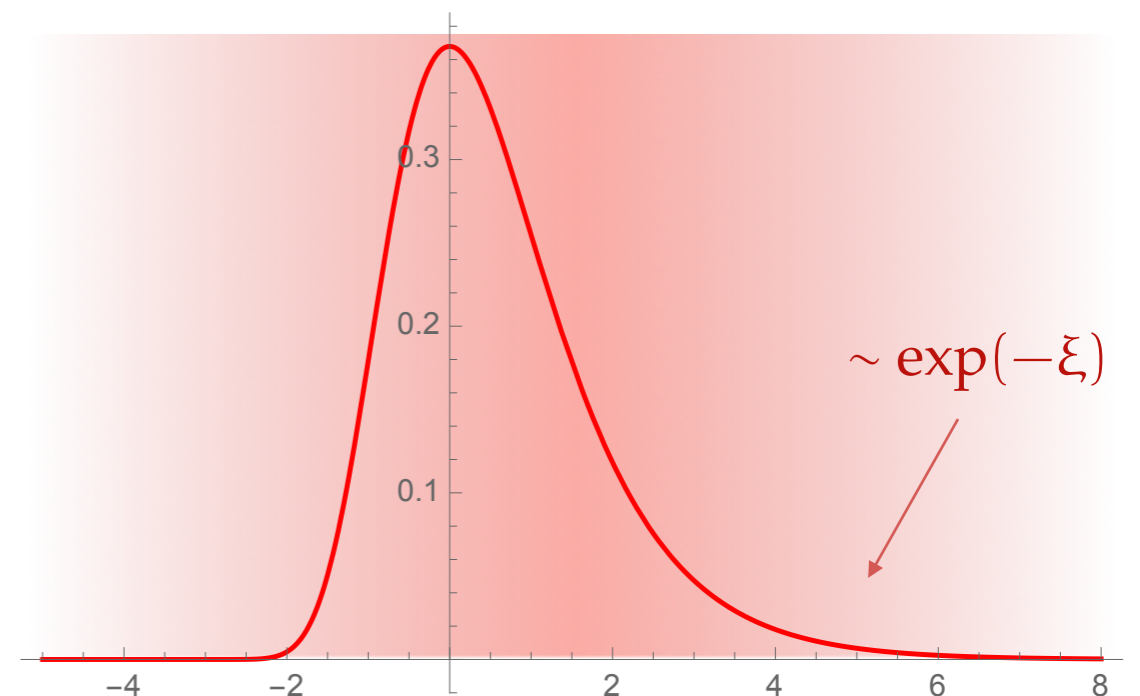
$$m_n = \max(y_1, \dots, y_n) = a_n + b_n \xi$$

$$a_n = \left( \frac{\ln(n)}{\beta} \right)^{1/\gamma}$$
$$b_n = (a_n)^{1-\gamma} \ll a_n$$

Prob. distribution of  $y$



Prob. distribution of  $\xi$



(Gumbel's law)

# 3. Independent random variables

## Summary

- The behavior of the sum and the max of independent  $n$  random variables with  $n$  large is well-understood since ~a century (Gauß's law, Gumbel's law, etc.)
- Universality phenomenon  
it depends very little on the random model considered  
(on the details of  $\rho(y)$ )

4

What happens for eigenvalues  
of large random matrices ?



# 4. Eigenvalues of large random matrices

- In random tilings  $\ell_1 > \dots > \ell_n$  are integers
- In random matrices  $E_1 \geq \dots \geq E_n$  are real numbers

In both cases, probability to find them around  $x_1, \dots, x_n$  at precision  $\delta \ll 1$  is proportional to  $\prod_{i < j} (x_i - x_j)^2$

→ these random variables are strongly correlated  
(  $x_i$  tends to avoid being close to  $x_j$  )

The previous laws (Gauß, Gumbel, etc.) do not apply !

# 4. Eigenvalues of large random matrices

Consider a random (hermitian) matrix of size  $n \times n$  with independent Gauß-distributed entries with variance  $\sigma^2$

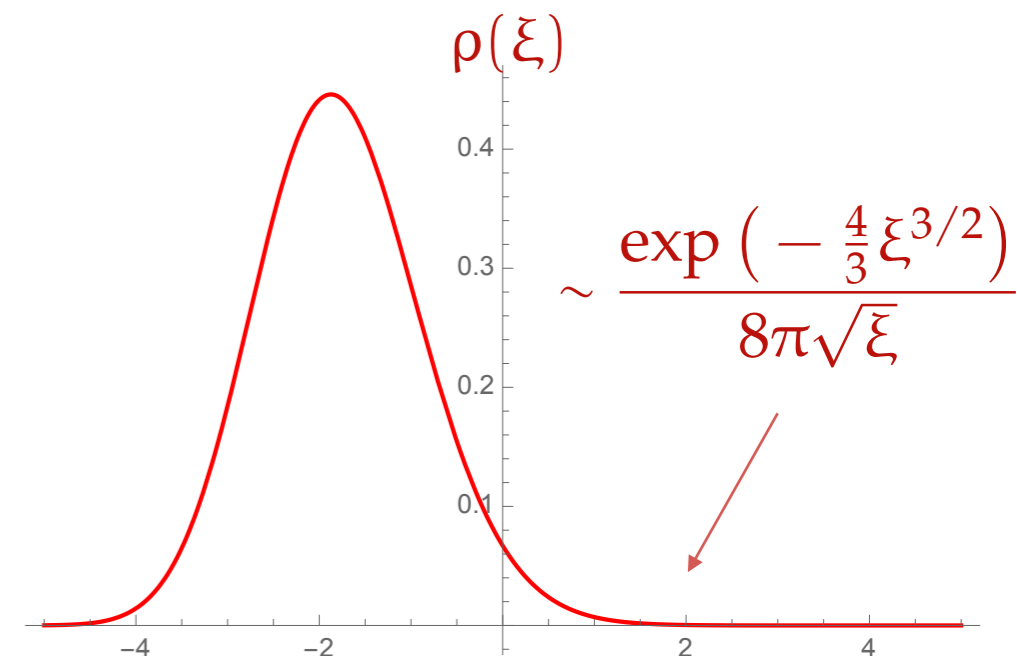
## Fluctuations of the maximum

Tracy-Widom 1992

$$E_1 = \max(E_1, \dots, E_n) \underset{n \text{ large}}{\approx} 2\sigma\sqrt{n} + \sigma n^{-1/6} \xi$$

$\xi$  random with Tracy-Widom distribution

- Quite different from Gumbel's law !
- Universal !
- Describes as well fluctuations of  $\ell_1$  on vertical section of large random tilings

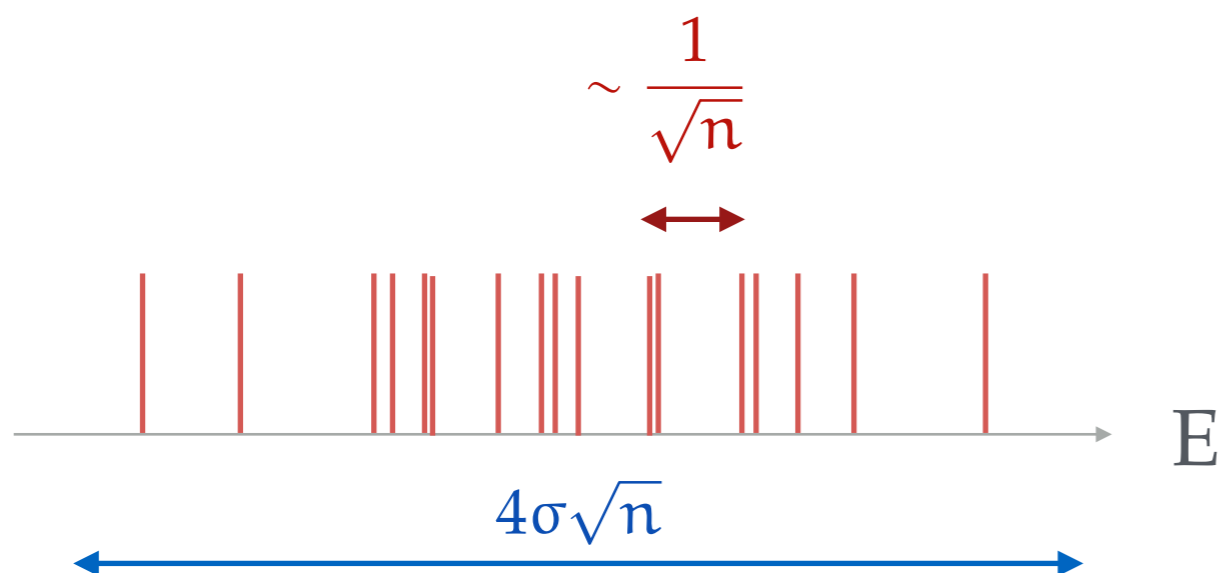


# 4. Eigenvalues of large random matrices

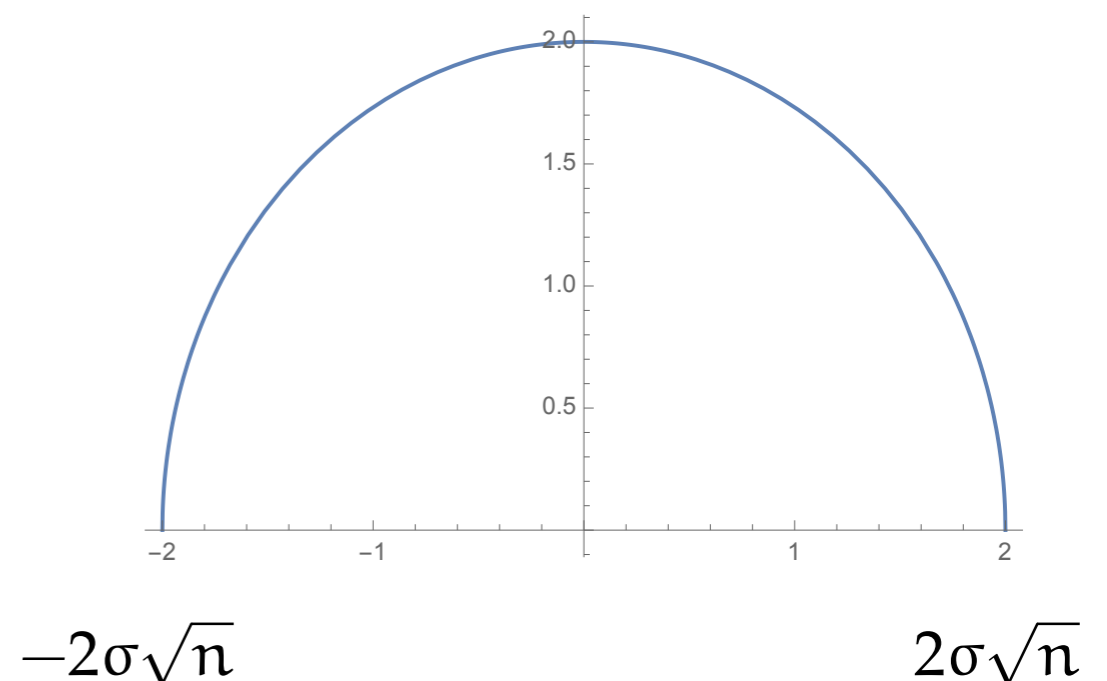
Consider a random (hermitian) matrix of size  $n \times n$  with independent Gauß-distributed entries with variance  $\sigma^2$

$n$  large : macroscopic distribution of eigenvalues becomes non-random but its shape is not universal (Wigner 1950)

*position of  $n$  eigenvalues*



*large  $n$  density of eigenvalues*

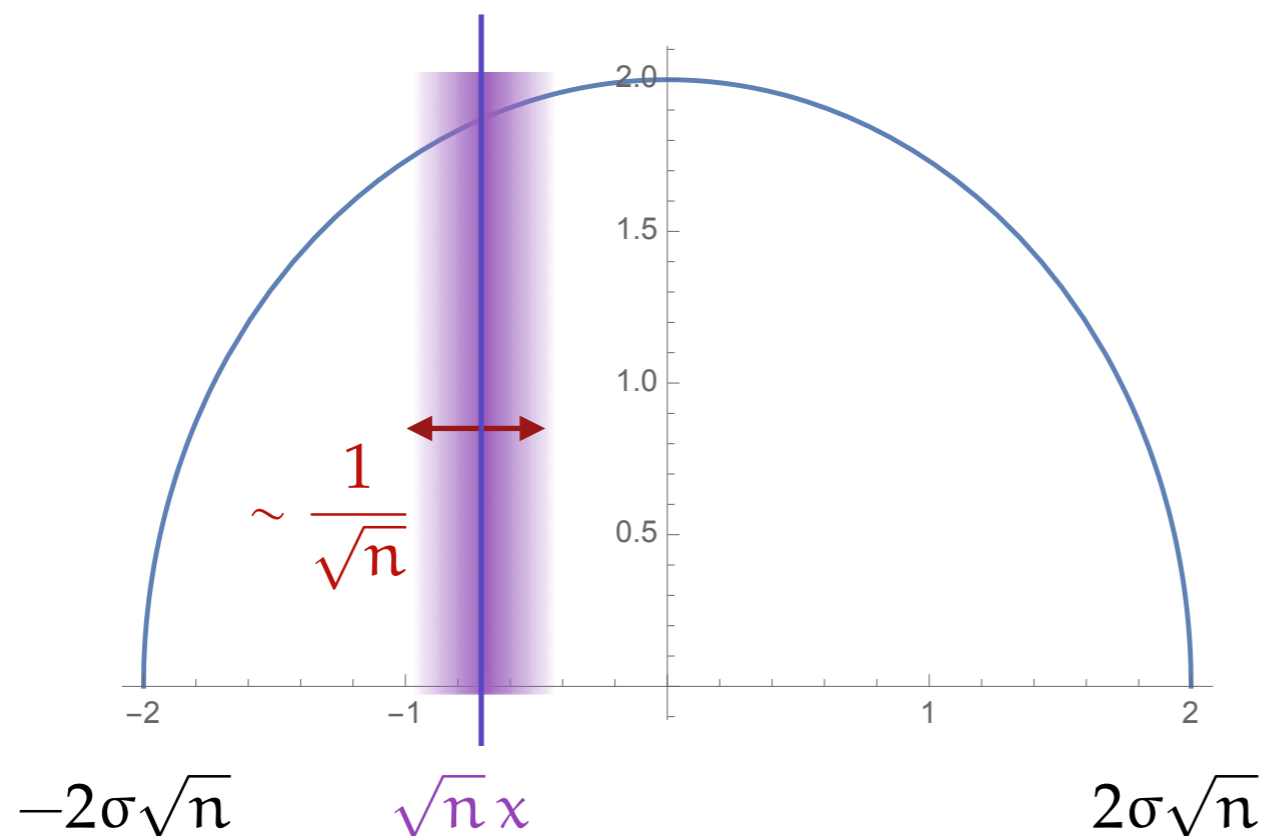


# 4. Eigenvalues of large random matrices

## Pair correlations

The probability distribution of finding a eigenvalue at  
at distance  $\frac{r}{\sqrt{n}}$  from an eigenvalue at  $\sqrt{n}x$

when  $n$  becomes large, is also known (Wigner, Dyson, Mehta 50s ...)



# 5

## Energy levels of heavy atomic nuclei

# 5. Heavy atomic nuclei

## In quantum mechanics

- the state of a system is represented by a point  
 $X(t) = (x_1(t), \dots, x_n(t))$  in n-dimensional space

- the time evolution is described by a matrix H

$$i\hbar \frac{X(t + \delta) - X(t)}{\delta} \underset{\delta \text{ small}}{\approx} H \cdot X(t) \quad (\text{Schrödinger's equation})$$

- H is not random: it is specified by the components of the system and their interactions

# 5. Heavy atomic nuclei

In general, understanding the evolution of the system  
= finding the states with fixed energy (eigenvectors of  $H$ )  
and the value of these energies (eigenvalues of  $H$ )

For heavy atomic nuclei (between 70 and 100 protons)  
 $H$  is very complicated !

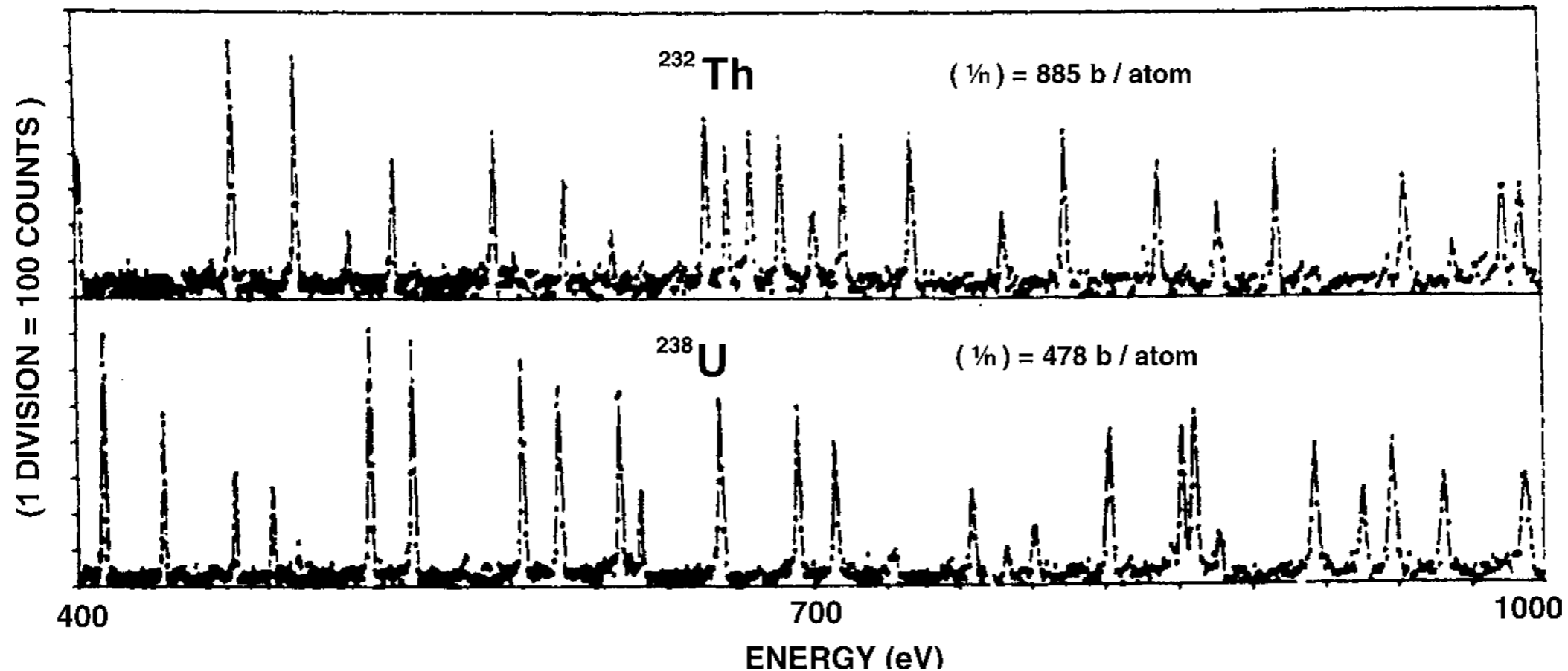
## **Wigner's idea (1950s)**

statistics of eigenvalues of  $H$  may look like the  
statistics of eigenvalues of a typical large matrix **chosen at random**

# 5. Heavy atomic nuclei

Eigenvalues of H for atomic nuclei

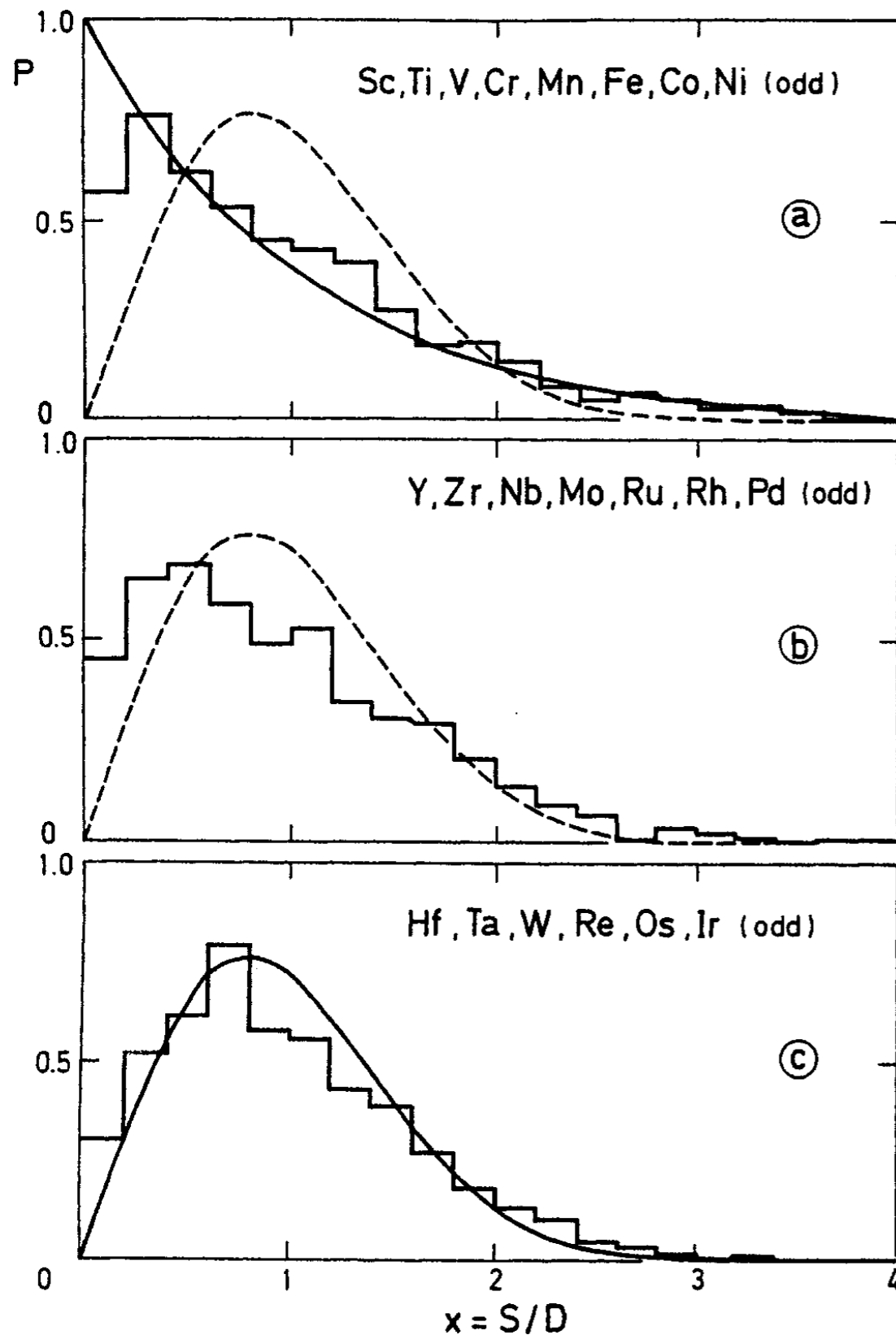
= energy of particles that it can absorb (resonances)



**Figure 1.1.** Slow neutron resonance cross-sections on thorium 232 and uranium 238 nuclei. Reprinted with permission from The American Physical Society, Rahn et al., Neutron resonance spectroscopy, X, *Phys. Rev. C* 6, 1854–1869 (1972). (from Mehta's book *Random matrices*)



# 5. Heavy atomic nuclei



## Pair correlations

- Cumulative histogram of distance between high energy resonances in certain heavy nuclei

vs.

- Predicted probability distributions for eigenvalues of a large random hermitian matrix

corresponding increase in strength of the spin dependent forces. Reprinted with permission from *Annales Academiae Scientiarum Fennicae*, Porter C.E. and Rosenzweig N., Statistical properties of atomic and nuclear spectra, *Annale Academiae Scientiarum Fennicae, Serie A VI, Physica 44*, 1-66 (1960).

(from Mehta's book *Random matrices*)

# 6

Other surprising apparitions  
of large random matrices

# 6. Other apparitions

## Summary

- **Universality phenomenon:** repulsive random variables can exhibit the same statistics as the eigenvalues of large random matrices
- **Proving universality** (for large classes of models of random matrices) has occupied mathematicians from the 60s until now, **is still an active topic of research**
- We have encountered applications to
  - random tilings
  - statistics of high energy resonances in heavy nuclei

# 6. Other apparitions

Many other apparitions of the theory of (large) random matrices

- Data analysis (economics, linguistics, phylogenetics, ...)
- Statistics of zeroes of the Riemann zeta function  
(related statistics of prime numbers)  
(Montgomery 1970)
- Random crystal growth / interface growth  
(Takeuchi-Sano 2010, Sasamoto-Spohn 2010)
- Statistics of distance between pine trees in Swedish forests  
(le Caer 1990)
- Statistics of bus waiting times in Cuernavaca (Mexico)  
(Krbalek-Seba 2000, Baik-Borodin-Deift-Suidan 2006)

*Thank you for your attention !*

