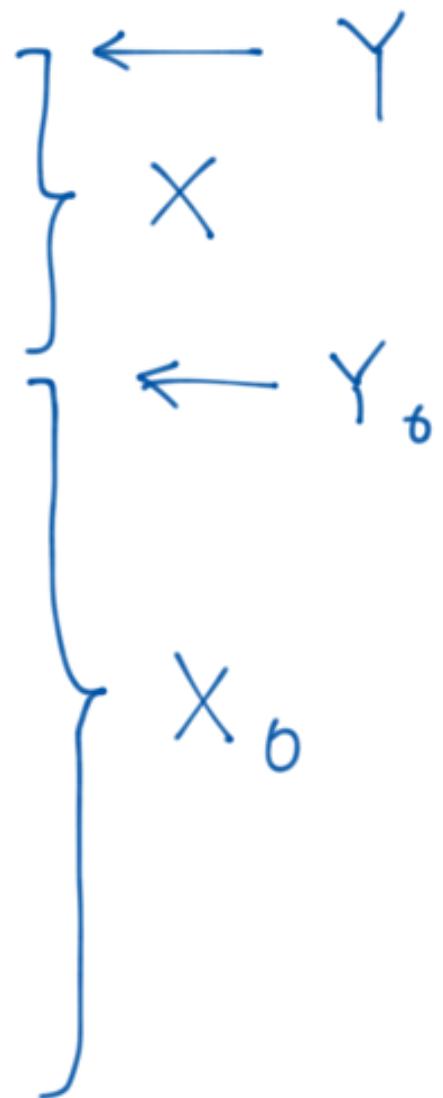
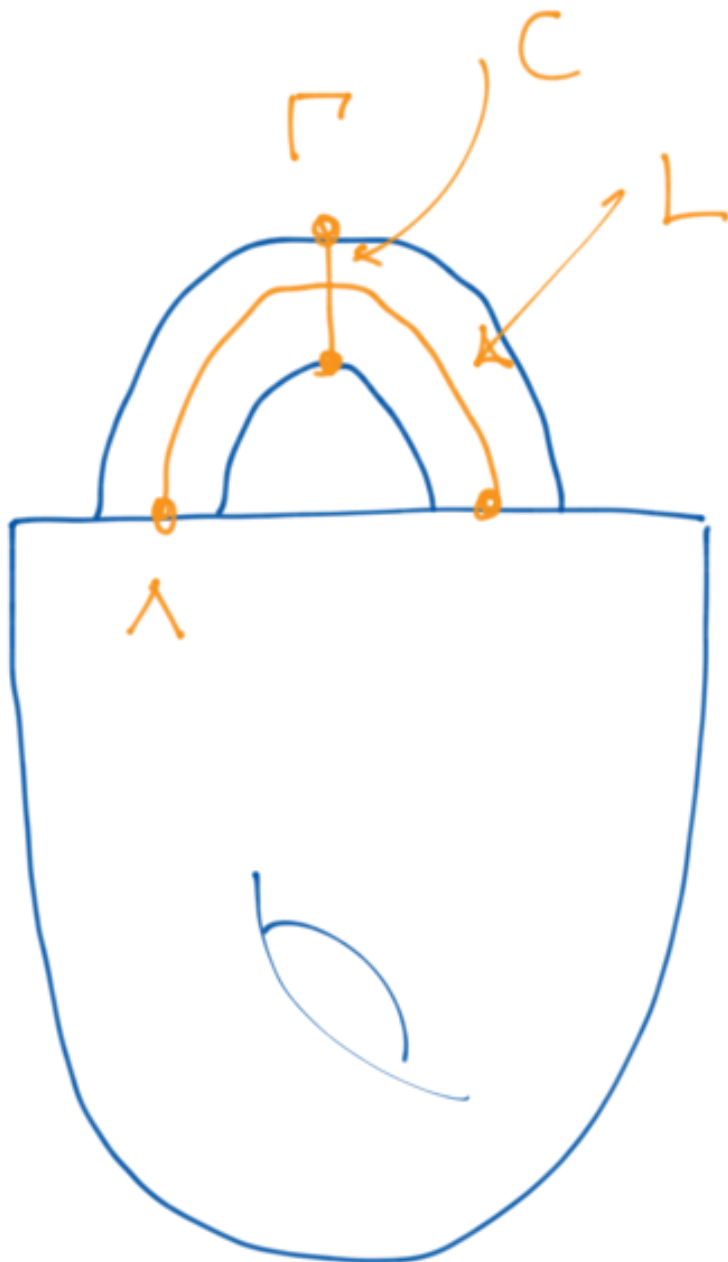


Lecture 1

Set up:



- **The main idea in these lectures is to express the symplectic cohomology of X and the wrapped Floer cohomology of C in terms of the Legendrian DGA of Λ .**

- **We will first focus on the linear chain complexes of symplectic cohomology and wrapped Floer cohomology and look at their product structures later.**

Contact homology

Y - contact $(2n-1)$ -manifold

α - contact 1-form

$\xi = \ker(\alpha)$ - contact plane field

R - Reeb vector field

$$d\alpha(R, \cdot) = 0; \quad \alpha(R) = 1$$

$\Lambda \subset Y$ - Legendrian submanifold.

For simplicity assume

$$c_1(\xi) = 0 \Rightarrow \det(\xi) \text{ trivial}$$

and that the Maslov
class $\mu \in H^1(\Lambda)$ vanishes

The orbit Differential Graded Algebra (DGA)

$Q(Y)$

is the unital graded algebra over the rational numbers generated by the good Reeb orbits in Y , graded by a shifted Conley-Zehnder index, and with differential which counts holomorphic spheres with one positive and several negative punctures.

γ - Reeb orbit

Pick complex trivialization \mathcal{T} of ξ along γ .

The linearized Reeb flow along γ defines a path of linear symplectomorphisms

$$\Psi_t : \mathbb{C}^{n-1} \rightarrow \mathbb{C}^{n-1}$$

The graph of Ψ_t is a path L_t of Lagrangian planes in $\mathbb{C}^n \oplus \mathbb{C}^{n-1}$. Let L'_t denote L_t closed up by the smallest positive rotation. The Conley-Zehnder index is

$$CZ_{\mathcal{T}}(\gamma) = \mu(L'_t) - \frac{n-1}{2}$$

where μ denotes the Maslov index (intersection number with the Maslov cycle).

A trivialization of $\det(\xi)$ determines a class of trivializations \mathcal{T} such that $\det(\mathcal{T}) \simeq \det(\xi)$, making CZ well-defined.

An orbit γ is called bad if it is an even multiple of another orbit β such that the parities of $CZ(\gamma)$ and $CZ(\beta)$ are different. (Note that parities are independent of complex trivializations.)

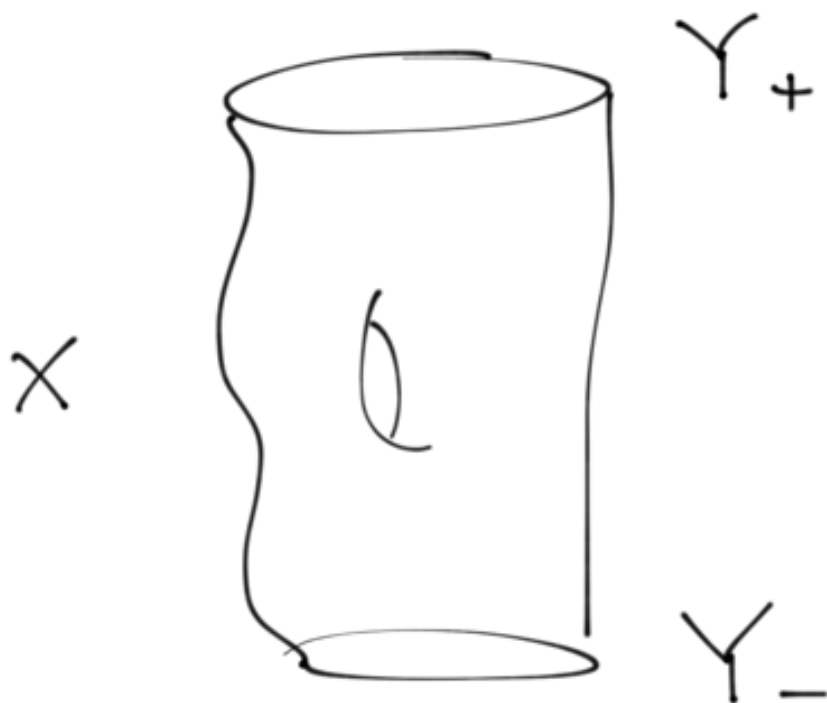
Orbits that are not bad are called good.

Take the contact form generic so that L_1 is transverse to the diagonal for each orbit.

$Q(Y)$ is generated by
all good Reeb orbits
 γ graded by

$$|\gamma| = CZ(\gamma) + (n-3).$$

Consider a Weinstein cobordism



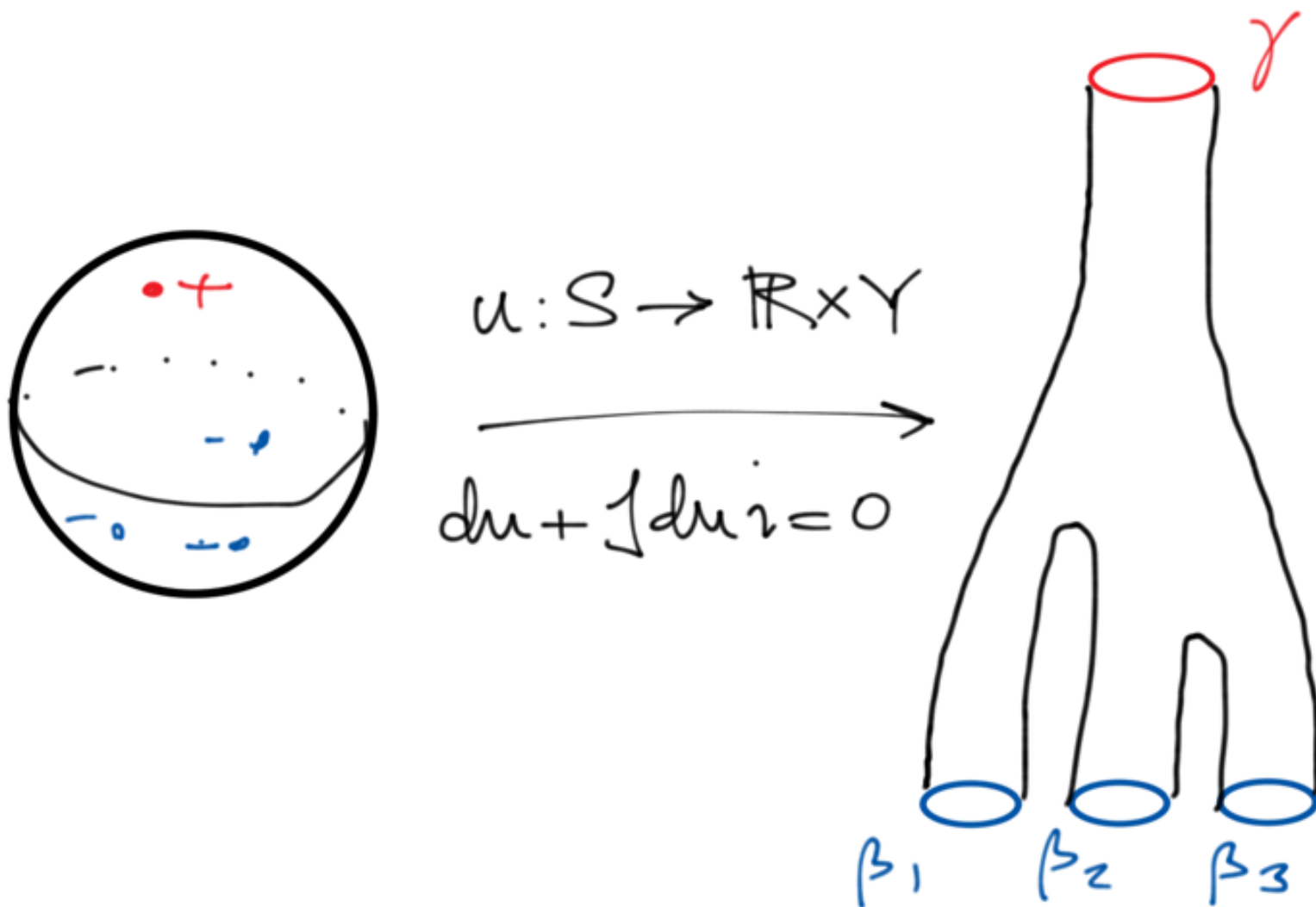
Pick an almost complex structure J on X which is adjusted to the contact structures in the positive and negative end.

$$J(\xi^\pm) \subset \xi^\pm ; \quad J(\partial_t^\pm) = \mathbb{R}^\pm$$

For γ a Reeb orbit in Y^+ and $\beta = \beta_1^{k_1} \beta_2^{k_2} \dots \beta_s^{k_s}$ a monomial of Reeb orbits in $\overline{Y^-}$ let

$$\mathcal{M}^X(\gamma, \beta)$$

denote the moduli space of holomorphic spheres:



The formal dimension of the moduli space is

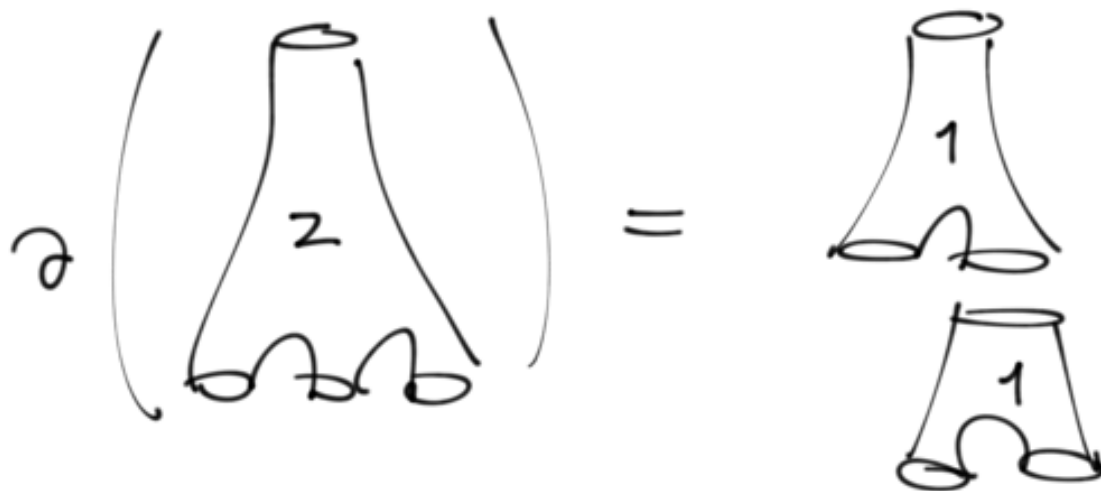
$$\begin{aligned} \dim(\mathcal{M}^X(\gamma, \beta)) &= |\gamma| - |\beta| = \\ &= |\gamma| - \sum_{j=1}^m k_j |\beta_j|. \end{aligned}$$

The differential $\partial: \mathcal{Q}(Y) \rightarrow \mathcal{Q}(Y)$ satisfies Leibniz rule and is defined on generators by the following curve count:

$$\begin{aligned} \partial(\gamma) &= \\ &= \sum_{|\gamma| - |\beta| = 1} |\mathcal{M}^{\mathbb{R} \times Y}(\gamma, \beta)| \frac{1}{k_1! \dots k_c!} \frac{1}{m(\beta_1)} \dots \frac{1}{m(\beta_s)} \beta \end{aligned}$$

where $m(\beta)$ is the multiplicity of β .

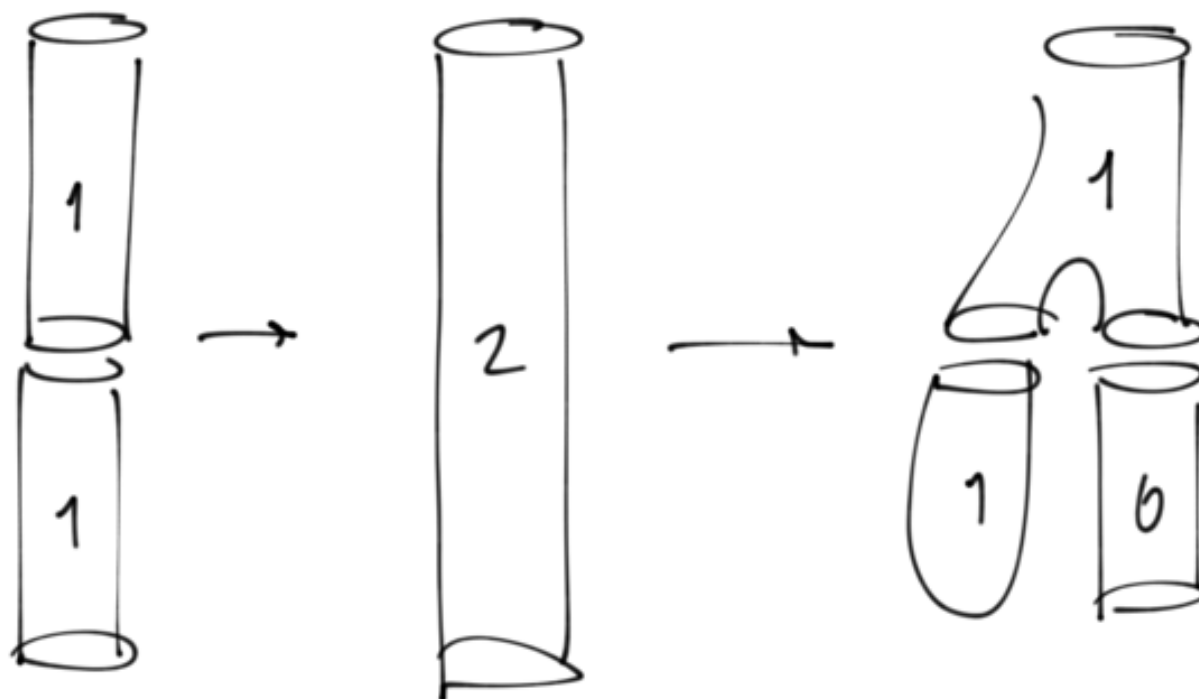
For a suitable perturbation scheme transversality holds and the boundary of the (reduced) 1-dim moduli space is by SFT-compactness and gluing in 1-1 correspondence with 2-level broken curves of total dimension 2:



This shows that ∂ is a differential:

$$\partial^2 = 0 .$$

Note in particular that there could be splittings of the form



This is the reason for a DGA rather than a linear chain complex.

Similarly the cobordism X defines a DGA chain map

$$\Phi_X : Q(Y^+) \longrightarrow Q(Y^-)$$

$$\Phi_X(\gamma) =$$

$$= \sum_{|\gamma| - |\beta| = 0} |M^X(\gamma; \beta)| \frac{1}{k_1! \dots k_s!} \frac{1}{K(\beta_1) \dots K(\beta_s)} \beta$$

Then $\Phi_X \circ \partial^+ = \partial^- \circ \Phi_X$
since



An augmentation is a DGA chain map

$$\Sigma: Q(Y) \rightarrow \mathbb{Q} \quad ;$$

$$\Sigma \circ \partial + \cancel{\partial \circ \Sigma}^0 = \Sigma \circ \partial = 0$$

If Y^- is empty then $Q(Y^-) = \mathbb{Q}$ and

$$\Sigma_X := \Phi_X: Q(Y^+) \rightarrow \mathbb{Q}$$

is an augmentation.

$$\partial \left(\text{cup} \right) = \text{cup} \Rightarrow \Sigma_X \circ \partial = 0$$

The Legendrian DGA

$\hat{A}(\Lambda)$

is a unital non-commutative DGA generated by Reeb chords of Λ , graded by Maslov index with differential that counts disks with one positive and several negative boundary punctures.

In general the chord algebra is an algebra over $\mathbb{Q}(Y)$, but an augmentation of $\mathbb{Q}(Y)$ can be used to reduce the coefficients to \mathbb{Q} .

c - Reeb chord with initial point c^- and final c^+

Pick a path λ'_c in Λ connecting c^+ to c^-

Then $c * \lambda'_c$ is a loop λ_c . Pick a trivialization T of ξ along λ_c compatible with $\det(\xi)$.

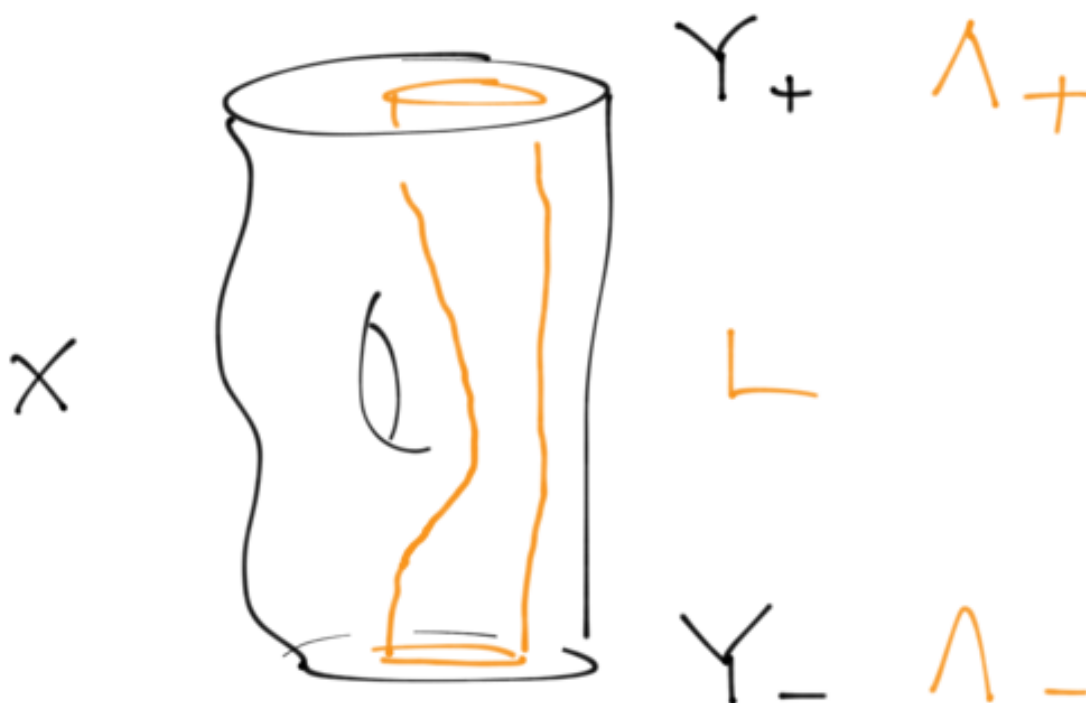
Consider the following loop of Lagrangian planes in ξ

$$L_c = T_{\lambda'_c} \Lambda * [d(e^{\mathbb{R}}) T_{c^-} \Lambda] * \text{Rot}^+$$

The grading of c is then

$$|c| = \mu(L_c) - 1.$$

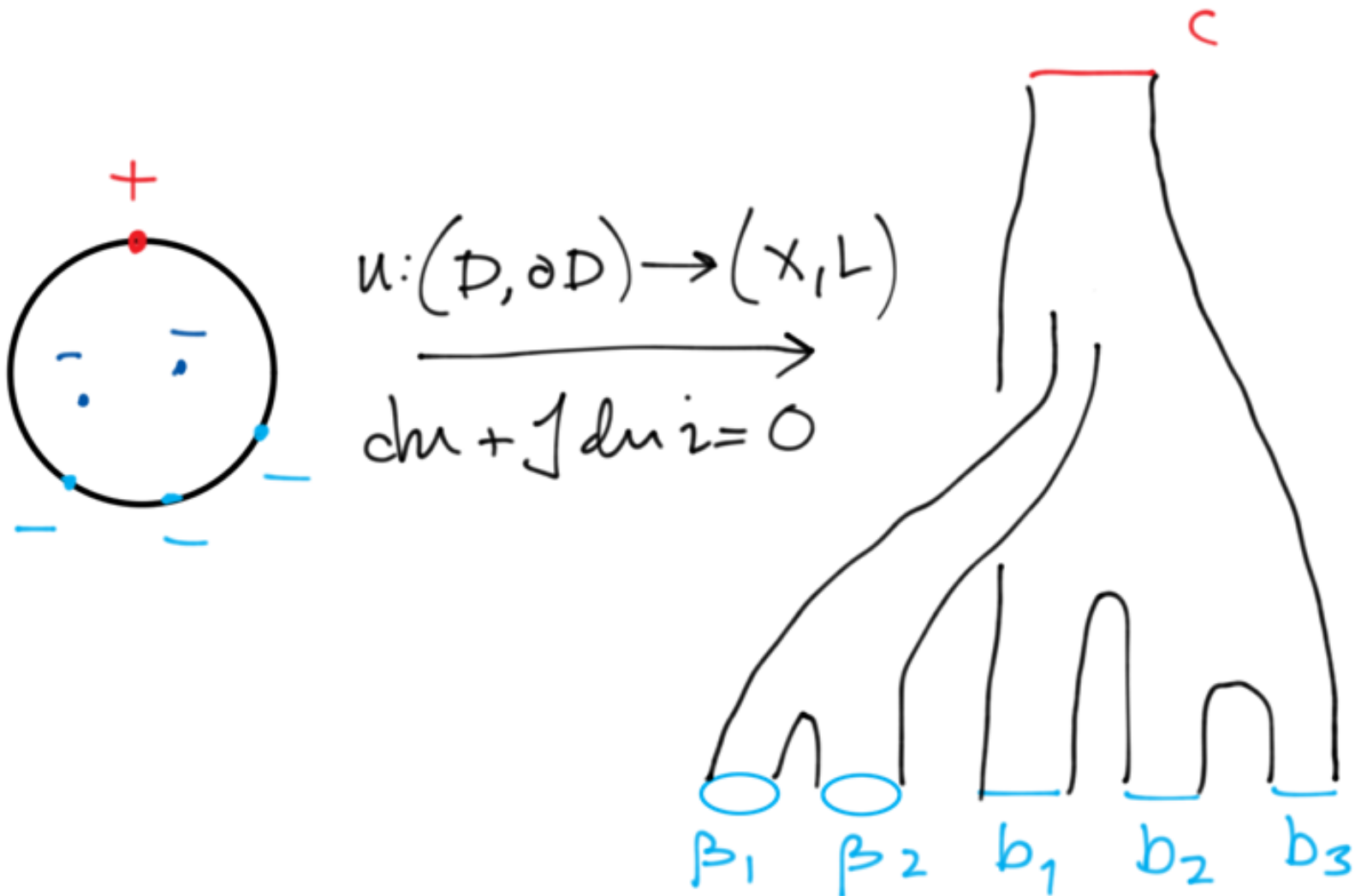
Consider a Weinstein cobordism with an exact Lagrangian cobordism inside



For c a Reeb chord of Λ^+ , for $\underline{b} = b_1 b_2 \dots b_r$, a monomial of Reeb chords of Λ^- , and for $\underline{\beta}$ a monomial of Reeb chords, let

$$\mathcal{M}^{(X, L)}(c, \underline{b}; \underline{\beta})$$

denote the moduli space of holomorphic disks:



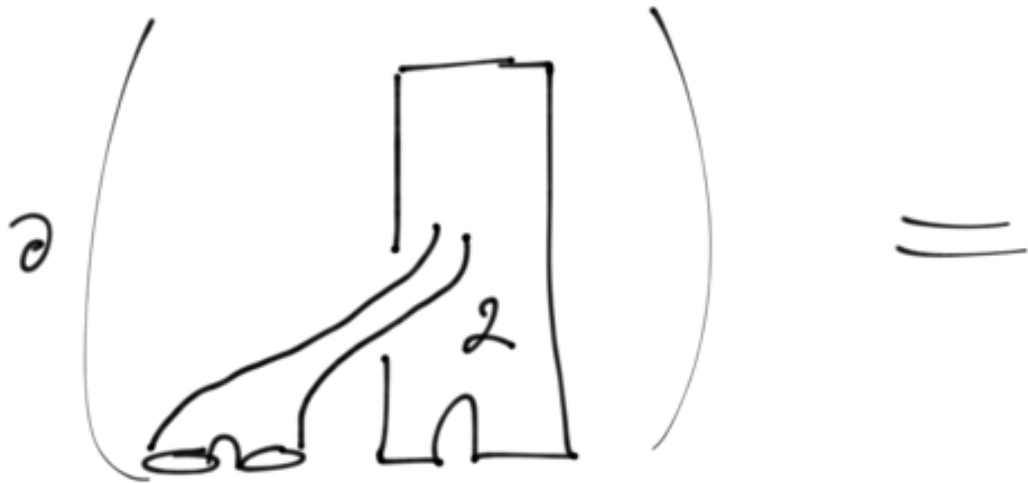
The formal dimension of the moduli space is

$$\dim(\mathcal{M}^{(X,L)}(c, \underline{b}; \underline{\beta})) = \\ = |c| - |\underline{b}| - |\underline{\beta}|.$$

The differential $\partial: A(\Lambda) \rightarrow A(\Lambda)$ satisfies Leibniz rule and is defined on generators by the following curve count:

$$\partial(c) = \\ = \sum_{|c| - |\underline{b}| - |\underline{\beta}| = 1} |\mathcal{M}^{(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)}(c, \underline{b}; \underline{\beta})| \frac{1}{k_1! \dots k_s!} \frac{1}{K(\beta_1)} \dots \frac{1}{K(\beta_s)} \underline{\beta} \underline{b}$$

This gives a differential:



$$\Rightarrow \partial^2 = 0 .$$

Similarly, the exact Lagrangian cobordism (X, L) defines a DGA map

$$\Phi_{(X, L)} : A(\Lambda^+) \rightarrow A(\Lambda^-)$$

$$\Phi_{(X, L)}(c) =$$

$$= \sum_{|c| - |b| - |\beta| = 0} |\mathcal{M}^{(X, L)}(c, b; \beta)| \frac{1}{k_1! \dots k_s!} \frac{1}{K(\beta_1)} \dots \frac{1}{K(\beta_s)} \overline{f}^\beta \underline{b}$$

$$\Phi_{(X, L)} \circ \partial^+ = \partial^- \circ \Phi_{(X, L)}$$

As before an augmentation is a chain map to \mathbb{Q}

$$\varepsilon: A(\Lambda) \longrightarrow \mathbb{Q}$$

$$\varepsilon \circ \partial = 0$$

and if (Y^-, Λ^-) is empty then

$$\varepsilon_{(X, L)} := \Phi_{(X, L)}$$

is an augmentation

The diagram shows the boundary of a 1-handle with a cap, represented as $\partial \left(\begin{array}{c} c \\ \text{1-handle} \\ 1 \end{array} \right)$, which is equal to a surface with three boundary components, represented as $\begin{array}{c} \text{1-handle} \\ 1 \\ \text{3 boundary components} \end{array}$. This leads to the equation $\partial \circ \varepsilon_{(X, L)} = 0$.

Examples

Standard contact S^{2n-1}

$$S = \left\{ |z_1|^2 + \frac{1}{a_2} |z_2|^2 + \dots + \frac{1}{a_n} |z_n|^2 = 1 \right\}$$

$0 < a_j \ll 1$; a_j 's lin indep over \mathbb{Q} .

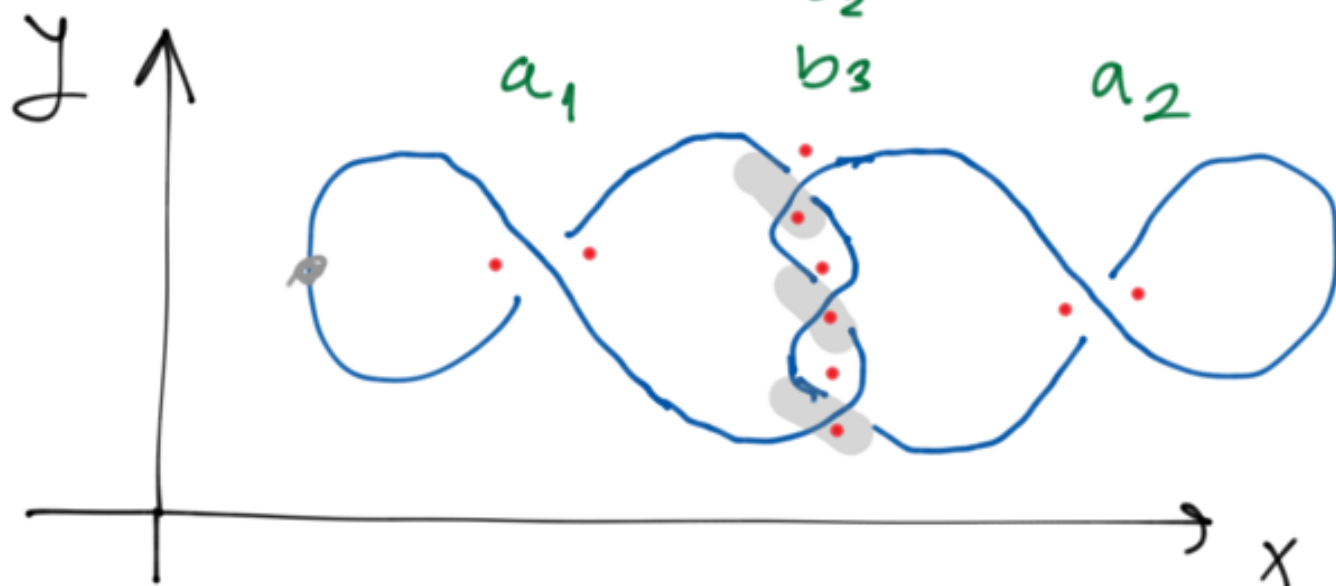
\Rightarrow effectively only one geometric orbit in $z_2 = \dots = z_n = 0, \gamma$

Orbit	γ	$\gamma^{(2)}$...	$\gamma^{(m)}$...
$ \cdot $	$2n-2$	$2n$		$2(n+m-1)$	

$$\Rightarrow H\mathbb{Q}(Y) = \mathbb{Q}[\gamma^{(j)}]_{j=1,2,\dots}$$

Legendrian trefoil in standard contact 3-space

$$\alpha = dz - y dx$$



$$|a_1| = |a_2| = 1$$

$$|b_1| = |b_2| = |b_3| = 0$$

$$\partial a_1 = -1 - b_1 - b_3 \rightarrow b_3 b_2 b_1$$

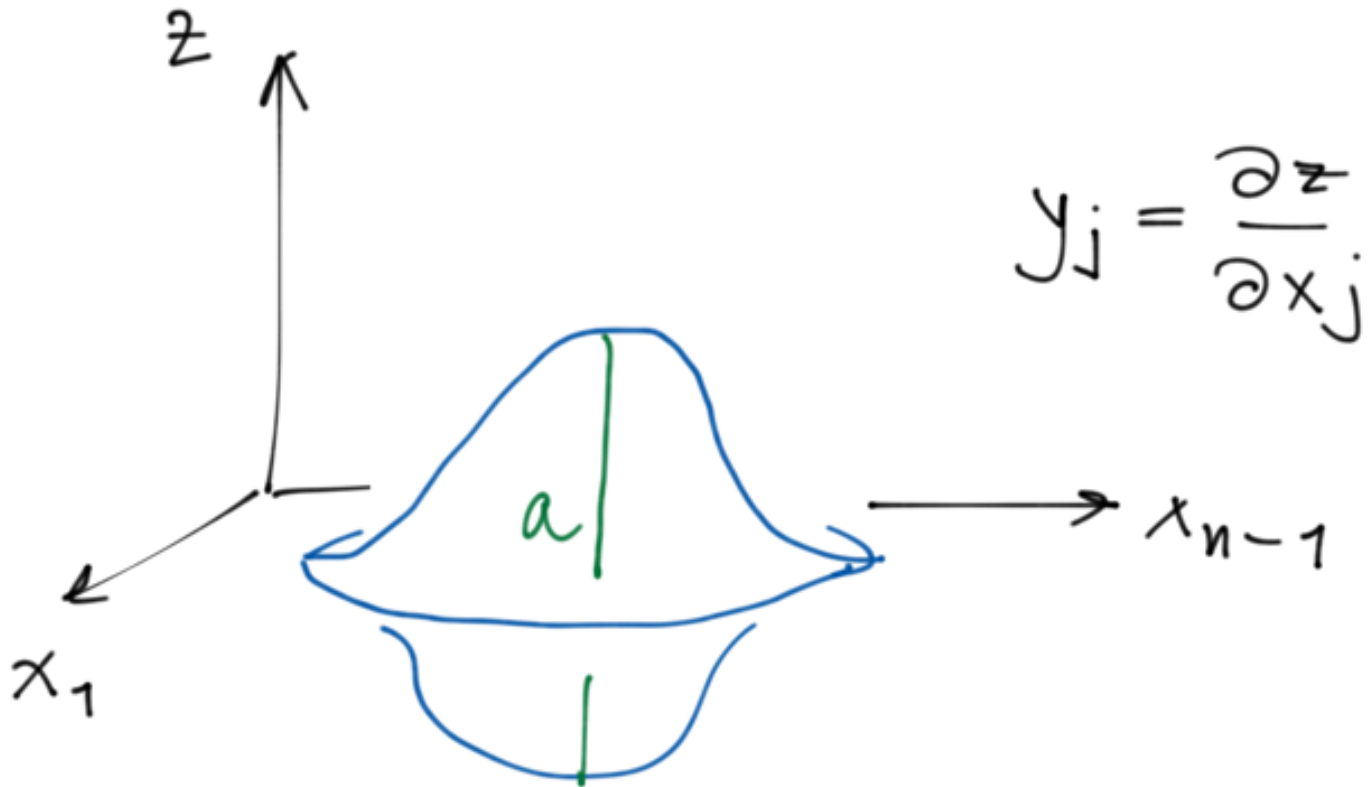
$$\partial a_2 = 1 + b_1 + b_3 + b_1 b_2 b_3$$

$$HA = \mathbb{Q}[b_1, b_2, b_3] / 1 + b_1 + b_3 + b_1 b_2 b_3$$

(Note that b_j 's commute in homology.)

Legendrian unknot of dimension $n-1$

$$S^{n-1} \text{ in } \mathbb{R}_{st}^{2n-1} \quad \alpha = dz - y_j dx_j$$



$$|a| = n-1$$

$$HA(\Lambda) = \mathbb{Q}[a] .$$

Augmentations

$$\varepsilon: Q(Y) \longrightarrow \mathbb{Q}, \quad \varepsilon \circ \partial = 0.$$

Linearized contact homology

Change variables in $Q(Y)$:

$$\psi(\gamma) = \gamma - \varepsilon(\gamma); \quad d = \psi^{-1} \circ \partial \circ \psi$$

$$Q^0 \subset Q^1 \subset Q^2 \subset \dots \subset Q^j \subset \dots \subset Q(Y)$$

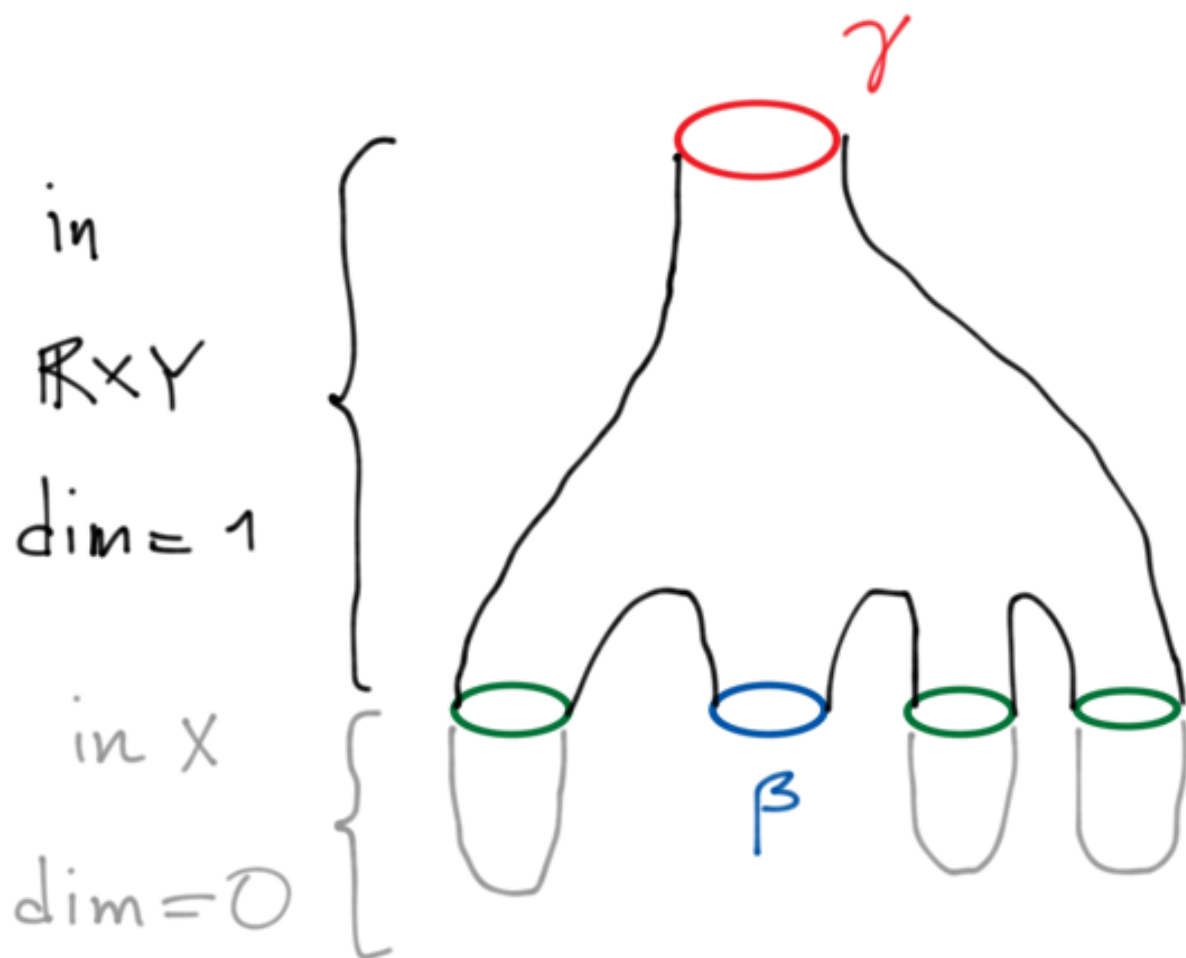
Q^j - polynomials of degree $\leq j$.

The differential d respects the degree filtration and the homology of the induced differential on Q^1 / Q^0 is called ε -linearized contact homology.

$$(C^{\text{lin}}(Y), d^{\text{lin}}) := (Q^1 / Q^0, \partial_1)$$

When $\varepsilon = \varepsilon_x$ the linearized differential has a geometric interpretation as a count of (partial) two-level curves:

$d^{\text{lin}} \gamma$ counts



Similarly, ϵ allows us to filter the Legendrian DGA by orbit degree and get a Legendrian DGA involving chords only.

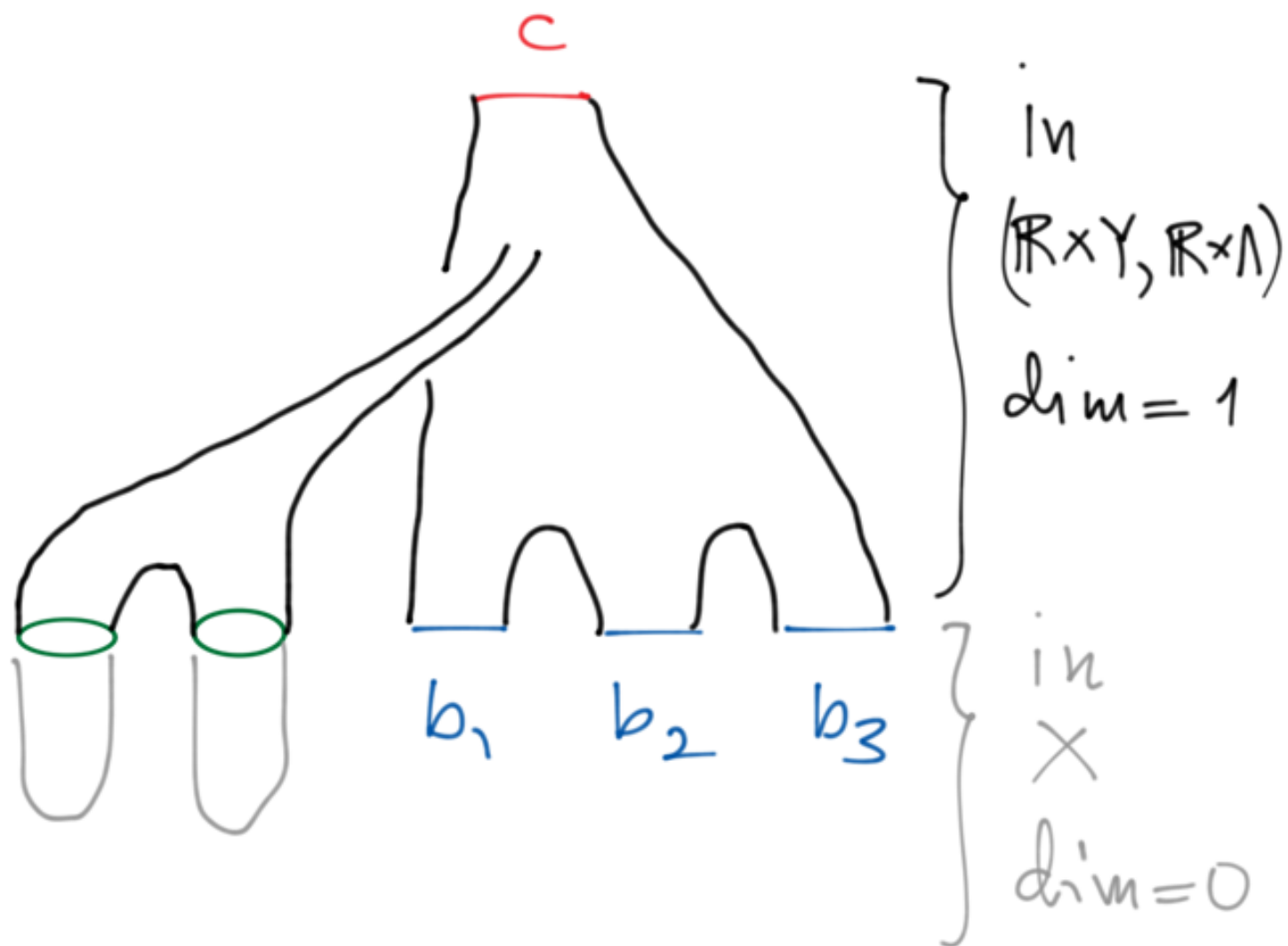
$$A(\Lambda) = \mathbb{Q} \langle \text{Reeb chords} \rangle.$$

$$d(c) =$$

$$\sum_{|c|-1, |b|=1} |\mathcal{M}^{\mathbb{R} \times Y}(c, \underline{b}; \underline{\beta})| \frac{1}{k_1! \dots k_s!} \frac{1}{K(\beta_1)} \dots \frac{1}{K(\beta_n)} \epsilon(\underline{\beta}) \underline{b}$$

When $\epsilon = \epsilon_x$ there is the following geometric interpretation:

dc counts (partial)
 d -level curves



If the Legendrian DGA admits an augmentation ϵ then we similarly get a linearized Legendrian DGA as the first page of a spectral sequence calculating the Legendrian contact homology.

$$\psi(c) = c - \Sigma(c) \quad ; \quad d = \psi^{-1} \partial \psi$$

$$A^0 \subset A^1 \subset \dots \subset A^j \subset \dots \subset A(\lambda)$$

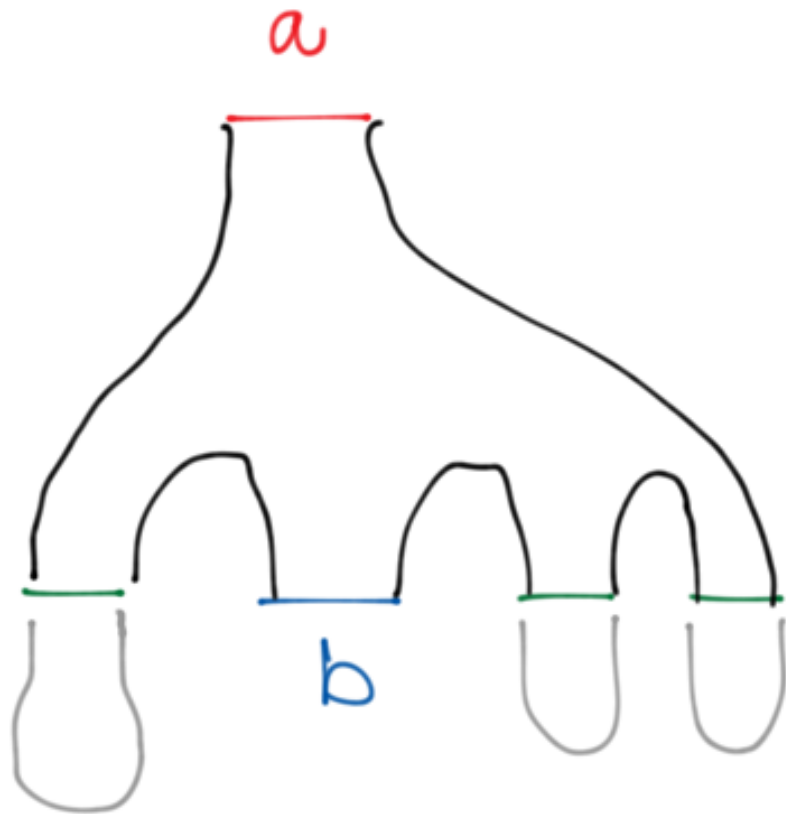
$$(C^{\text{lin}}(\Lambda), d^{\text{lin}}) = (A^1/A_0, d_1)$$

$$d^{\text{lin}}(c) =$$

$$\sum_{|c| - |b| = 1} |M^{(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)}(c, \underline{b}^l, b, \underline{b}^r)| (\Sigma(\underline{b}^l) \Sigma(\underline{b}^r)) b$$

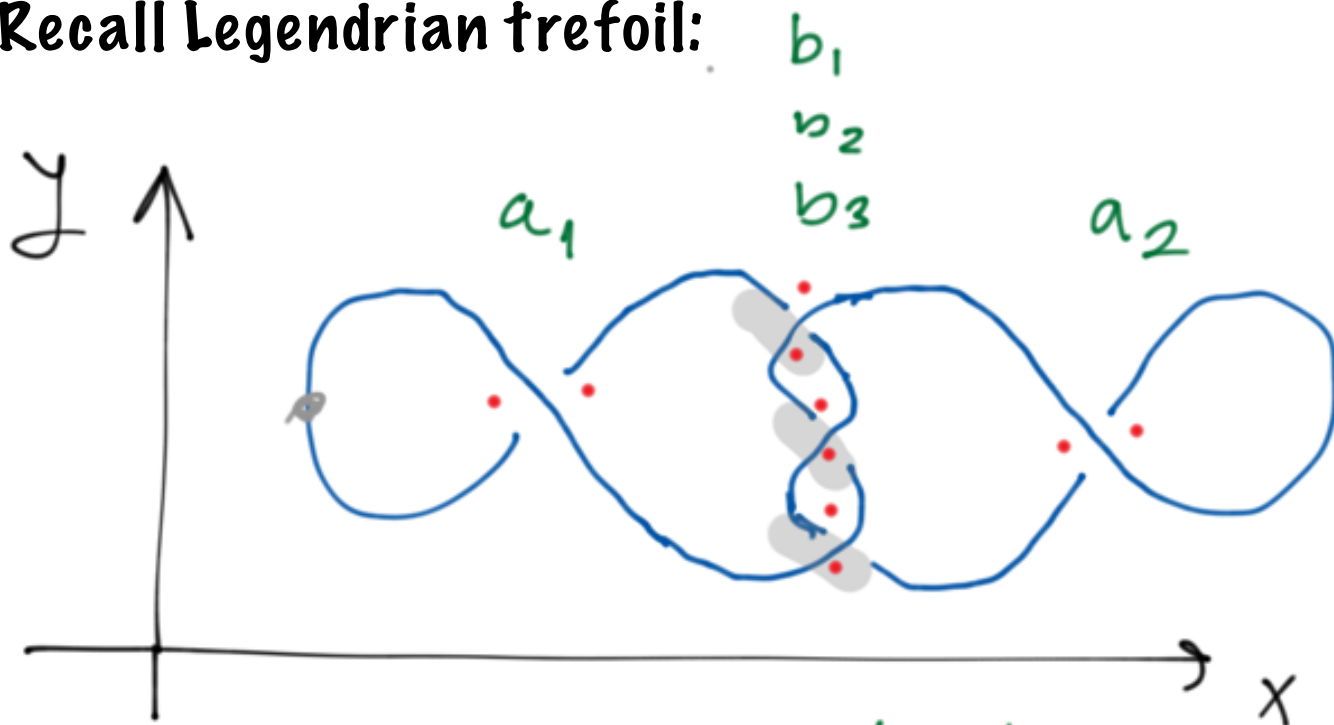
If the augmentation $\epsilon = \epsilon_L$ is induced by an exact Lagrangian we get the following geometric interpretation:

$d^{\text{lin}}(c)$ counts



$\left. \begin{array}{l} I_n \\ (\mathbb{R} \times Y, \mathbb{R} \times \lambda) \\ \dim = 1 \end{array} \right\}$
 $\left. \begin{array}{l} I_n \\ (X, L) \\ \dim = 0 \end{array} \right\}$

Recall Legendrian trefoil:



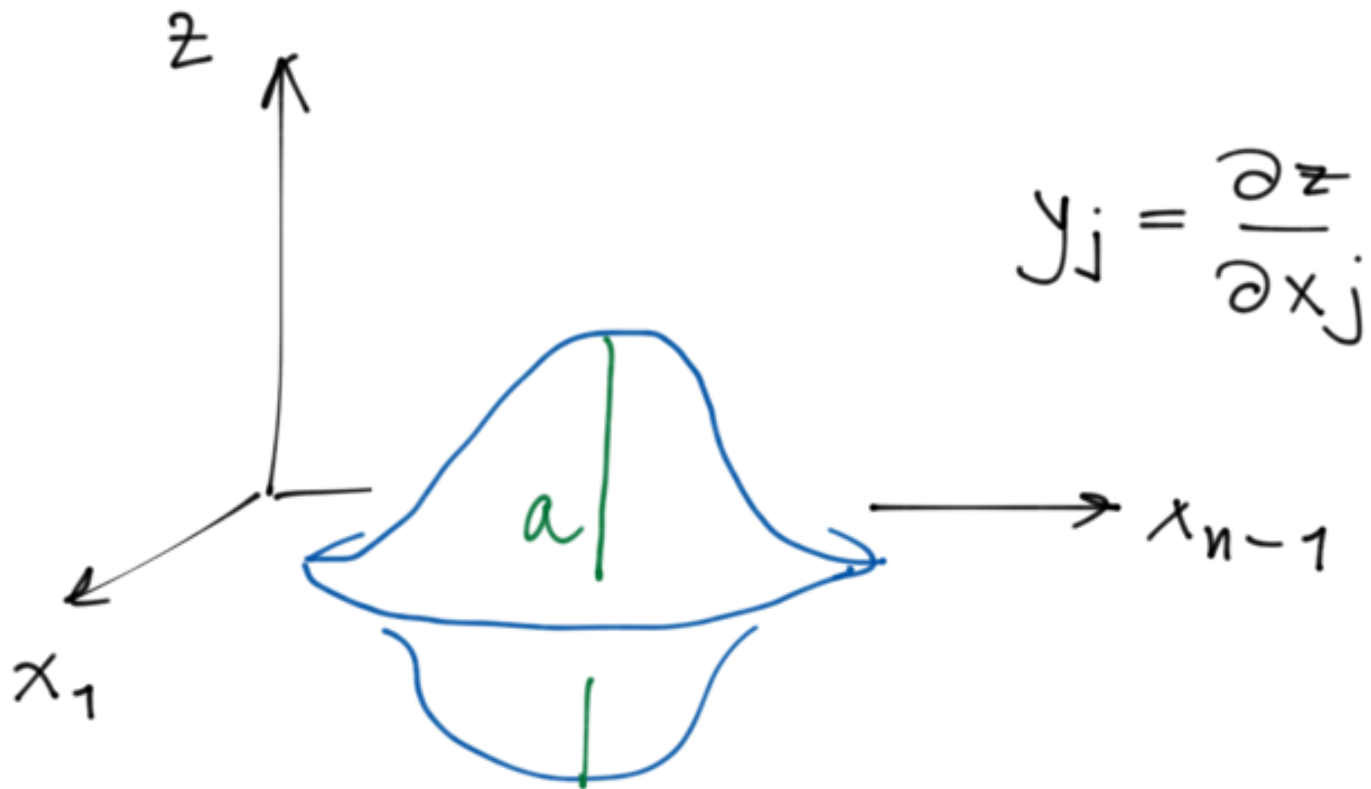
$$\partial a_1 = -1 - b_1 - b_3 \rightarrow b_3 b_2 b_1$$

$$\partial a_2 = 1 + b_1 + b_3 + b_1 b_2 b_3$$

	Σ_1	Σ_2	Σ_3	Σ_4	Σ_5
b_1	1	0	1	0	1
b_2	0	0	1	1	1
b_3	0	1	0	1	1

$$C\#_{*}^{\text{lin}} = \begin{cases} \mathbb{Q}, & * = 1 \\ \mathbb{Q}^2, & * = 0 \\ 0, & \text{otherwise} \end{cases}$$

Recall Legendrian unknot of dimension $n-1$:



$$|a| = n-1, \quad \partial a = 0$$

$$CH_*^{\text{lin}} = \begin{cases} \mathbb{Q}, & * = n-1 \\ 0, & \text{otherwise} \end{cases}$$

Linearized Legendrian homology and 2-copies

$\Lambda \subset Y$, neighborhood

$$N(\Lambda) \cong U \subset j^1(\Lambda)$$

$$\alpha \longleftarrow dz - p dq$$

$$\Lambda_0 := \Lambda$$

$$\Lambda_1 := j^1(f)$$

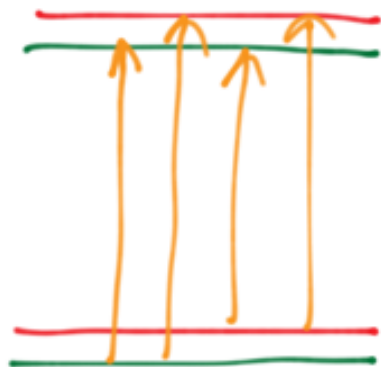
$f: \Lambda \rightarrow \mathbb{R}$ small positive Morse fctn.

$$\text{Ch}(\Lambda_0 \cup \Lambda_1) =$$

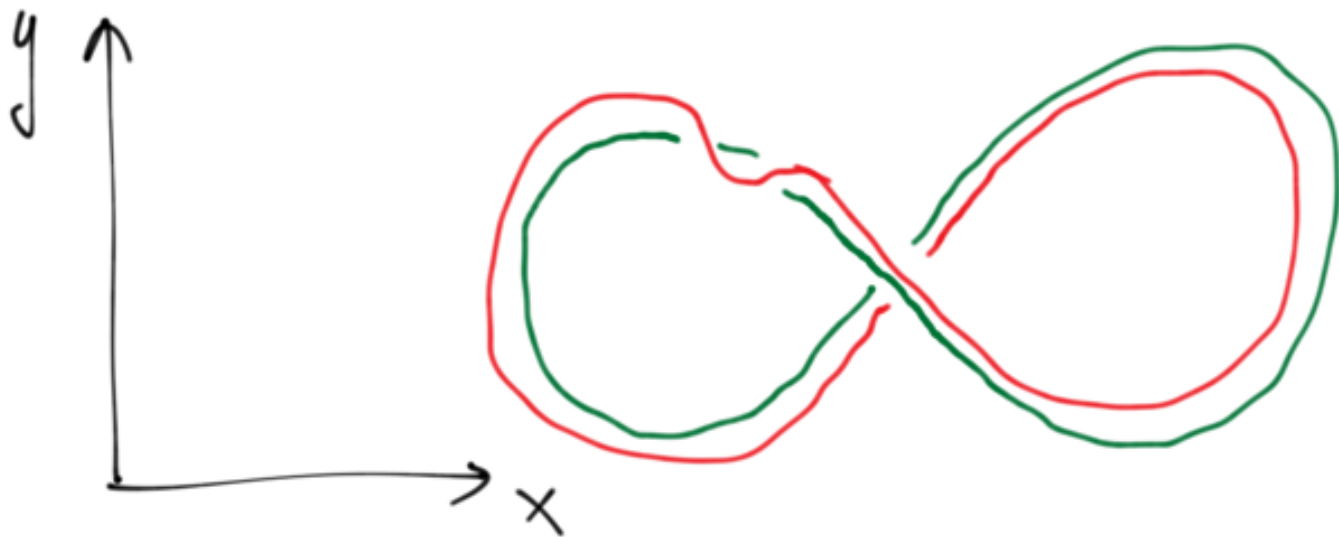
$$C_{00} \cup C_{11} \cup C_{01} \cup C_{10}$$

$$\cup \text{Crit}(f).$$

$$C_{ij} \approx \text{Ch}(\Lambda)$$

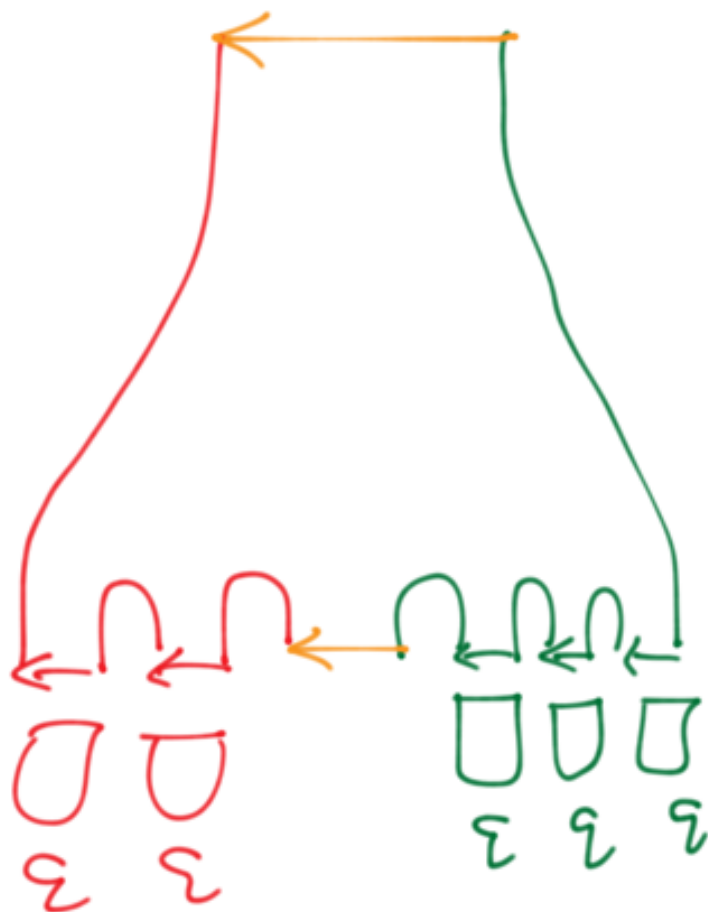


Example: $\alpha = dz - y dx$

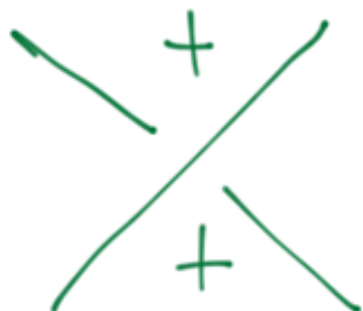


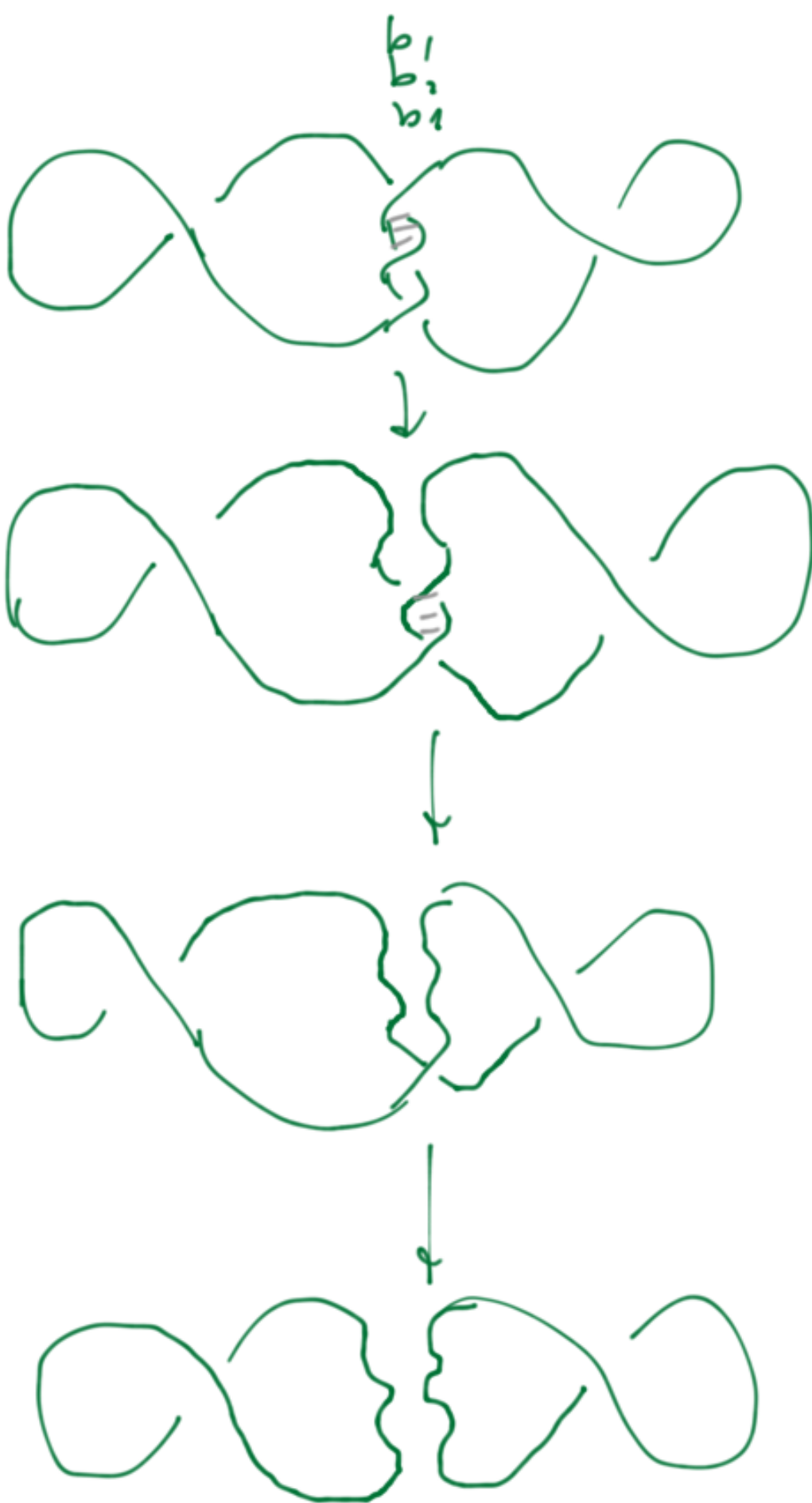
The differential of the 2-copy respects the filtration given by degree in mixed chords. To reproduce the linearized Legendrian homology we count disks with one mixed 10 chord:

$$d_{10} : C^{10} \rightarrow C^{10} \quad \text{counts}$$



The tori can be drawn explicitly and their chain maps computed via flow trees.





b_1	b_2	b_3
↓	↓	↓
1	b_2+1	b_3
↓	↓	↓
1	0	b_3+1
	↓	
1	0	0